



We still believe in supersymmetry

You must be joking

A Higgs Boson at 96 GeV?!

Sven Heinemeyer, IFT/IFCA (CSIC, Madrid/Santander)

Hong Kong, 01/2020

- Motivation & the “excesses”
- General analysis
- SUSY realizations
- Conclusions

1. Motivation & the “excesses”

Two Facts:

We have an SM-like Higgs discovery!

The SM cannot be the ultimate theory!

Conclusion: It cannot be “the SM Higgs”!

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Q': Which model?

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Conclusion: It cannot be “the SM Higgs”!

Q: Does the BSM physics have any (relevant) impact on the Higgs?

Q': Which model?

A1: check changed properties

A2: check for additional Higgs bosons

A2': check for additional Higgs bosons above and below 125 GeV

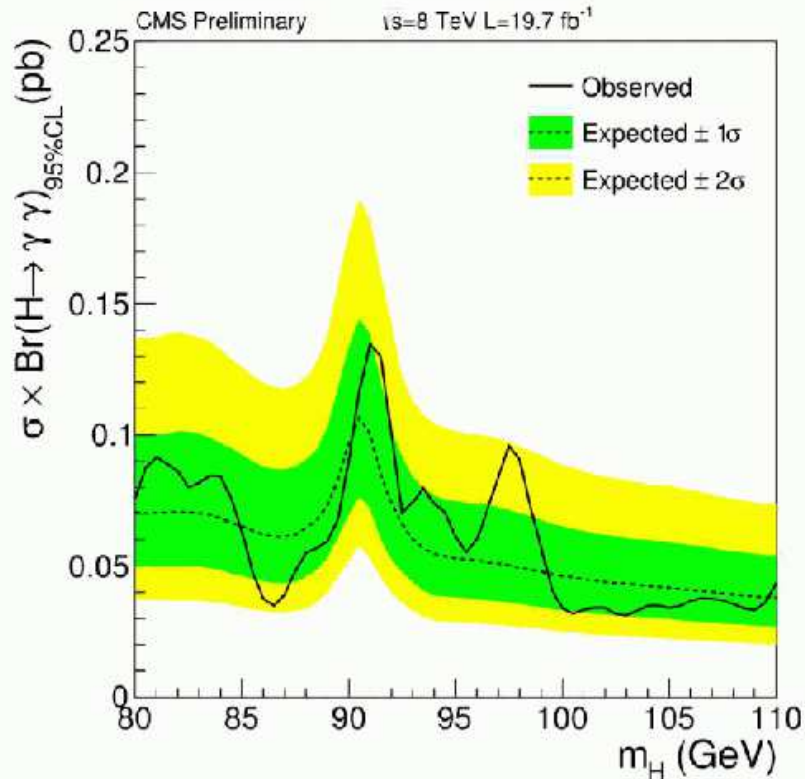
⇒ some results of searches below 125 GeV

$h \rightarrow \gamma\gamma$ (65-110 GeV) Run 1

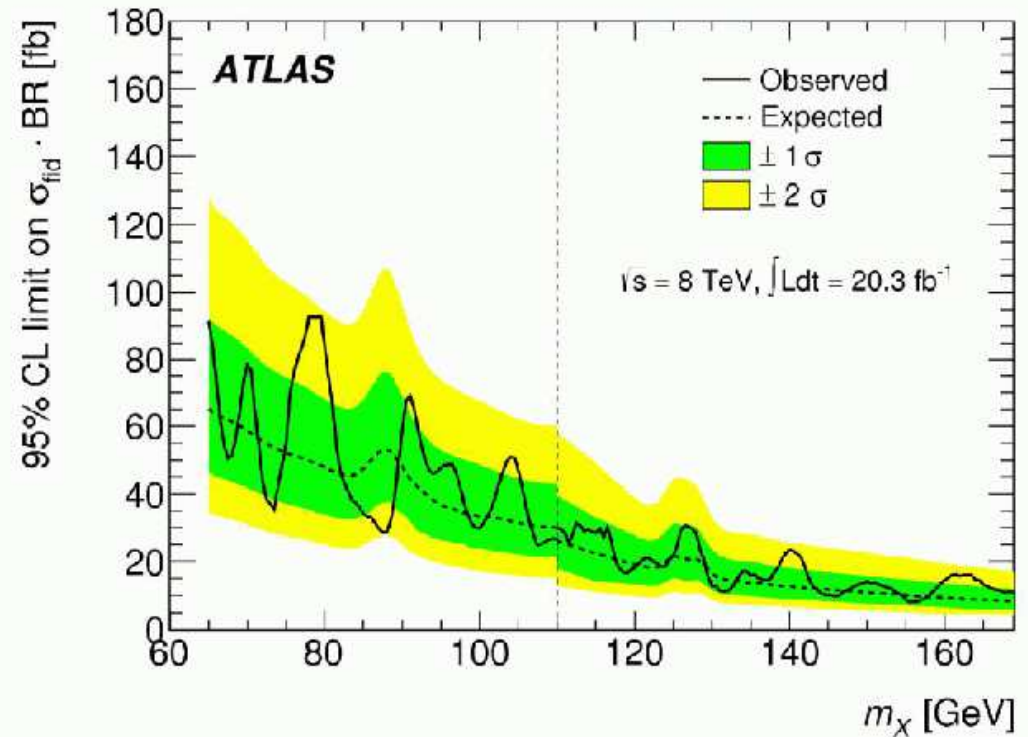


CMS PAS HIG-14-037

PRL 113 171801 (2014)



• $\sim 2\sigma$ excursion @ ~ 97.5 GeV



• $\sim 2\sigma$ excursion @ ~ 80 GeV

S. Gascon-Shotkin HDays17, Santander, ES Sept. 22 2017

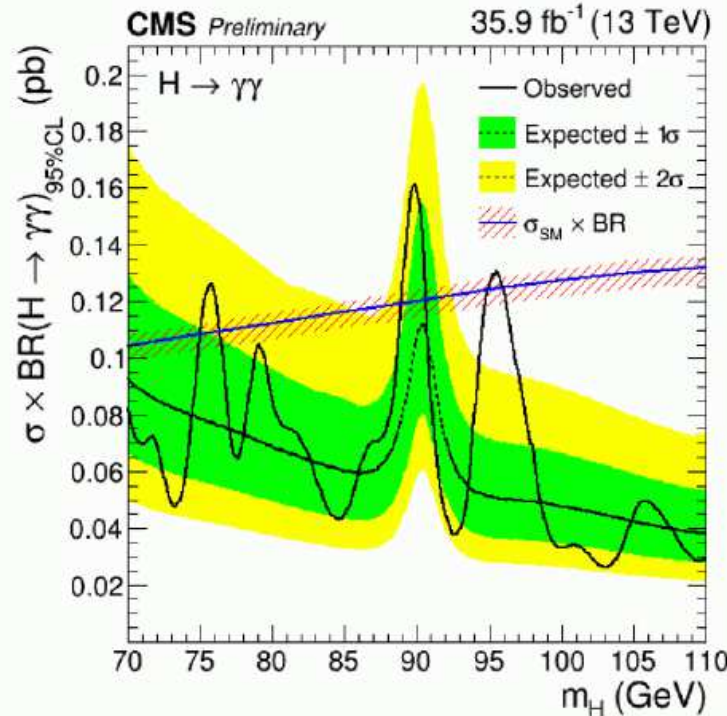
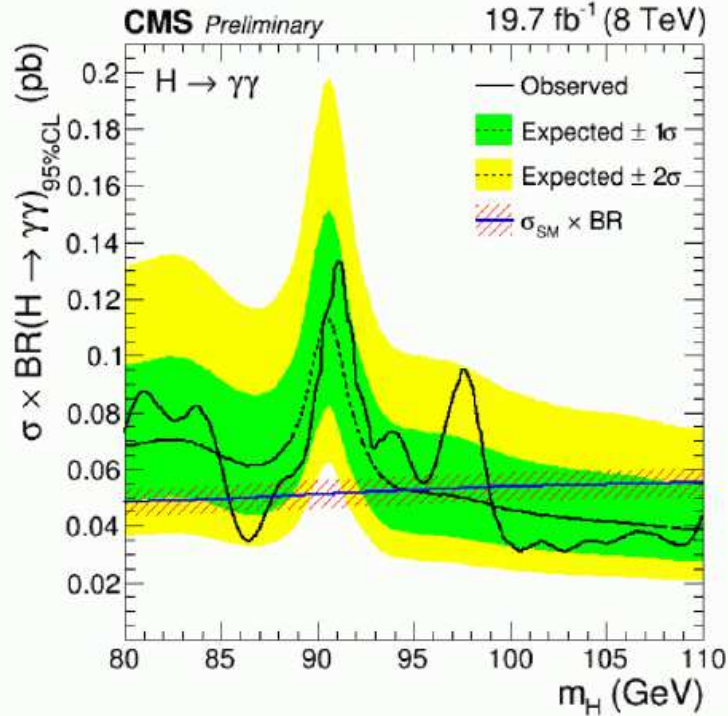
18



$h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+2



CMS PAS HIG-17-013



8 TeV:
 minimum(maximum)
 limit on $\sigma \times \text{Br}$:
 31(133) fb at
 $m=102.8(91.1)\text{GeV}$

13 TeV:
 minimum(maximum)
 limit on $\sigma \times \text{Br}$:
 26(161) fb at
 $m=103.0(89.9)\text{GeV}$

- 8 TeV limits on $\sigma \times \text{Br}$ redone with 0.1 GeV step. Production processes assumed in SM proportions. No significant excess with respect to expected limits observed.

S. Gascon-Shotkin HDays17, Santander, ES Sept. 22 2017

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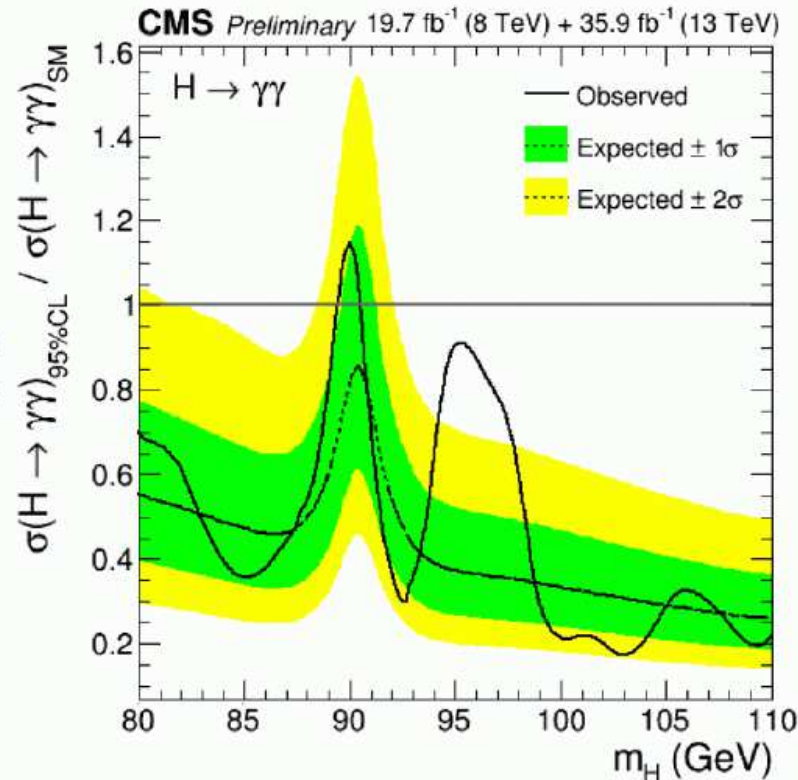


$h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+ 2



CMS PAS HIG-17-013

All experimental + theoretical systematic uncertainties assumed uncorrelated except for those on signal acceptance due to scale variations + those on production cross sections (assumed 100% correlated).



8 TeV+13 TeV:
 minimum(maximum) limit
 on $(\sigma \times Br) / (\sigma \times Br)_{SM}$:
 0.17(1.15) at
 $m=103.0(90.0)\text{GeV}$

- Combined 8 TeV+13 TeV $\sigma \times Br$ limit normalized to SM expectation (production processes assumed in SM proportions). No significant excess with respect to expected limits observed.

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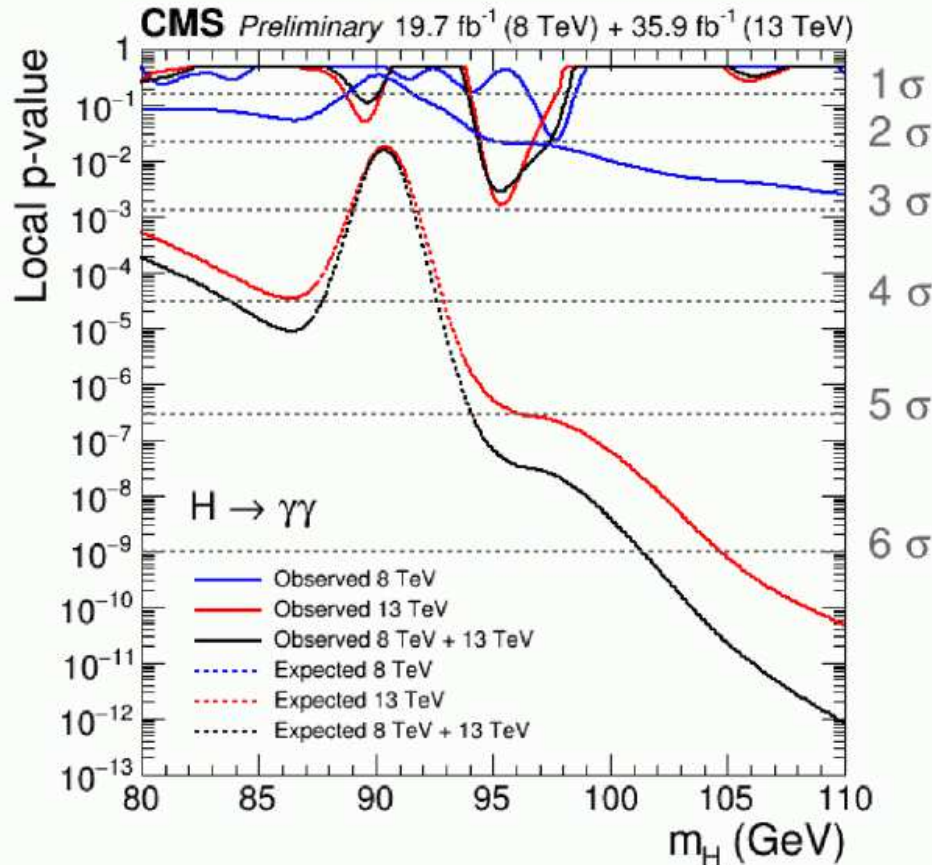
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$h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+ 2



CMS PAS HIG-17-013



8 TeV: Excess with $\sim 2.0 \sigma$ local significance at $m=97.6$ GeV

13 TeV: Excess with $\sim 2.9 \sigma$ local (1.47σ global) significance at $m=95.3$ GeV

8TeV+13 TeV: Excess with $\sim 2.8 \sigma$ local (1.3σ global) significance at $m=95.3$ GeV

More data are required to ascertain the origin of this excess

- Expected and observed local p-values for 8 TeV, 13 TeV and their combination

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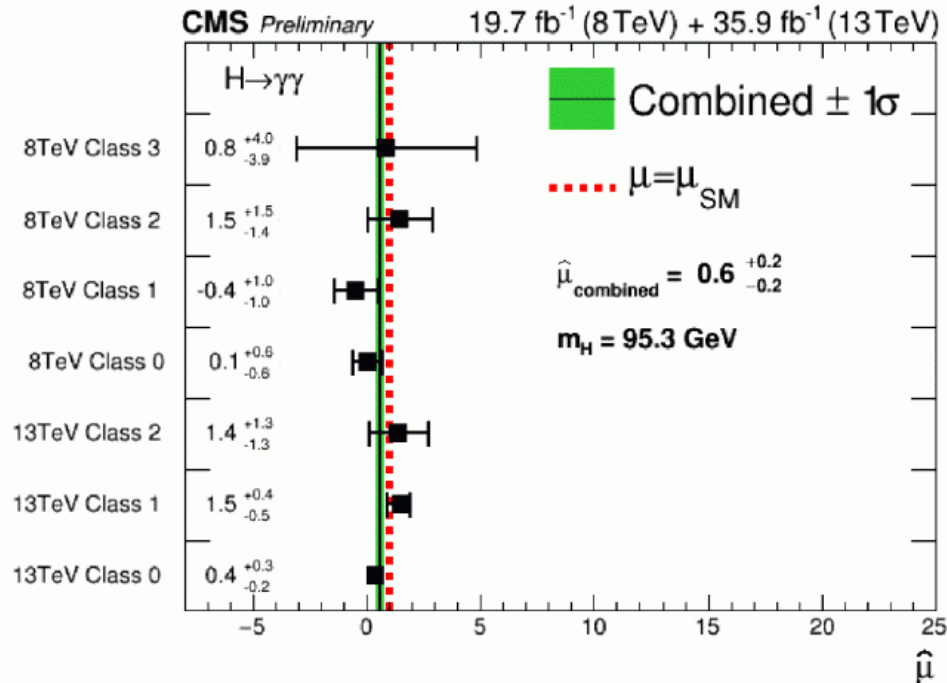
30



$h \rightarrow \gamma\gamma$ (70-110 GeV) Runs 1+2



CMS PAS HIG-17-013



Excess here mostly driven by class 1 (&2) at 13 TeV

χ^2 probability for the seven individual values to be compatible with a single signal hypothesis: 41%

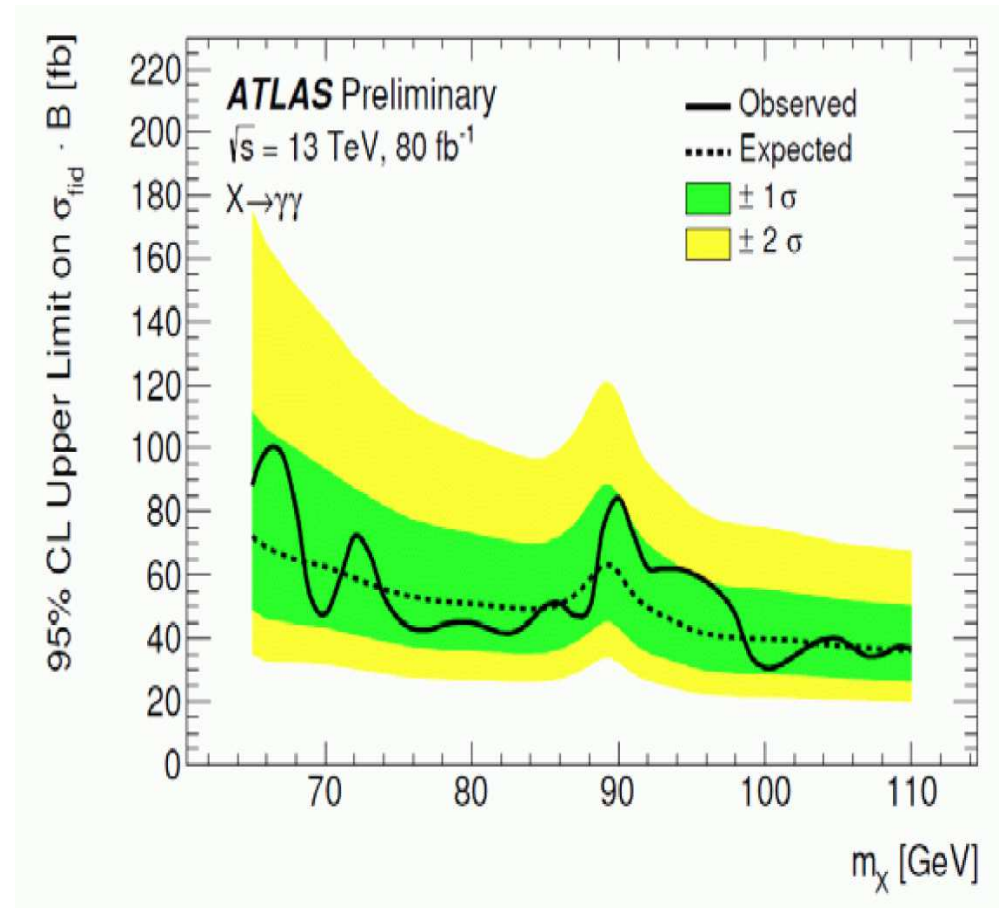
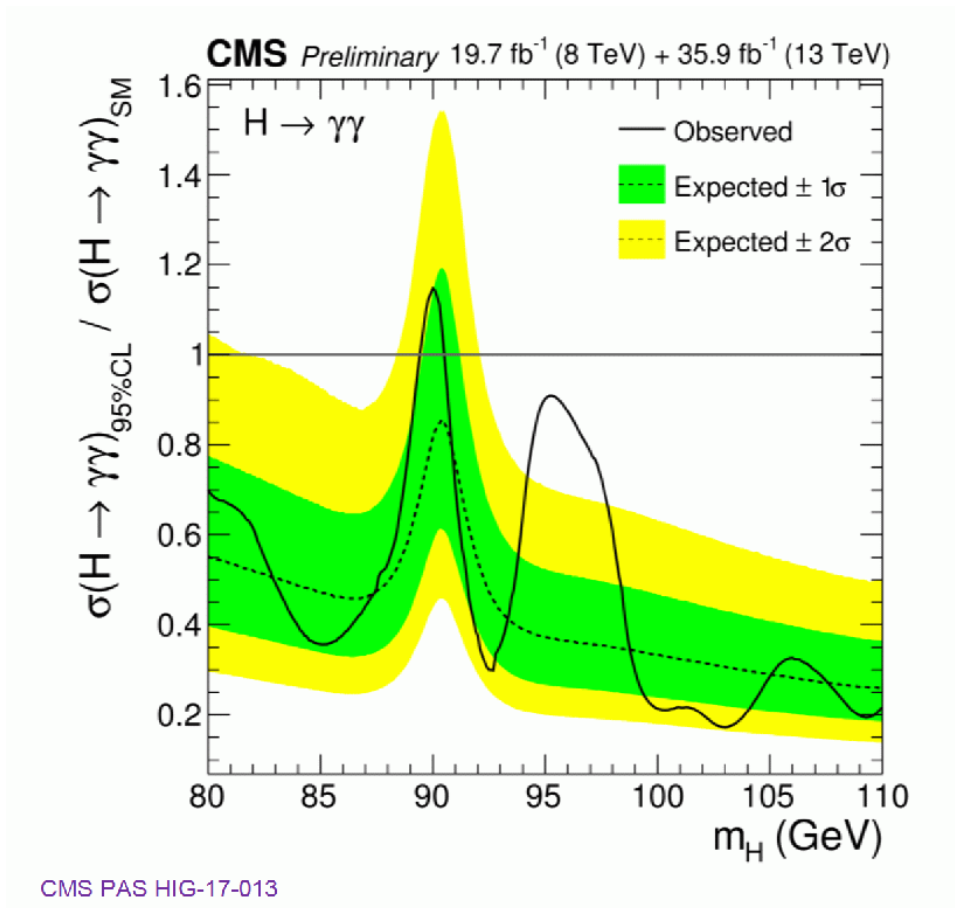
- ‘Signal’ strengths for the 7 event classes and overall, in the 8 TeV+13TeV combination, fixing $m_H=95.3 \text{ GeV}$
- More data are required to ascertain the origin of this excess

S. Gascon-Shotkin HDays17, Santander, ES Sept. 22 2017

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$$\mu_{CMS}(96 \text{ GeV}) = [\sigma(pp \rightarrow h_1) \times BR(h_1 \rightarrow \gamma\gamma)]_{exp/SM} = 0.6 \pm 0.2$$

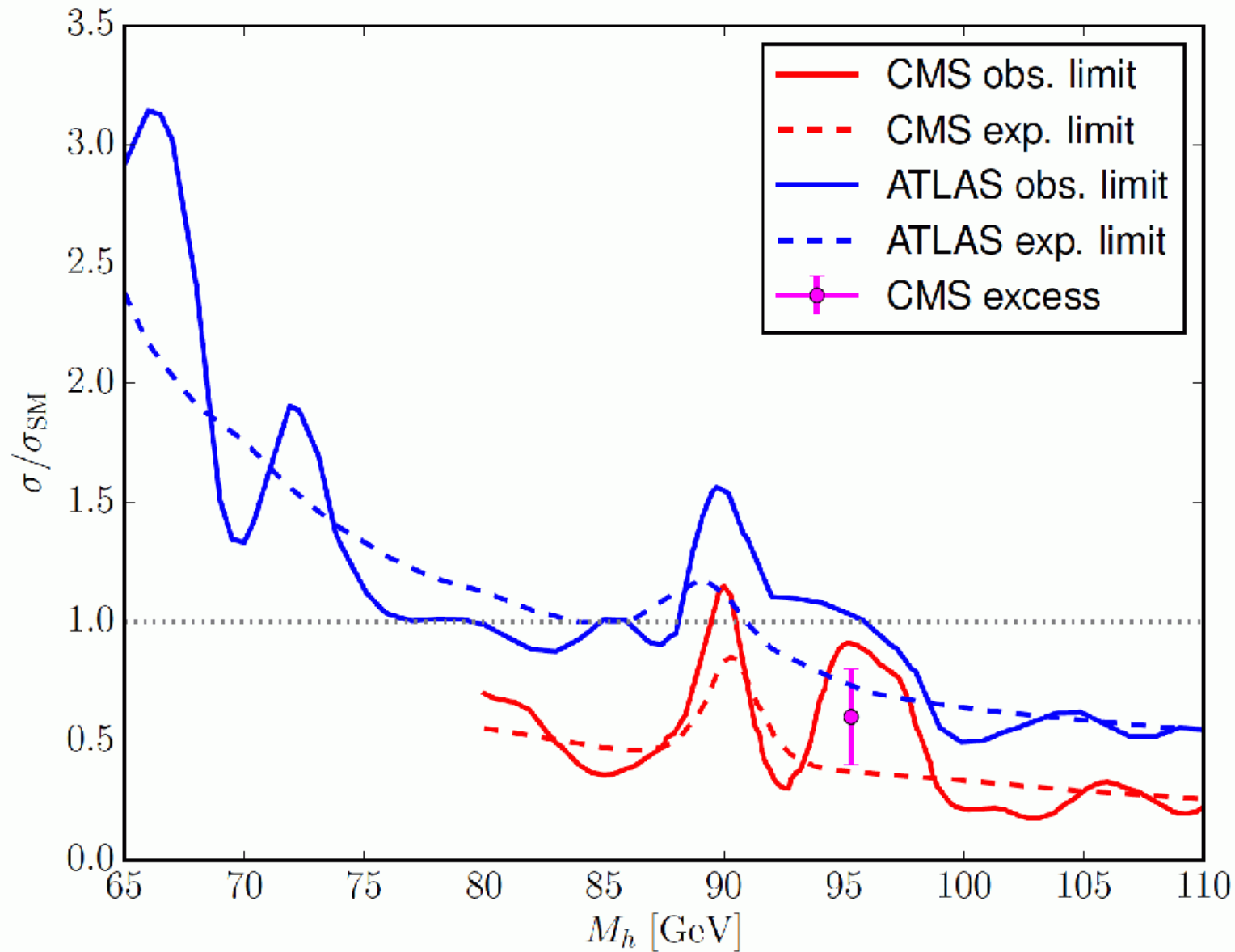
What about ATLAS?



Note: ATLAS gives fiducial cross section! Conversion factor: 1/0.45

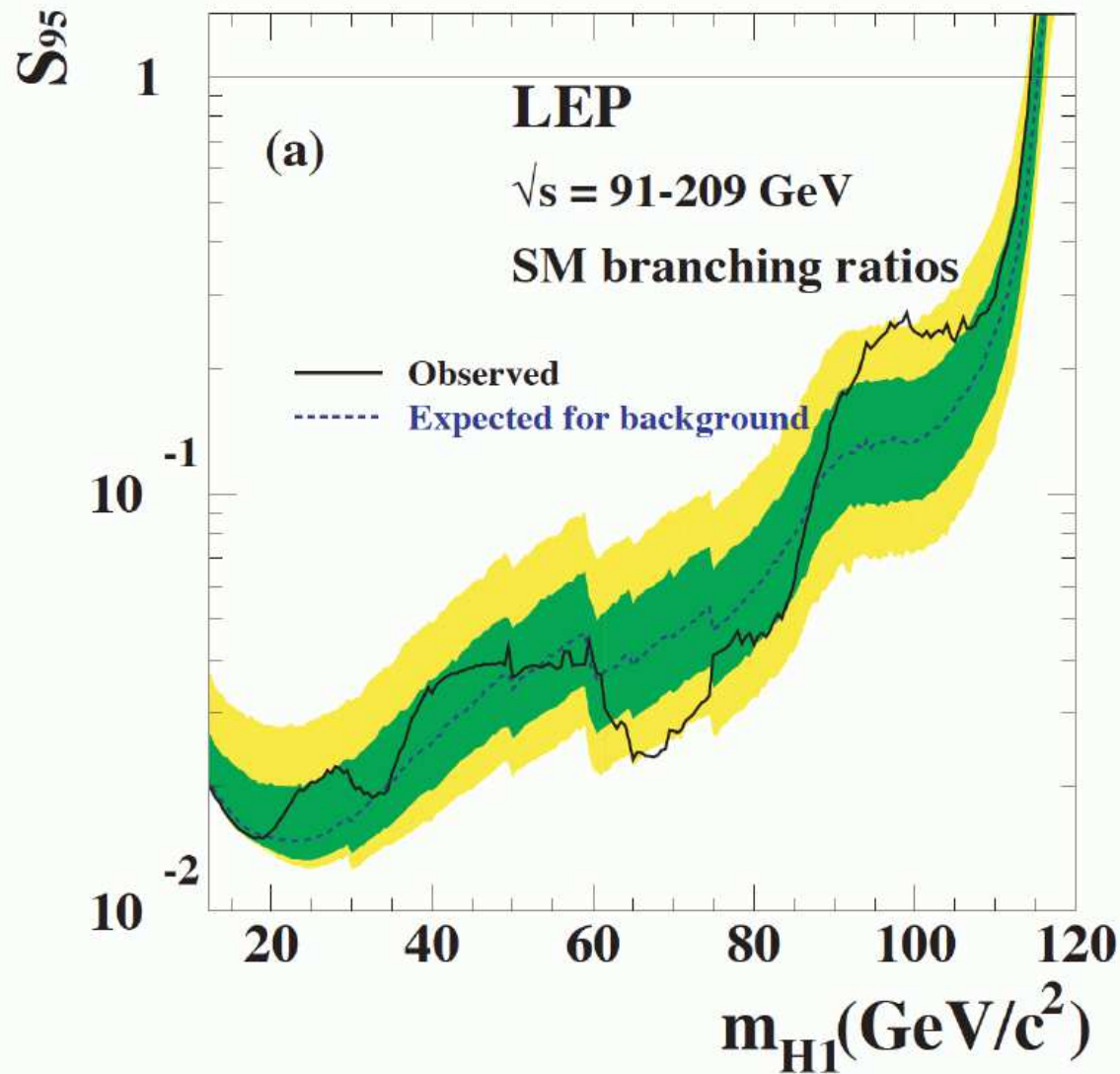
⇒ ATLAS exclusion limit even weaker than CMS!

Q: why does ATLAS has same sensitivity with twice amount of data?



⇒ if there is something, it would look exactly like this!

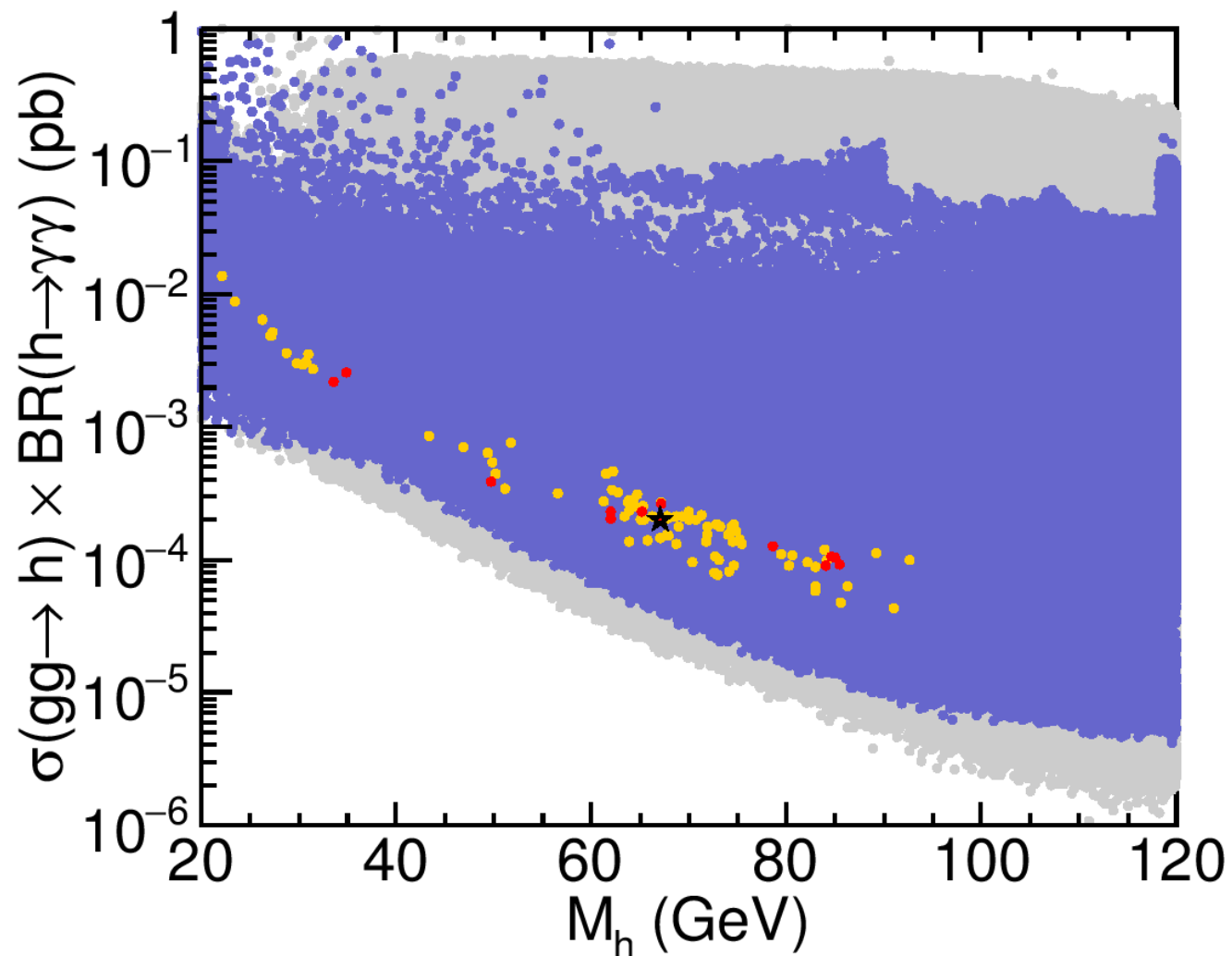
What was seen at LEP?



$$\mu_{\text{LEP}}(98 \text{ GeV}) = \left[\sigma(e^+e^- \rightarrow Zh_1) \times \text{BR}(h_1 \rightarrow b\bar{b}) \right]_{\text{exp/SM}} = 0.117 \pm 0.057$$

What about the MSSM?

[P. Bechtle, H. Haber, S.H., O. Stål, T. Stefaniak, G. Weiglein, L. Zeune '16]



⇒ too small rates!

⇒ problem: 2HDM structure too “rigid”

2. General analysis

MSSM: too small rates!

⇒ problem: 2HDM structure too “rigid”

More general Ansatz:

- richer Higgs structure
 - ⇒ add (at least) another Higgs singlet
- drop SUSY for now
 - ⇒ allow for more flexibility
 - ⇒ but check for hints towards SUSY
- check explicit SUSY scenarios later

More general Ansatz: **N2HDM**

[*T. Biekötter, M. Chakraborti, S.H. '19*]

Fields:

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + \rho_1 + i\eta_1) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + \rho_2 + i\eta_2) \end{pmatrix}, \quad \Phi_S = v_S + \rho_S$$

Potential:

$$\begin{aligned} V = & m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + h.c.) + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 \\ & + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \frac{\lambda_5}{2} [(\Phi_1^\dagger \Phi_2)^2 + h.c.] \\ & + \frac{1}{2} m_S^2 \Phi_S^2 + \frac{\lambda_6}{8} \Phi_S^4 + \frac{\lambda_7}{2} (\Phi_1^\dagger \Phi_1) \Phi_S^2 + \frac{\lambda_8}{2} (\Phi_2^\dagger \Phi_2) \Phi_S^2 \end{aligned}$$

Z_2 symmetry: $\Phi_1 \rightarrow \Phi_1$, $\Phi_2 \rightarrow -\Phi_2$, $\Phi_S \rightarrow \Phi_S$

Physical states: h_1, h_2, h_3 (CP -even), A (CP -odd), H^\pm (charged)

Extension of the Z_2 symmetry to fermions determines four types:

	u -type	d -type	leptons
type I	Φ_2	Φ_2	Φ_2
type II	Φ_2	Φ_1	Φ_1
type III (lepton-specific)	Φ_2	Φ_2	Φ_1
type IV (flipped)	Φ_2	Φ_1	Φ_2

\Rightarrow exactly as in 2HDM

Three neutral \mathcal{CP} -even Higgses:

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \rho_1 \\ \rho_2 \\ \rho_S \end{pmatrix}, \quad R = \begin{pmatrix} c_{\alpha_1} c_{\alpha_2} & s_{\alpha_1} c_{\alpha_2} & s_{\alpha_2} \\ -(c_{\alpha_1} s_{\alpha_2} s_{\alpha_3} + s_{\alpha_1} c_{\alpha_3}) & c_{\alpha_1} c_{\alpha_3} - s_{\alpha_1} s_{\alpha_2} s_{\alpha_3} & c_{\alpha_2} s_{\alpha_3} \\ -c_{\alpha_1} s_{\alpha_2} c_{\alpha_3} + s_{\alpha_1} s_{\alpha_3} & -(c_{\alpha_1} s_{\alpha_3} + s_{\alpha_1} s_{\alpha_2} c_{\alpha_3}) & c_{\alpha_2} c_{\alpha_3} \end{pmatrix}$$

Coupling to massive gauge bosons: (identical for all four types)

$$c_{h_i VV} = c_\beta R_{i1} + s_\beta R_{i2}$$

$$h_1 \quad c_{\alpha_2} c_{\beta - \alpha_1}$$

$$h_2 \quad -c_{\beta - \alpha_1} s_{\alpha_2} s_{\alpha_3} + c_{\alpha_3} s_{\beta - \alpha_1}$$

$$h_3 \quad -c_{\alpha_3} c_{\beta - \alpha_1} s_{\alpha_2} - s_{\alpha_3} s_{\beta - \alpha_1}$$

Coupling to fermions: (same pattern as in 2HDM)

	u -type ($c_{h_i tt}$)	d -type ($c_{h_i bb}$)	leptons ($c_{h_i \tau\tau}$)
type I	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$
type II	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i1}}{c_\beta}$
type III (lepton-specific)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$
type IV (flipped)	$\frac{R_{i2}}{s_\beta}$	$\frac{R_{i1}}{c_\beta}$	$\frac{R_{i2}}{s_\beta}$

“Physical” input parameters:

$$\alpha_{1,2,3}, \quad \tan \beta, \quad v, \quad v_S, \quad m_{h_{1,2,3}}, \quad m_A, \quad M_{H^\pm}, \quad m_{12}^2$$

Needed to fit the two excesses: $m_{h_1} \sim 96 \text{ GeV}$, $m_{h_2} \sim 125 \text{ GeV}$

- $c_{h_1 VV}^2$ strongly reduced for μ_{LEP}
- $c_{h_1 bb}$ reduced to enhance $\text{BR}(h_1 \rightarrow \gamma\gamma)$
- $c_{h_1 tt}$ not reduced for μ_{CMS}
- $c_{h_1 \tau\tau}$ possibly reduced to enhance $\text{BR}(h_1 \rightarrow \gamma\gamma)$

	Decrease $c_{h_1 b\bar{b}}$	No decrease $c_{h_1 t\bar{t}}$	No enhancement $c_{h_1 \tau\bar{\tau}}$
type I	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$
type II	$(\frac{R_{11}}{c_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{11}}{c_\beta}) :-)$
type III	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{11}}{c_\beta}) :-)$
type IV	$(\frac{R_{11}}{c_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$	$(\frac{R_{12}}{s_\beta}) :-)$

Type II and IV: $c_{h_1 bb}$ and $c_{h_1 tt}$ independent

Type II bonus: $c_{h_1 \tau\tau}$ can be suppressed (together with $c_{h_1 bb}$)

\Rightarrow only type II and IV can fit CMS and LEP excesses

⇒ Parameter scan ⇒ ScannerS

Constraints:

- Tree-level perturbativity ⇒ ScannerS
- Minimum of potential is global minimum ⇒ ScannerS
- Higgs searches at LEP, Tevatron, LHC ⇒ HiggsBounds
- SM-like Higgs properties ⇒ HiggsSignals (N2HDECAY, SusHi)
 $\chi_{\text{red}}^2 := \chi^2/n_{\text{obs}}$
- Flavor physics (mainly $\text{BR}(B_s \rightarrow X_s \gamma)$, ΔM_{B_s}) ⇒ SuperIso bounds
- Electroweak precision data (T and S) ⇒ ScannerS

Fitting the excesses:

$$\mu_{\text{LEP}} = 0.117 \pm 0.057, \quad \mu_{\text{CMS}} = 0.6 \pm 0.2$$

$$\begin{aligned}\mu_{\text{LEP}} &= \frac{\sigma_{\text{N2HDM}}(e^+e^- \rightarrow Zh_1)}{\sigma_{\text{SM}}(e^+e^- \rightarrow ZH)} \cdot \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow b\bar{b})}{\text{BR}_{\text{SM}}(H \rightarrow b\bar{b})} \\ &= |c_{h_1VV}|^2 \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow b\bar{b})}{\text{BR}_{\text{SM}}(H \rightarrow b\bar{b})}\end{aligned}$$

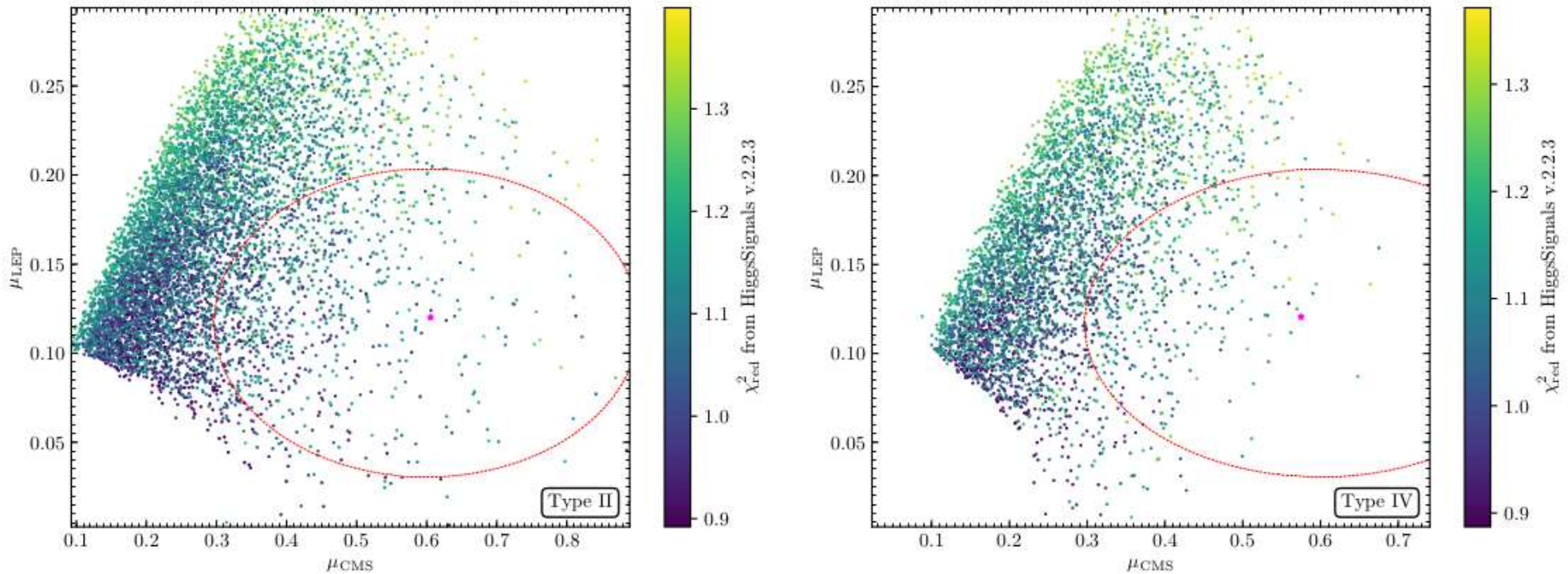
$$\begin{aligned}\mu_{\text{CMS}} &= \frac{\sigma_{\text{N2HDM}}(gg \rightarrow h_1)}{\sigma_{\text{SM}}(gg \rightarrow H)} \cdot \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow \gamma\gamma)}{\text{BR}_{\text{SM}}(H \rightarrow \gamma\gamma)} \\ &= |c_{h_1tt}|^2 \frac{\text{BR}_{\text{N2HDM}}(h_1 \rightarrow \gamma\gamma)}{\text{BR}_{\text{SM}}(H \rightarrow \gamma\gamma)}\end{aligned}$$

$$\chi_{\text{CMS-LEP}}^2 = \frac{(\mu_{\text{LEP}} - 0.117)^2}{(0.057)^2} + \frac{(\mu_{\text{CMS}} - 0.6)^2}{(0.2)^2}$$

⇒ “best-fit point”

Fitting the excesses:

[T. Biekötter, M. Chakraborti, S.H. '19]



⇒ excesses well fitted, with good χ_{red}^2 : 0.9 – 1.3

⇒ preferred M_{H^\pm} : 650 GeV – 950 GeV (lower limit: flavor constr.)

⇒ preferred $\tan \beta$: 0.8 – 3.8

Best-fit point in type II:

m_{h_1}	m_{h_2}	m_{h_3}	m_A	M_{H^\pm}	
96.5263	125.09	535.86	712.578	737.829	
$\tan \beta$	α_1	α_2	α_3	m_{12}^2	v_S
1.26287	1.26878	-1.08484	-1.24108	80644.3	272.72
$\text{BR}_{h_1}^{bb}$	$\text{BR}_{h_1}^{gg}$	$\text{BR}_{h_1}^{\tau\tau}$	$\text{BR}_{h_1}^{\gamma\gamma}$	$\text{BR}_{h_1}^{WW}$	$\text{BR}_{h_1}^{ZZ}$
0.5048	0.2682	$5.09 \cdot 10^{-2}$	$2.582 \cdot 10^{-3}$	$1.37 \cdot 10^{-2}$	$1.753 \cdot 10^{-3}$
$\text{BR}_{h_2}^{bb}$	$\text{BR}_{h_2}^{gg}$	$\text{BR}_{h_2}^{\tau\tau}$	$\text{BR}_{h_2}^{\gamma\gamma}$	$\text{BR}_{h_2}^{WW}$	$\text{BR}_{h_2}^{ZZ}$
0.5916	0.0771	$6.36 \cdot 10^{-2}$	$2.153 \cdot 10^{-3}$	0.2087	$2.610 \cdot 10^{-3}$

⇒ surprisingly large $\text{BR}_{h_1}^{\gamma\gamma}$

Best-fit point in type IV:

m_{h_1}	m_{h_2}	m_{h_3}	m_A	M_{H^\pm}	
97.8128	125.09	485.998	651.502	651.26	
$\tan \beta$	α_1	α_2	α_3	m_{12}^2	v_S
1.3147	1.27039	-1.02829	-1.32496	41034.1	647.886
$\text{BR}_{h_1}^{bb}$	$\text{BR}_{h_1}^{gg}$	$\text{BR}_{h_1}^{\tau\tau}$	$\text{BR}_{h_1}^{\gamma\gamma}$	$\text{BR}_{h_1}^{WW}$	$\text{BR}_{h_1}^{ZZ}$
0.4074	0.20714	0.248324	$2.139 \cdot 10^{-3}$	$1.347 \cdot 10^{-2}$	$1.579 \cdot 10^{-3}$
$\text{BR}_{h_2}^{bb}$	$\text{BR}_{h_2}^{gg}$	$\text{BR}_{h_2}^{\tau\tau}$	$\text{BR}_{h_2}^{\gamma\gamma}$	$\text{BR}_{h_2}^{WW}$	$\text{BR}_{h_2}^{ZZ}$
0.5363	0.09388	$7.58 \cdot 10^{-2}$	$2.247 \cdot 10^{-3}$	0.2267	$2.836 \cdot 10^{-2}$

⇒ substantially larger $\text{BR}_{h_1}^{\tau\tau}$ than in type II

What can we learn from future measurements?

- LHC h_{125} coupling measurements
- HL-LHC h_{125} coupling measurements
- ILC (or other e^+e^- collider) h_{125} coupling measurements

- direct production of ϕ_{96} at the LHC
- direct production of ϕ_{96} at the HL-LHC
- direct production of ϕ_{96} at the ILC (or other e^+e^- coll.)

- production of other BSM Higgs bosons at the LHC/HL-LHC/ILC/...

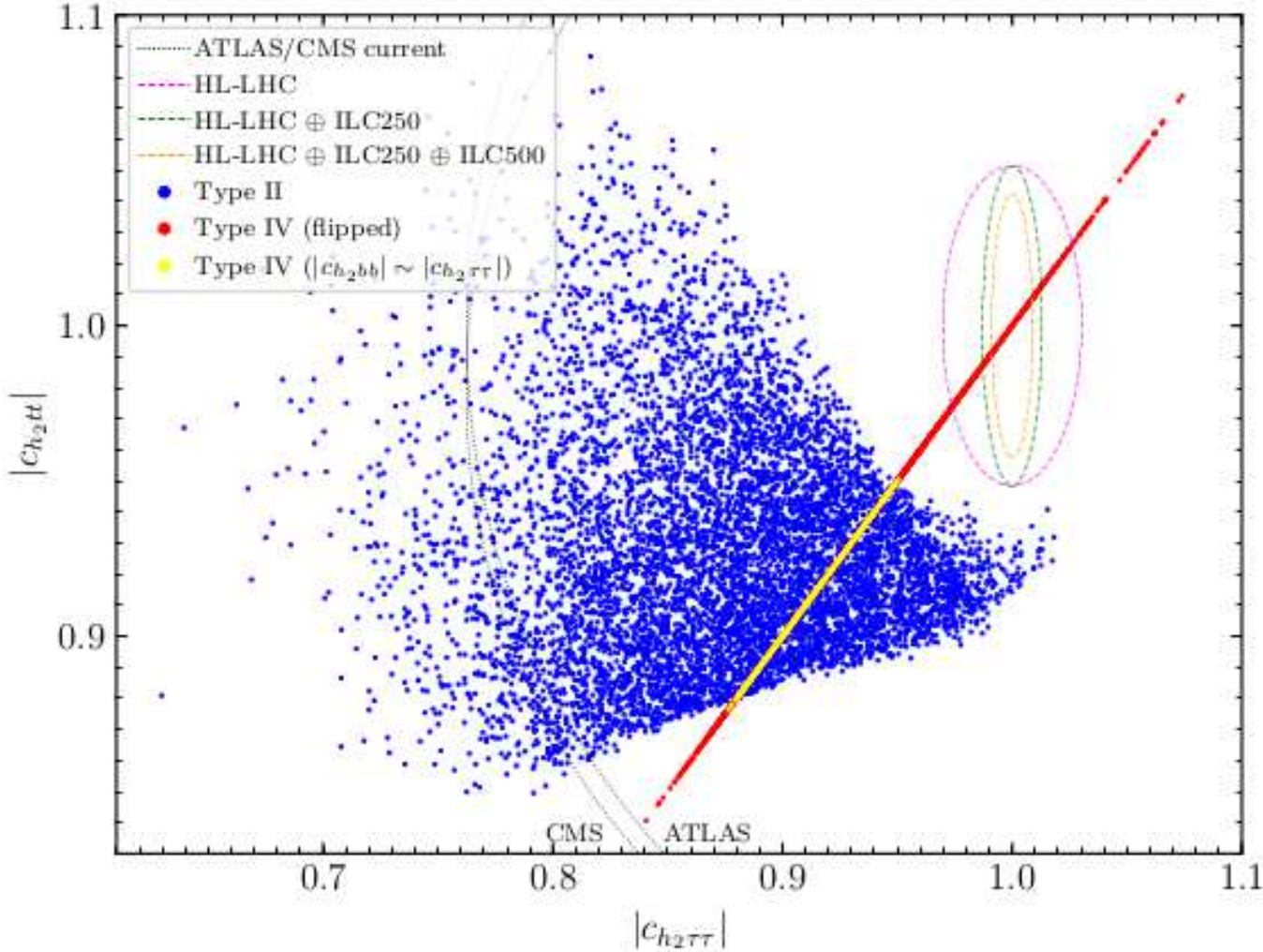
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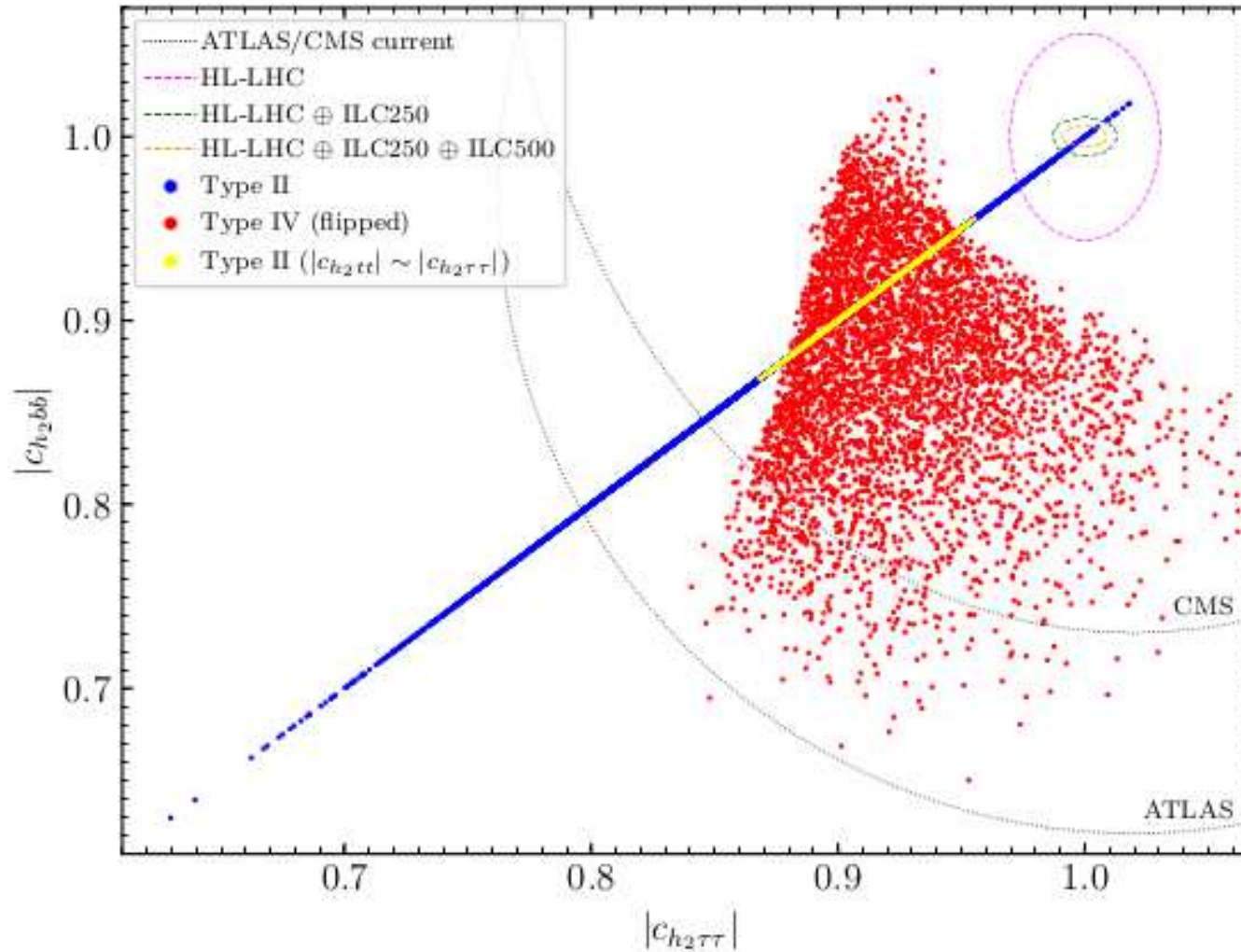
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Future measurements: \Rightarrow HL-LHC/ILC Higgs coupling measurements



\Rightarrow type II shows deviation from SM

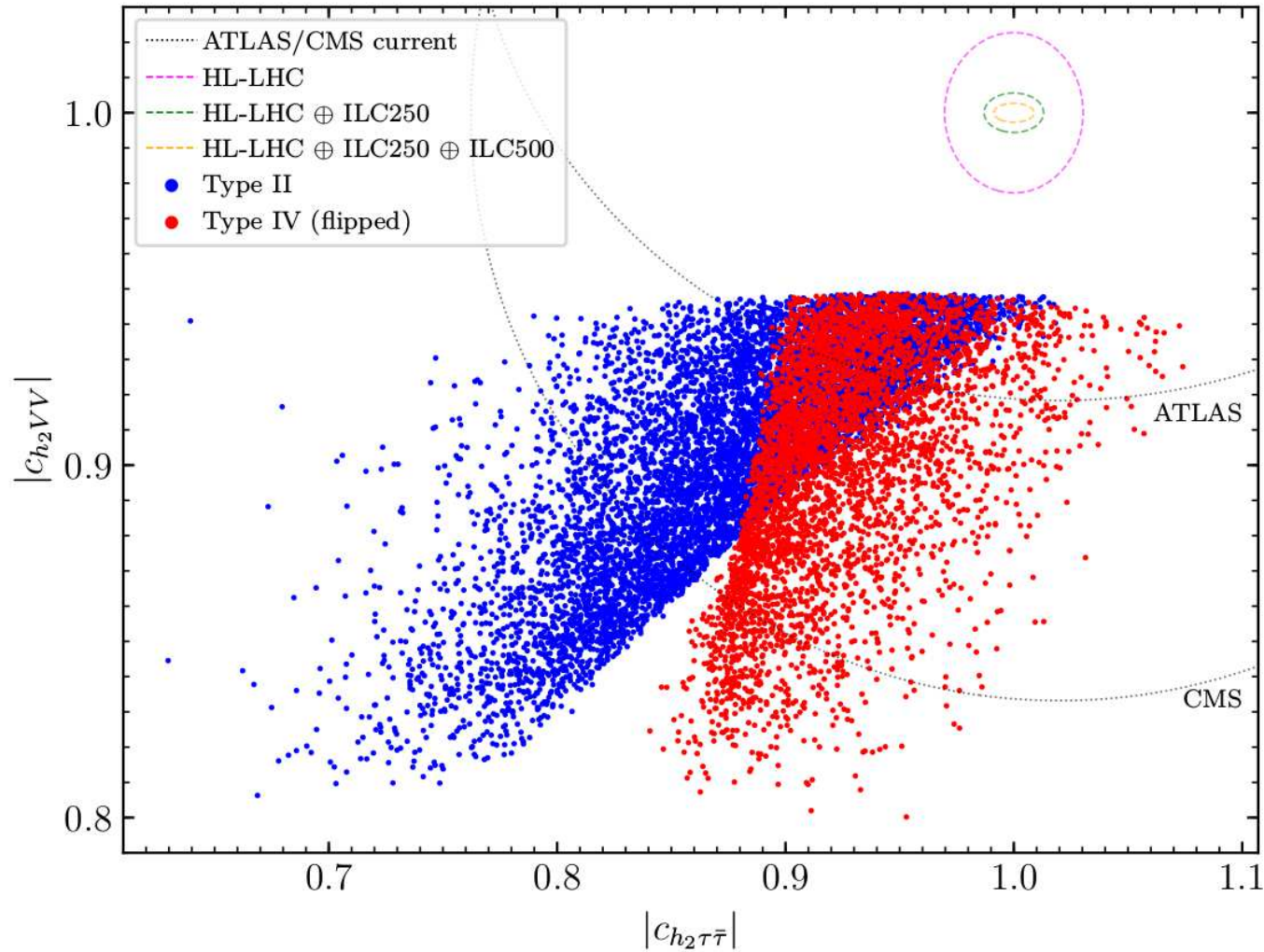
Future measurements: \Rightarrow HL-LHC/ILC Higgs coupling measurements



\Rightarrow type IV shows deviations from SM

\Rightarrow N2HDM can always be distinguished from SM!

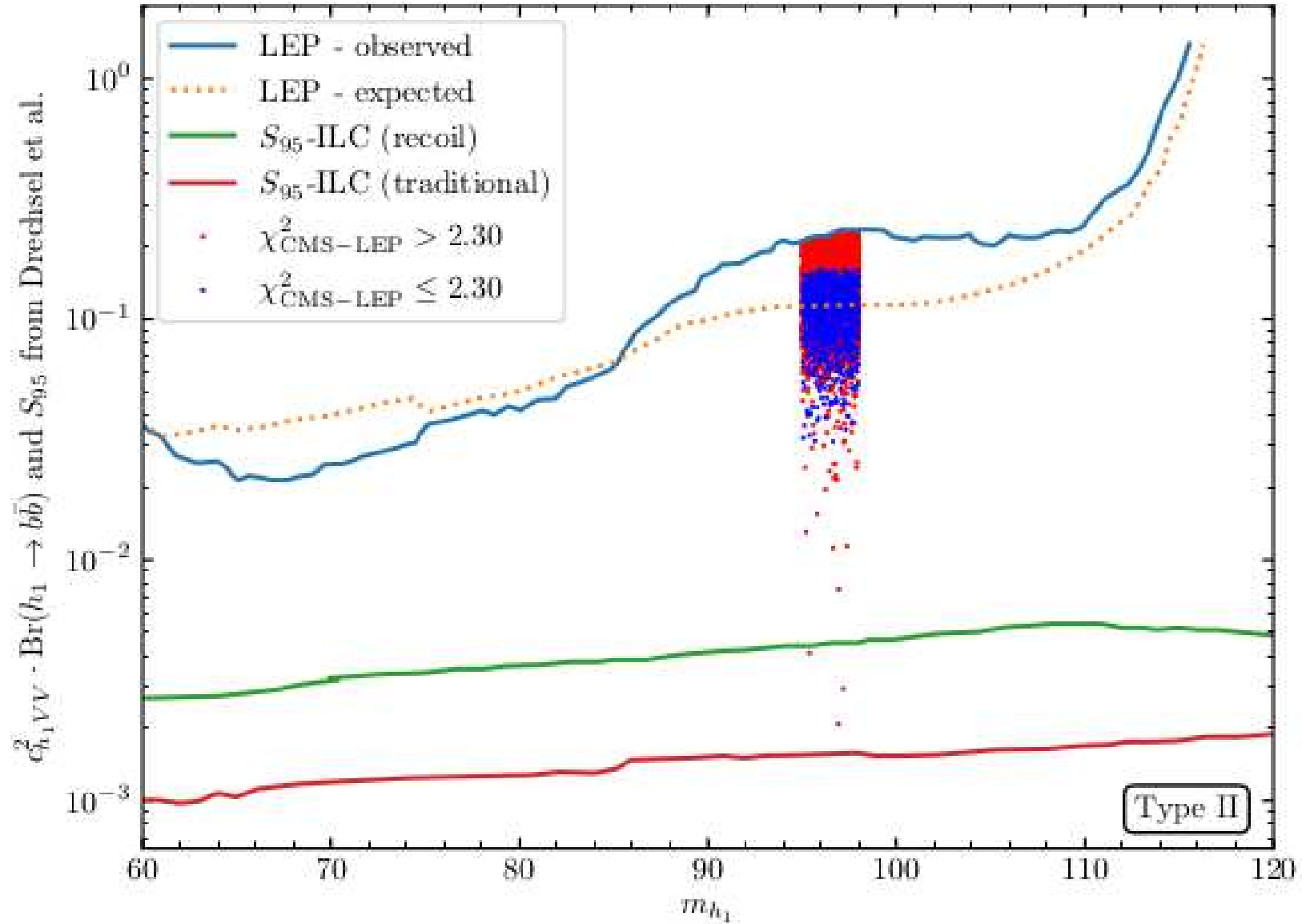
Future measurements: \Rightarrow HL-LHC/ILC Higgs coupling measurements



\Rightarrow type II and IV show strong deviations from SM

\Rightarrow N2HDM can always be distinguished from SM!

Next project? \Rightarrow ILC production of the light scalar



\Rightarrow new state easily in the reach of the ILC (or other e^+e^- collider)

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- $\mu\nu$ SJM
- ...

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⇒ models with an additional singlet??

- NMSSM
- $\mu\nu$ SSM
- ...

Q: Can the models fit the excesses **despite** the additional SUSY constraints on the Higgs sector **???**

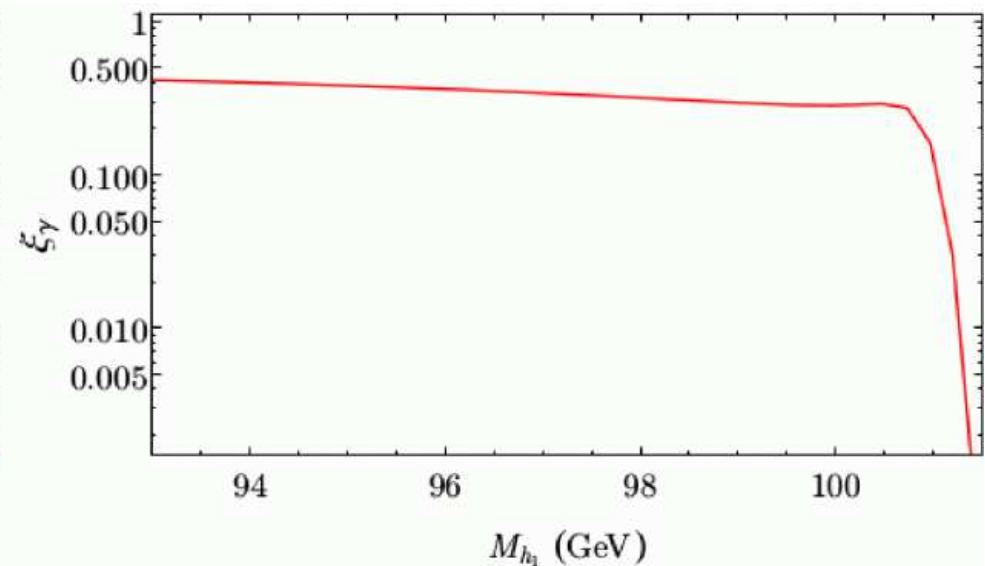
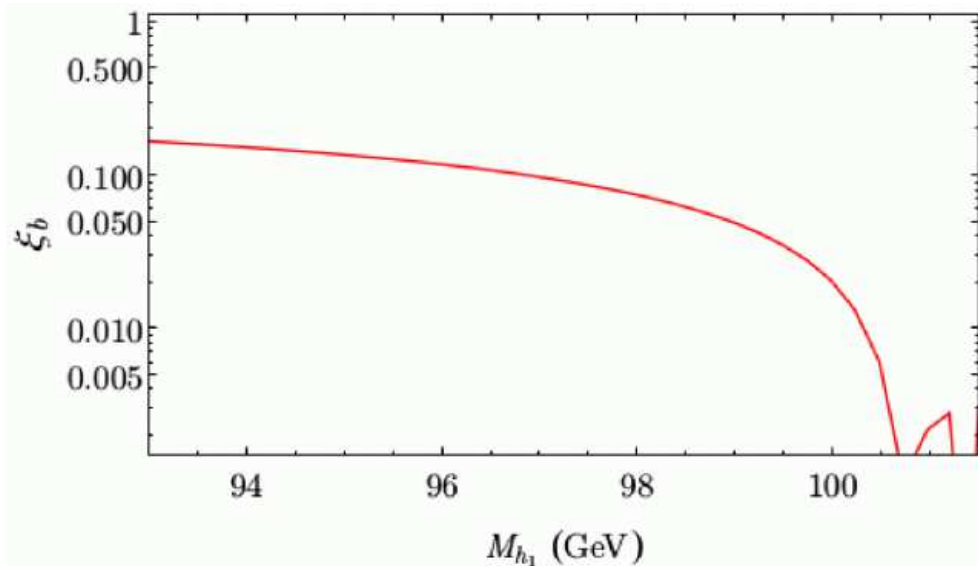
What about the NMSSM?

[F. Domingo, S.H., S. Passehr, G. Weiglein '18]

Parameters:

$\lambda = 0.6$, $\kappa = 0.035$, $\tan\beta = 2$, $\mu_{\text{eff}} = (397 + 15x)$ GeV, $M_{H^\pm} = 1$ TeV,
 $A_\kappa = -325$ GeV, $M_{\text{SUSY}} = 1$ TeV, $A_t = A_b = 0$

$$\xi_b \equiv \frac{\Gamma[h_1 \rightarrow ZZ] \cdot \text{BR}[h_1 \rightarrow b\bar{b}]}{\Gamma[H_{\text{SM}}(M_{h_1}) \rightarrow ZZ] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \rightarrow b\bar{b}]} \sim \frac{\sigma[e^+e^- \rightarrow Z(h_1 \rightarrow b\bar{b})]}{\sigma[e^+e^- \rightarrow Z(H_{\text{SM}}(M_{h_1}) \rightarrow b\bar{b})]}$$
$$\xi_\gamma \equiv \frac{\Gamma[h_1 \rightarrow gg] \cdot \text{BR}[h_1 \rightarrow \gamma\gamma]}{\Gamma[H_{\text{SM}}(M_{h_1}) \rightarrow gg] \cdot \text{BR}[H_{\text{SM}}(M_{h_1}) \rightarrow \gamma\gamma]} \sim \frac{\sigma[gg \rightarrow h_1 \rightarrow \gamma\gamma]}{\sigma[gg \rightarrow H_{\text{SM}}(M_{h_1}) \rightarrow \gamma\gamma]}.$$



⇒ both excesses can be fitted simultaneously (at $1 - 1.5\sigma$)!

What about the $\mu\nu$ SSM?

$\mu\nu$ SSM: [*D. Lopez-Fogliani, C. Muñoz '06*]

$\mu\nu$ SSM: NMSSM + well motivated RPV (in simple terms)
 \Rightarrow EW scale seesaw to reproduce the neutrino data

What about the $\mu\nu$ SSM?

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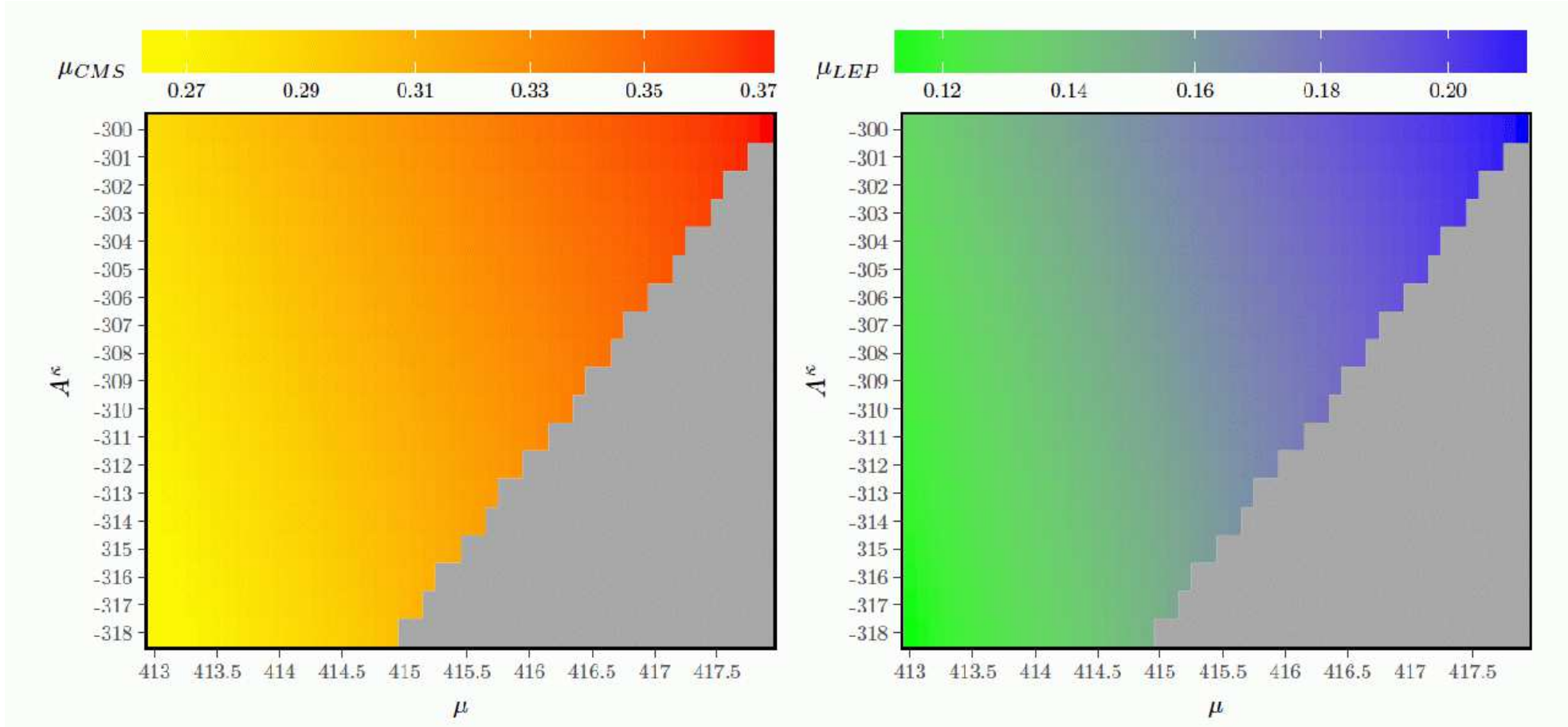
Can the $\mu\nu$ SSM explain the two excesses?

[T. Biekötter, S.H., C. Muñoz '17]

v_{iL}	Y_i^ν	A_i^ν	$\tan\beta$	μ	λ	A^λ	κ	A^κ	M_1
$\sqrt{2} \cdot 10^{-5}$	10^{-7}	-1000	2	[413; 418]	0.6	956.035	0.035	[-300; -318]	100
M_2	M_3	$m_{\tilde{Q}_{iL}}^2$	$m_{\tilde{u}_{iR}}^2$	$m_{\tilde{d}_{iR}}^2$	A_1^u	$A_{2,3}^{u,d}$	$(m_e^2)_{ii}$	A_{33}^e	$A_{11,22}^e$
200	1500	800^2	800^2	800^2	0	0	800^2	0	0

Can the $\mu\nu$ SSM explain the two excesses?

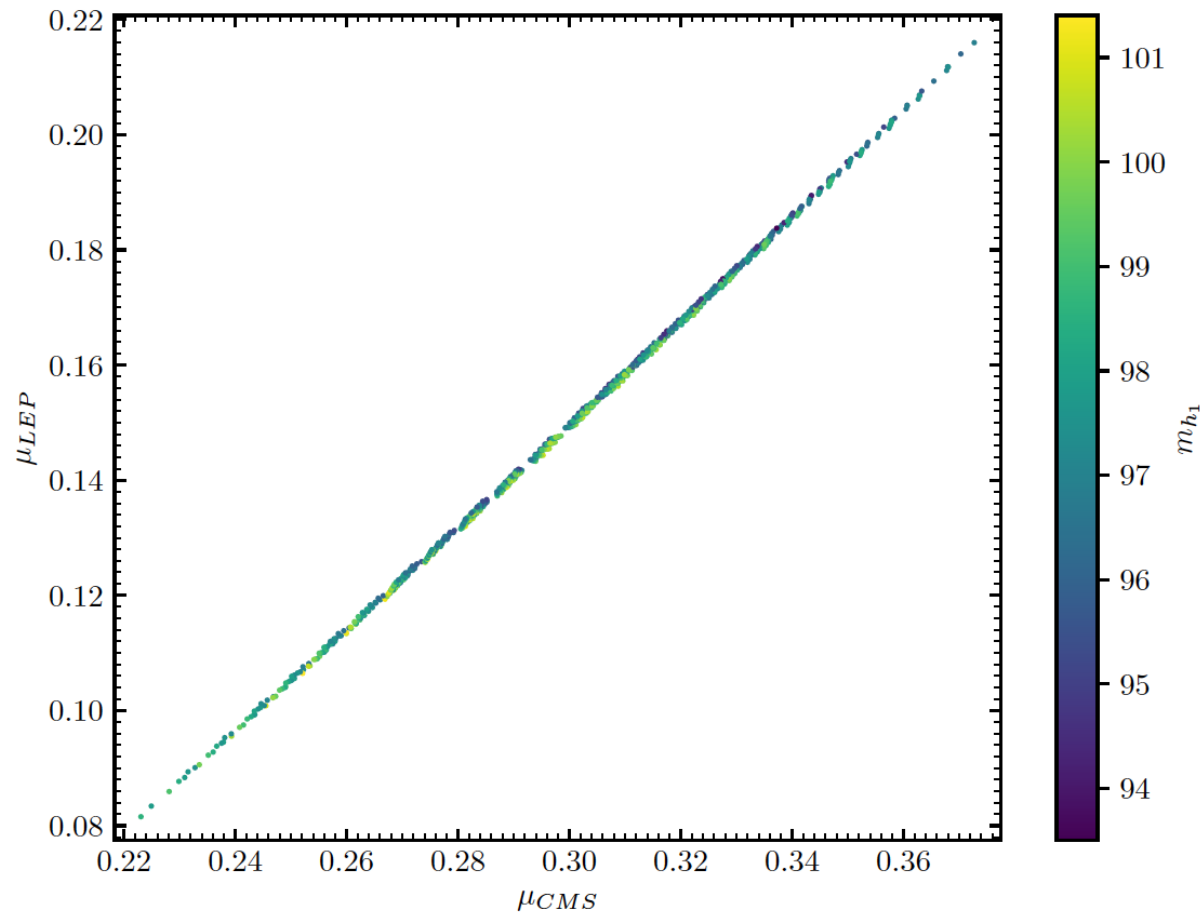
[*T. Biekötter, S.H., C. Muñoz '17*]



⇒ YES, WE CAN! :-)
at the 1 – 1.5 σ level

Why can SUSY explain the excesses only at $1 - 1.5 \sigma$?

[*T. Biekötter, S.H., C. Muñoz '19*]



⇒ SUSY enforces strong correlation!

⇒ note: ATLAS limits and CMS “observation”
will likely result in a lower μ_{LHC} !

4. Conclusinos

- Interesting excess in $\phi \rightarrow \gamma\gamma$ at 96 GeV at CMS
Compatible with $\phi \rightarrow b\bar{b}$ at 98 GeV at LEP
Compatible with ATLAS limits
- MSSM cannot explain the $\gamma\gamma$ excess \Rightarrow 2HDM Higgs structure too rigid
- More general Ansatz: N2DHM
 - type II and IV can fit the excesses
 - type II fits best (as predicted by SUSY)
- Measurements of h_{125} couplings at e^+e^- colliders
 - can firmly establish a difference w.r.t. the SM
 - can distinguish between type II and IV
- Direct production of ϕ_{96} at e^+e^- colliders
 - production easily possible
 - detailed study of ϕ_{96} possible
- SUSY realizations:
 - NMSSM can explain both excesses at $1 - 1.5 \sigma$
 - $\mu\nu$ SSM can explain both excesses at $1 - 1.5 \sigma$
 - SUSY enforces strong correlation between μ_{CMS} and μ_{LEP}
 - better agreement with μ_{LHC}

Higgs Days at Santander 2020

Theory meets Experiment

28 September - 02 October

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Further Questions?

