On the future of Higgs, electroweak and diboson measurements at lepton colliders

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Based on *J.B., G. Durieux, C. Grojean, J. Gu and A. Paul,* JHEP12 (2019) 117 (arXiv:1907.04311 [hep-ph])

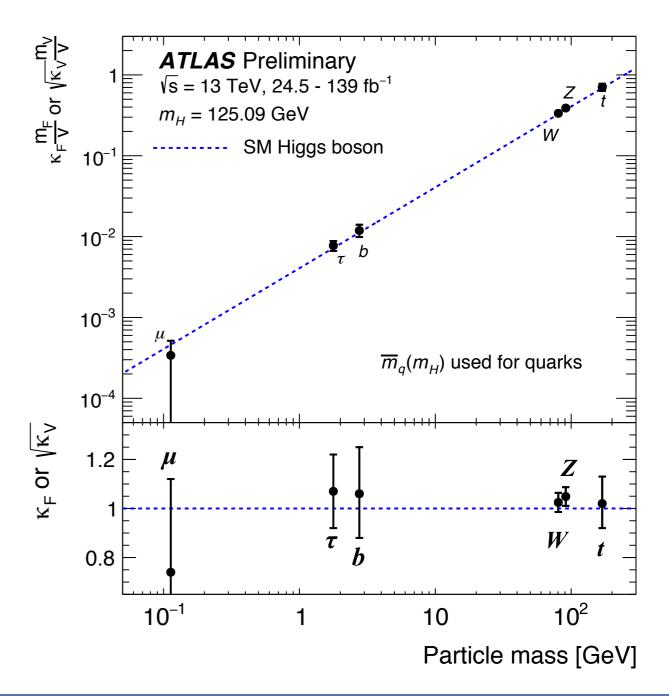


 The discovery of a Higgs-like boson in 2012 provided experimental confirmation of the last particle predicted by the SM...

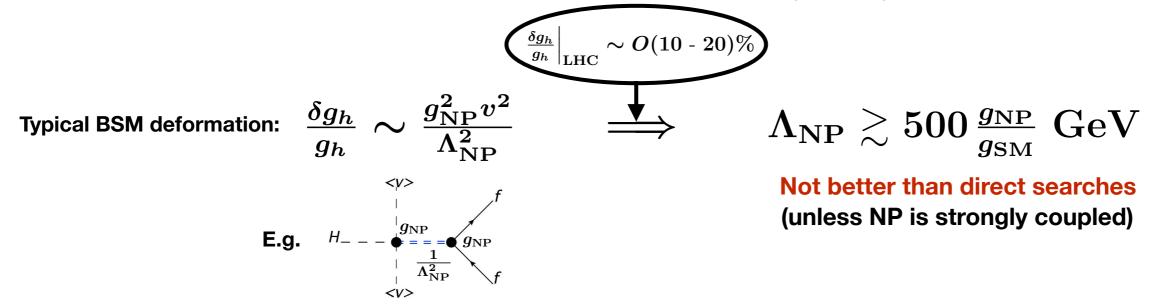
 So far, its properties also seem to be very SM-like... to ~10-20% accuracy.

 This precision, however, is still not enough to learn too much about possible sources of New Physics (NP)...

H Couplings



 Indeed, the measurements of the Higgs properties have not provided much information (constraints) on New Physics yet:



Higgs couplings also provide information about Naturalness:

$$\delta m_H^2 = \cdots + \cdots + \cdots \sim 0$$
 R. Rattazzi's talk at ESU symposium, Granada
$$\frac{\delta g_h}{g_h} \sim \frac{m_h^2}{\Delta m_h^2} \equiv \epsilon_T \equiv \text{fine tuning}$$

⇒ Higgs precision physics is a key tool to learn from BSM

 Therefore, a key component of the physics program at future particle colliders has revolved around the possible improvements on the knowledge of properties of the Higgs ...

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 One should however not forget about improving our measurements of the properties of the other particles of the EW sector:

Electroweak Precision Observables (EWPO)

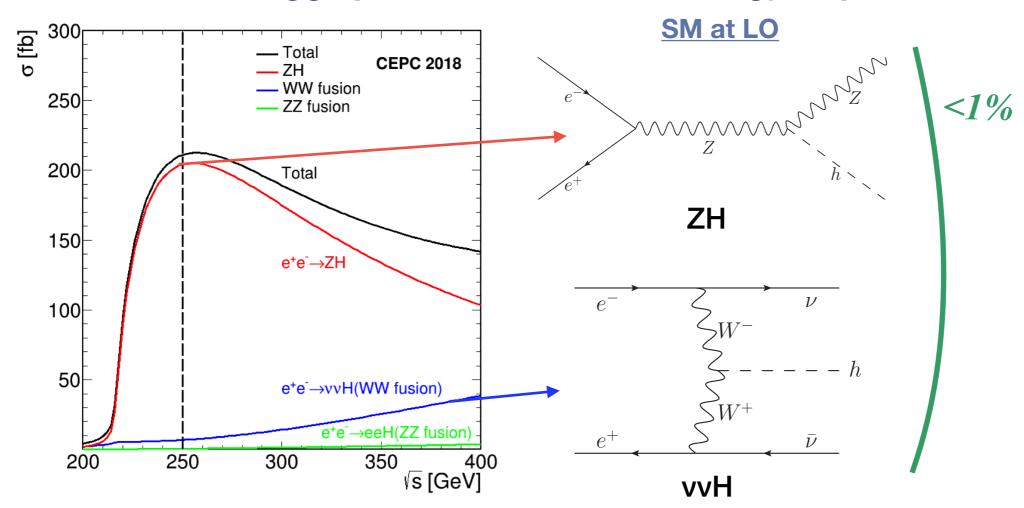
Measurements of cross sections and asymmetries, coming mostly from e⁺e⁻ colliders (LEP/SLD)

Provide precise measurements of properties (mass & couplings) of the W & Z bosons

 Therefore, a key component of the physics program at future particle colliders has revolved around the possible improvements on the knowledge of properties of the Higgs ...

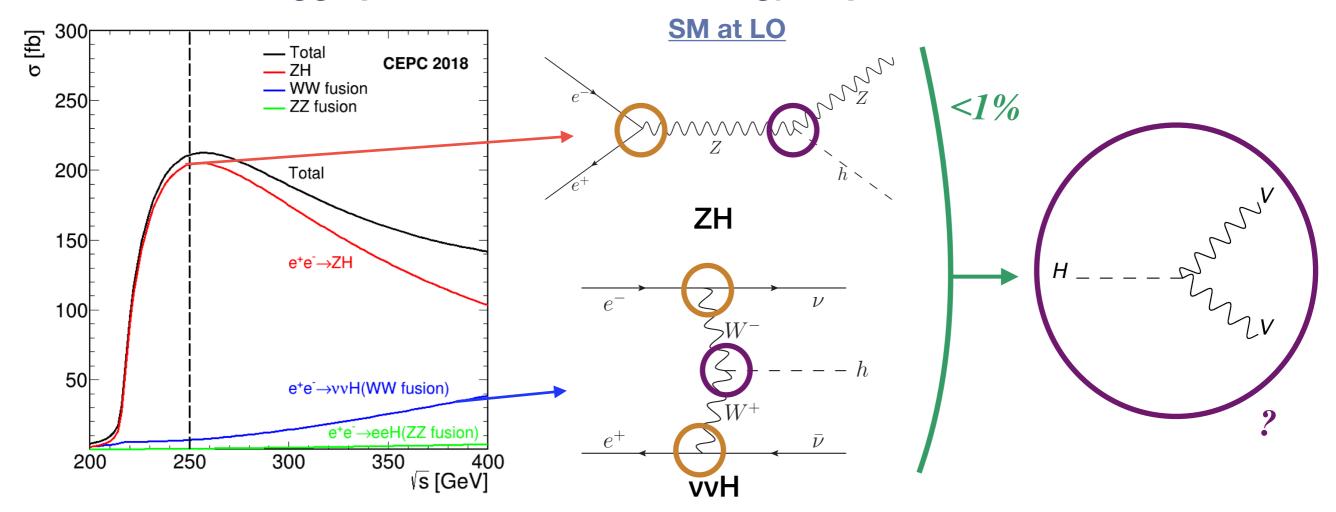
- One should however not forget about improving our measurements of the properties of the other particles of the EW sector:
 - ...even if the Higgs is the primary goal, one cannot separate the determination of its properties from the knowledge of the properties of the other SM particles...

Higgs production at "low-energy" lepton colliders



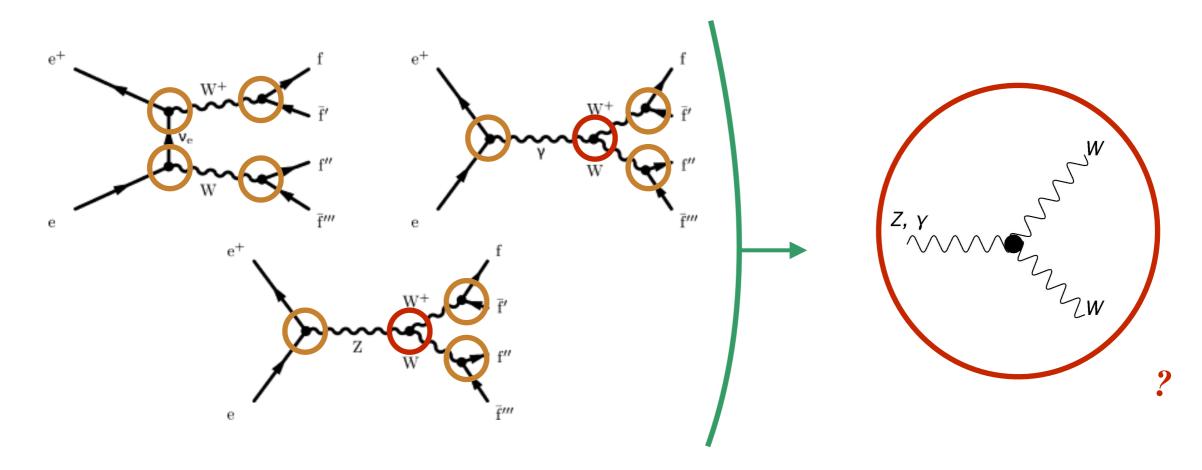
 Precision of Higgs measurements expected to be close to per mille level in several cases

Higgs production at "low-energy" lepton colliders



- Precision of Higgs measurements expected to be close to per mille level in several cases
- Is the knowledge of the EW interactions from LEP/SLD enough to neglect EW uncertainties in the extraction of Higgs properties?
- To answer this, first we need to set the (B)SM interpretation "framework"

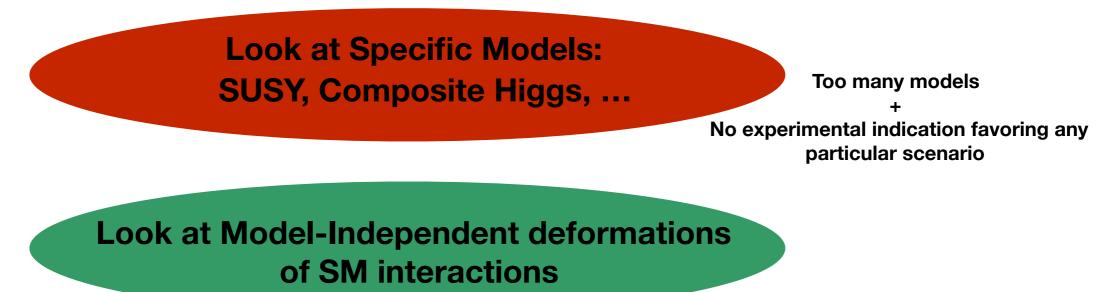
Similar considerations for WW production at lepton colliders



- At LEP2 aTGC determination was built on the assumption that the EW *Vff* are SM-like. Justified by the precise LEP/SLD Z-pole constraints
- Is the knowledge of the EW interactions from LEP/SLD enough to neglect EW uncertainties in the extraction of aTGC at future colliders?
- To answer this, first we need to set the (B)SM interpretation "framework"



• Two possible approaches to compare indirect sensitivity to BSM:



We will follow the model-independent approach using an EFT:



- The dimension 6 SMEFT:
- SM Particles and Symmetries at Low-energies. Assume H in SU(2) doublet
- Expansion in canonical mass dimension

$$\mathcal{L}_{ ext{Eff}} = \sum_{d=4}^{\infty} rac{1}{\Lambda^{d-4}} \mathcal{L}_d = \mathcal{L}_{ ext{SM}} + rac{1}{\Lambda} \mathcal{L}_5 + rac{1}{\Lambda^2} \mathcal{L}_6 + \cdots$$
 $\mathcal{L}_d = \sum_i C_i^d \mathcal{O}_i \qquad [\mathcal{O}_i] = d \longrightarrow \left(rac{q}{\Lambda}
ight)^{d-4}$ Effects suppressed by $q = v, E < \Lambda$

- LO new physics effects "start" at dimension 6: 59 B & L preserving operators B.Grzadkowski, M.Iskrynski, M.Misiak, J.Rosiek, JHEP 1010 (2010) 085 (2499 counting flavor)
- Describes correlations of BSM effects in different types of observables,

 We focus on EW/Higgs: Assume CP-even. 4-Fermion and dipole operators tested better by other processes (no EW/Higgs) and are neglected.

We also restrict the analysis to flavour preserving processes/interactions

Dimension-6 effects in Higgs/EW physics at LO

$$\begin{split} \Delta \mathscr{L}_{6}^{\text{hVV}} &= \frac{h}{v} \left[2\delta c_{w} m_{W}^{2} W_{\mu}^{+} W_{\mu}^{-} + \delta c_{z} m_{Z}^{2} Z_{\mu} Z_{\mu} + c_{w\Box} g^{2} \left(W_{\mu}^{-} \partial_{v} W_{\mu\nu}^{+} + \text{h.c.} \right) + c_{z\Box} g^{2} Z_{\mu} \partial_{v} Z_{\mu\nu} + c_{\gamma\Box} g g' Z_{\mu} \partial_{v} A_{\mu\nu} \right. \\ &+ c_{ww} \frac{g^{2}}{2} W_{\mu\nu}^{+} W_{\mu\nu}^{-} + c_{gg} \frac{g_{s}^{2}}{4} G_{\mu\nu}^{a} G_{\mu\nu}^{a} + c_{\gamma\gamma} \frac{e^{2}}{4} A_{\mu\nu} A_{\mu\nu} + c_{z\gamma} \frac{e\sqrt{g^{2} + g'^{2}}}{2} Z_{\mu\nu} A_{\mu\nu} + c_{zz} \frac{g^{2} + g'^{2}}{4} Z_{\mu\nu} Z_{\mu\nu} \right] \\ &\delta c_{w} = \delta c_{z} + 4\delta m, \\ c_{ww} = c_{zz} + 2\sin^{2} \theta_{w} c_{z\gamma} + \sin^{4} \theta_{w} c_{\gamma\gamma}, \\ c_{w\Box} = \frac{1}{g^{2} - g'^{2}} \left[g^{2} c_{z\Box} + g'^{2} c_{zz} - e^{2} \sin^{2} \theta_{w} c_{\gamma\gamma} - (g^{2} - g'^{2}) \sin^{2} \theta_{w} c_{z\gamma} \right] \\ c_{\gamma\Box} = \frac{1}{g^{2} - g'^{2}} \left[2g^{2} c_{z\Box} + (g^{2} + g'^{2}) c_{zz} - e^{2} c_{\gamma\gamma} - (g^{2} - g'^{2}) c_{z\gamma} \right], \end{split}$$

$$\Delta \mathcal{L}^{\text{aTGC}} = ie\delta \kappa_{\gamma} A^{\mu\nu} W_{\mu}^{+} W_{\nu}^{-} + ig\cos\theta_{w} \left[\delta g_{1Z} (W_{\mu\nu}^{+} W^{-\mu} - W_{\mu\nu}^{-} W^{+\mu}) Z^{\nu} + (\delta g_{1Z} - \frac{g^{\prime 2}}{g^{2}} \delta \kappa_{\gamma}) Z^{\mu\nu} W_{\mu}^{+} W_{\nu}^{-} \right] + \frac{ig\lambda_{z}}{m_{W}^{2}} \left(\sin\theta_{w} W_{\mu}^{+\nu} W_{\nu}^{-\rho} A_{\rho}^{\mu} + \cos\theta_{w} W_{\mu}^{+\nu} W_{\nu}^{-\rho} Z_{\rho}^{\mu} \right),$$

$$\Delta \mathscr{L}_{6}^{\mathrm{hff}} = -\frac{h}{v} \sum_{f \in u,d,e} \hat{\delta} y_f m_f \bar{f} f + \mathrm{h.c.}$$

$$\Delta \mathcal{L}_{6}^{vff,hvff} = \frac{g}{\sqrt{2}} \left(1 + 2\frac{h}{v} \right) W_{\mu}^{+} \left(\hat{\delta} g_{L}^{W\ell} \bar{\mathbf{v}} \bar{\gamma}_{\mu} e + \hat{\delta} g_{L}^{Wq} \bar{\mathbf{u}} \gamma_{\mu} d + \hat{\delta} g_{R}^{Wq} \bar{\mathbf{u}} \gamma_{\mu} d + \text{h.c.} \right)$$

$$+ \sqrt{g^{2} + g'^{2}} \left(1 + 2\frac{h}{v} \right) Z_{\mu} \left[\sum_{f=u,d,e,v} \hat{\delta} g_{L}^{Zf} \bar{f} \gamma_{\mu} f + \sum_{f=u,d,e} \hat{\delta} g_{R}^{Zf} \bar{f} \gamma_{\mu} f \right]$$

Dimension-6 effects in Higgs/EW physics at LO



$$\Delta \mathcal{L}_{6}^{\text{hVV}} = \frac{h}{v} \left[2\delta c_{w} m_{W}^{2} W_{\mu}^{+} W_{\mu}^{-} + \delta c_{z} n_{Z}^{2} Z_{\mu} Z_{\mu} + c_{w\square} g^{2} \left(W_{\mu}^{-} \partial_{\nu} W_{\mu\nu}^{+} + \text{h.c.} \right) + c_{z\square} g^{2} Z_{\mu} \partial_{\nu} Z_{\mu\nu} + c_{\gamma\square} g g' Z_{\mu} \partial_{\nu} A_{\mu\nu} \right]$$

$$+c_{ww}\frac{g^{2}}{2}W_{\mu\nu}^{+}W_{\mu\nu}^{-} + c_{gg}\frac{g_{s}^{2}}{4}G_{\mu\nu}^{a}G_{\mu\nu}^{a} + c_{\gamma\gamma}\frac{e^{2}}{4}A_{\mu\nu}A_{\mu\nu} + c_{z\gamma}\frac{e\sqrt{g^{2}+g^{\prime}\,^{2}}}{2}Z_{\mu\nu}A_{\mu\nu} + c_{zz}\frac{g^{2}+g^{\prime}\,^{2}}{4}Z_{\mu\nu}Z_{\mu\nu}$$

Only 7 independent parameters

$$\delta c_w = \delta c_z + 4\delta m$$

$$c_{ww} = c_{zz} + 2\sin^2\theta_w c_{z\gamma} + \sin^4\theta_w c_{\gamma\gamma},$$

$$c_{w\Box} = \frac{1}{\varrho^2 - \varrho'^2} \left[g^2 c_{z\Box} + g'^2 c_{zz} - e^2 \sin^2 \theta_w c_{\gamma\gamma} - (g^2 - g'^2) \sin^2 \theta_w c_{z\gamma} \right]$$

$$c_{\gamma\Box} = \frac{1}{g^2 - g'^2} \left[2g^2 c_{z\Box} + (g^2 + g'^2) c_{zz} - e^2 c_{\gamma\gamma} - (g^2 - g'^2) c_{z\gamma} \right],$$

"basis"



Where to test these?

$$H o VV' pp o HV \ gg o H ext{ } e^+e^- o HZ$$

(Tree level)

Dimension-6 effects in Higgs/EW physics at LO

 $\Delta \mathcal{L}_{6}^{\text{hVV}} = \frac{h}{v} \left[2\delta c_{w} m_{W}^{2} W_{\mu}^{+} W_{\mu}^{-} + \delta c_{z} m_{Z}^{2} Z_{\mu} Z_{\mu} + c_{w\square} g^{2} \left(W_{\mu}^{-} \partial_{\nu} W_{\mu\nu}^{+} + \text{h.c.} \right) + c_{z\square} g^{2} Z_{\mu} \partial_{\nu} Z_{\mu\nu} + c_{\gamma\square} g g' Z_{\mu} \partial_{\nu} A_{\mu\nu} \right. \\ \left. + c_{ww} \frac{g^{2}}{2} W_{\mu\nu}^{+} W_{\mu\nu}^{-} + c_{gg} \frac{g_{s}^{2}}{4} G_{\mu\nu}^{a} G_{\mu\nu}^{a} + c_{\gamma\gamma} \frac{e^{2}}{4} A_{\mu\nu} A_{\mu\nu} + c_{z\gamma} \frac{e\sqrt{g^{2} + g'^{2}}}{2} Z_{\mu\nu} A_{\mu\nu} + c_{zz} \frac{g^{2} + g'^{2}}{4} Z_{\mu\nu} Z_{\mu\nu} \right]$

$$\delta c_w = \delta c_z + 4\delta m,$$

$$c_{ww} = c_{zz} + 2\sin^2\theta_w c_{z\gamma} + \sin^4\theta_w c_{\gamma\gamma},$$

$$1 \qquad 1 \qquad 2 \qquad 2$$

CC

 $\Delta \mathcal{L}^{aTGC} = ie\delta \kappa_{\gamma} A^{\mu\nu} W_{\mu}^{+} W_{\nu}^{-} + ig\cos\theta_{w} \left[\delta g_{1Z} (W_{\mu\nu}^{+} W^{-\mu} - W_{\mu\nu}^{-} W^{+\mu}) Z^{\nu} + (\delta g_{1Z} - \frac{g^{\prime 2}}{g^{2}} \delta \kappa_{\gamma}) Z^{\mu\nu} W_{\mu}^{+} W_{\nu}^{-} \right]$

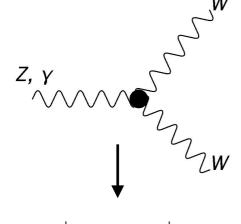
$$+\frac{i \varepsilon \lambda_z}{m_W^2} \left(\sin \theta_w W_\mu^{+\nu} W_\nu^{-\rho} A_\rho^\mu + \cos \theta_w W_\mu^{+\nu} W_\nu^{-\rho} Z_\rho^\mu \right),$$

Only λ_z is independent

 δg_{1Z} and $\delta \kappa_{V}$ related to HVV couplings



\mathbf{W} Where to test these?



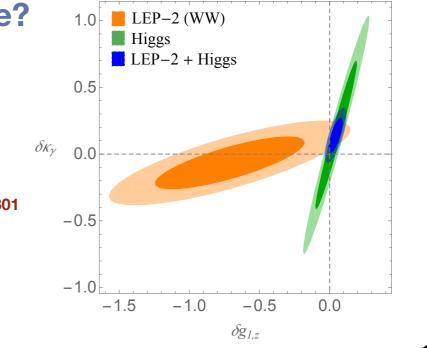
 $e^+e^- \to W^+W^-$

 $pp o W^+W^-, WZ, W\gamma$



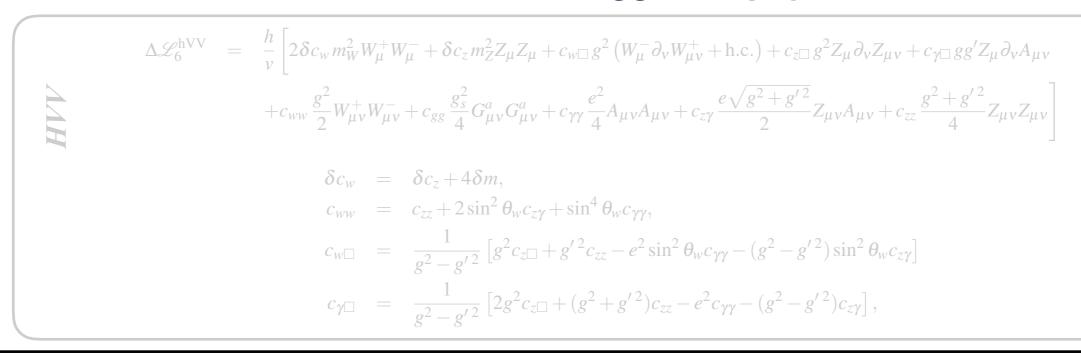
A. Falkowski et al., PRL 116 (2016) 011801

(Tree level)



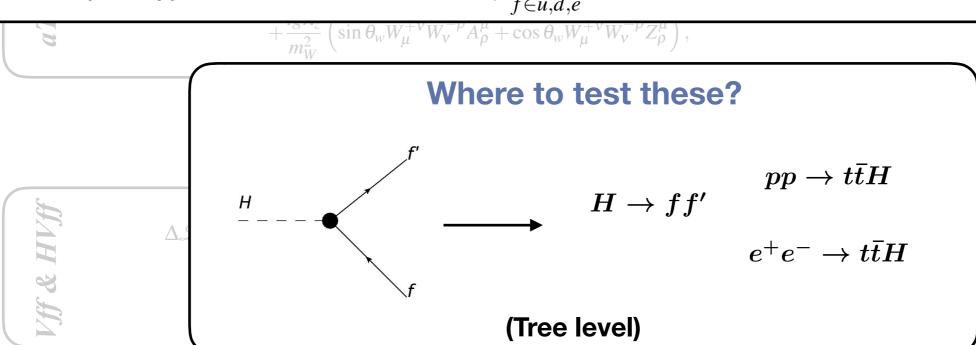
Higgs "basis" parameter

Dimension-6 effects in Higgs/EW physics at LO



5 NP parameters (t,b,c,τ,μ)

$$\Delta \mathcal{L}_{6}^{\text{hff}} = -\frac{h}{v} \sum_{f \in u,d,e} \hat{\delta} y_f m_f \bar{f} f + \text{h.c.}$$



Dimension-6 effects in Higgs/EW physics at LO

 $\Delta \mathcal{L}_{6}^{\text{hVV}} = \frac{h}{v} \left[2\delta c_{w} m_{W}^{2} W_{\mu}^{+} W_{\mu}^{-} + \delta c_{z} m_{Z}^{2} Z_{\mu} Z_{\mu} + c_{w\square} g^{2} \left(W_{\mu}^{-} \partial_{\nu} W_{\mu\nu}^{+} + \text{h.c.} \right) + c_{z\square} g^{2} Z_{\mu} \partial_{\nu} Z_{\mu\nu} + c_{\gamma\square} g g' Z_{\mu} \partial_{\nu} A_{\mu\nu} \right]$

$$\left. + c_{ww} \frac{g^2}{2} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{gg} \frac{g_s^2}{4} G_{\mu\nu}^a G_{\mu\nu}^a + c_{\gamma\gamma} \frac{e^2}{4} A_{\mu\nu} A_{\mu\nu} + c_{z\gamma} \frac{e\sqrt{g^2 + {g^\prime}^2}}{2} Z_{\mu\nu} A_{\mu\nu} + c_{zz} \frac{g^2 + {g^\prime}^2}{4} Z_{\mu\nu} Z_{\mu\nu} \right]$$

$$\delta c_w = \delta c_z + 4\delta m,$$

$$c_{ww} = c_{zz} + 2\sin^2\theta_w c_{z\gamma} + \sin^4\theta_w c_{\gamma\gamma},$$

$$c_{w\Box} = \frac{1}{g^2 - g'^2} \left[g^2 c_{z\Box} + g'^2 c_{zz} - e^2 \sin^2 \theta_w c_{\gamma\gamma} - (g^2 - g'^2) \sin^2 \theta_w c_{z\gamma} \right]$$

 $\Delta \mathcal{L}_{6}^{vff,hvff} = \frac{g}{\sqrt{2}} \left(1 + 2\frac{h}{v} \right) W_{\mu}^{+} \left(\hat{\delta} g_{L}^{W\ell} \bar{v} \bar{\gamma}_{\mu} e + \hat{\delta} g_{L}^{Wq} \bar{u} \gamma_{\mu} d + \hat{\delta} g_{R}^{Wq} \bar{u} \gamma_{\mu} d + \text{h.c.} \right)$

$$+ \sqrt{g^2 + g'^2} \left(1 + 2\frac{h}{v} \right) Z_{\mu} \left[\sum_{f=u,d,e,v} \hat{\delta} g_L^{Zf} \bar{f} \gamma_{\mu} f + \sum_{f=u,d,e} \hat{\delta} g_R^{Zf} \bar{f} \gamma_{\mu} f \right]$$

$$\hat{\delta}g_L^{W\ell} = \hat{\delta}g_L^{Z\nu} - \hat{\delta}g_L^{Ze}$$

$$\hat{\delta}g_L^{Wq} = \hat{\delta}g_L^{Zu}V_{CKM} - V_{CKM}\hat{\delta}g_L^{Zd}$$

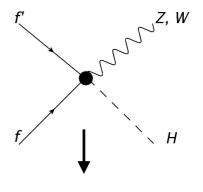
paramete

Higgs "basis"

Where to test these?

"Same" EFT interaction





$$e^+e^- o HZ$$

Dimension-6 effects in Higgs/EW physics at LO

 $\Delta \mathcal{L}_{6}^{\text{hVV}} = \frac{h}{v} \left[2\delta c_{w} m_{W}^{2} W_{\mu}^{+} W_{\mu}^{-} + \delta c_{z} m_{Z}^{2} Z_{\mu} Z_{\mu} + c_{w\square} g^{2} \left(W_{\mu}^{-} \partial_{\nu} W_{\mu\nu}^{+} + \text{h.c.} \right) + c_{z\square} g^{2} Z_{\mu} \partial_{\nu} Z_{\mu\nu} + c_{\gamma\square} g g' Z_{\mu} \partial_{\nu} A_{\mu\nu} \right]$

SMEFT fit

- -Hff and Vff (HVff) diagonal in the physical basis
- -Vff (HVff) flavour universality respected by first 2 quark families

-For H & EW exploration purposes only -Cumbersome from model-building point of view to avoid FCNC

Parameter counting in the parameterization of LHCHXSWG-INT-2015-001

$$egin{aligned} \{\delta m,\ c_{gg},\ \delta c_z,\ c_{\gamma\gamma},\ c_{z\gamma},\ c_{zz},\ c_{z\square},\ \delta y_t,\ \delta y_c,\ \delta y_b,\ \delta y_ au,\ \delta y_\mu,\ \lambda_z\} +\ &\left\{ (\delta g_L^{Zu})_{q_i}, (\delta g_L^{Zd})_{q_i}, (\delta g_L^{W\ell
u})_\ell, (\delta g_L^{Ze})_\ell, (\delta g_R^{Zu})_{q_i}, (\delta g_R^{Zd})_{q_i}, (\delta g_R^{Ze})_\ell
ight\}_{\substack{q_1=q_2,\ \ell=e,\mu, au\\ (q_i
eq t)}} \end{aligned}$$

15 Vff/HVff

5 SM + 28 New Physics Parameters

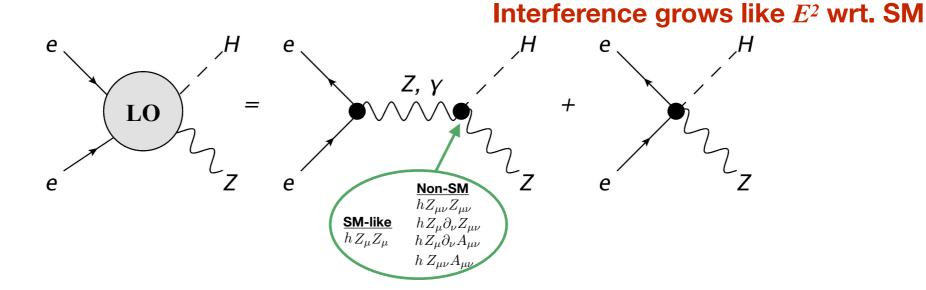
To study the impact of EW measurements in the determination of Higgs properties we will compare with an scenario with "perfect EW" data

i.e. EW interactions known to be SM-like with infinite precision:

$$\text{``Perfect EW"} \rightarrow \ \left\{\delta m, (\delta g_L^{Zu})_{q_i}, (\delta g_L^{Zd})_{q_i}, (\delta g_L^{W\ell\nu})_\ell, (\delta g_L^{Ze})_\ell, (\delta g_R^{Zu})_{q_i}, (\delta g_R^{Zd})_{q_i}, (\delta g_R^{Ze})_\ell\right\} \equiv 0$$

<u>Higgs production in the EFT framework</u>

 New type of contributions: apart from new HVV' tensor structures, virtual exchange of BSM particles can generate contact interactions



 Remember, these HZff terms are connected to modifications of Zff couplings, e.g.

$$\phi^{\dagger}i\stackrel{\leftrightarrow}{D}_{\!\mu}\phi \; \overline{e_R}\gamma^{\mu}e_R \sim rac{ev^2}{2sc}Z_{\mu}\overline{e_R}\gamma^{\mu}e_R + rac{ev}{sc}HZ_{\mu}\overline{e_R}\gamma^{\mu}e_R + \dots$$

Uncertainty on (H)Zee introduces growing-with-E "contamination" in the extraction of HZZ interactions from ZH processes

 \Rightarrow Use future EWPO (Z-pole data) to constrain $Zee \rightarrow HZee$

Presentation of SMEFT fit results

- Compare Future Collider sensitivity to BSM EW/Higgs deformations in a basis-independent way
 - <u>Project EFT fit</u> results from dim-6 Wilson coefficients into (pseudo)
 observable Higgs <u>effective couplings</u>, aTGC and effective <u>Zff/Wff</u> couplings
 defined from EWPO:

Dim 6 Wilson coeffs.

Effective couplings

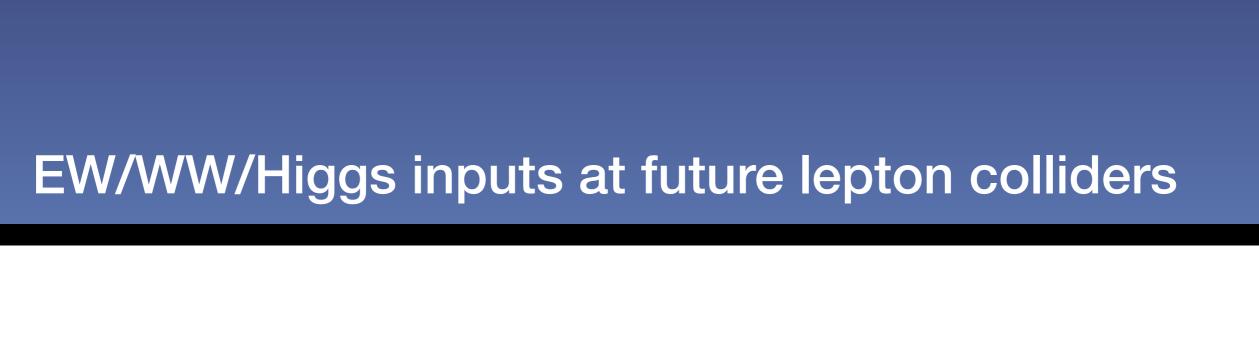
$$egin{array}{ll} rac{C_i}{\Lambda^2} & \longrightarrow & g_x^{ ext{eff}} \equiv g_x^{ ext{eff}} \left(rac{C_i}{\Lambda^2}
ight) = \left.g_x^{ ext{eff}}
ight|_{ ext{SM}} + \sum lpha_i^x rac{C_i}{\Lambda^2} + O(rac{C^2}{\Lambda^4})
ight.$$

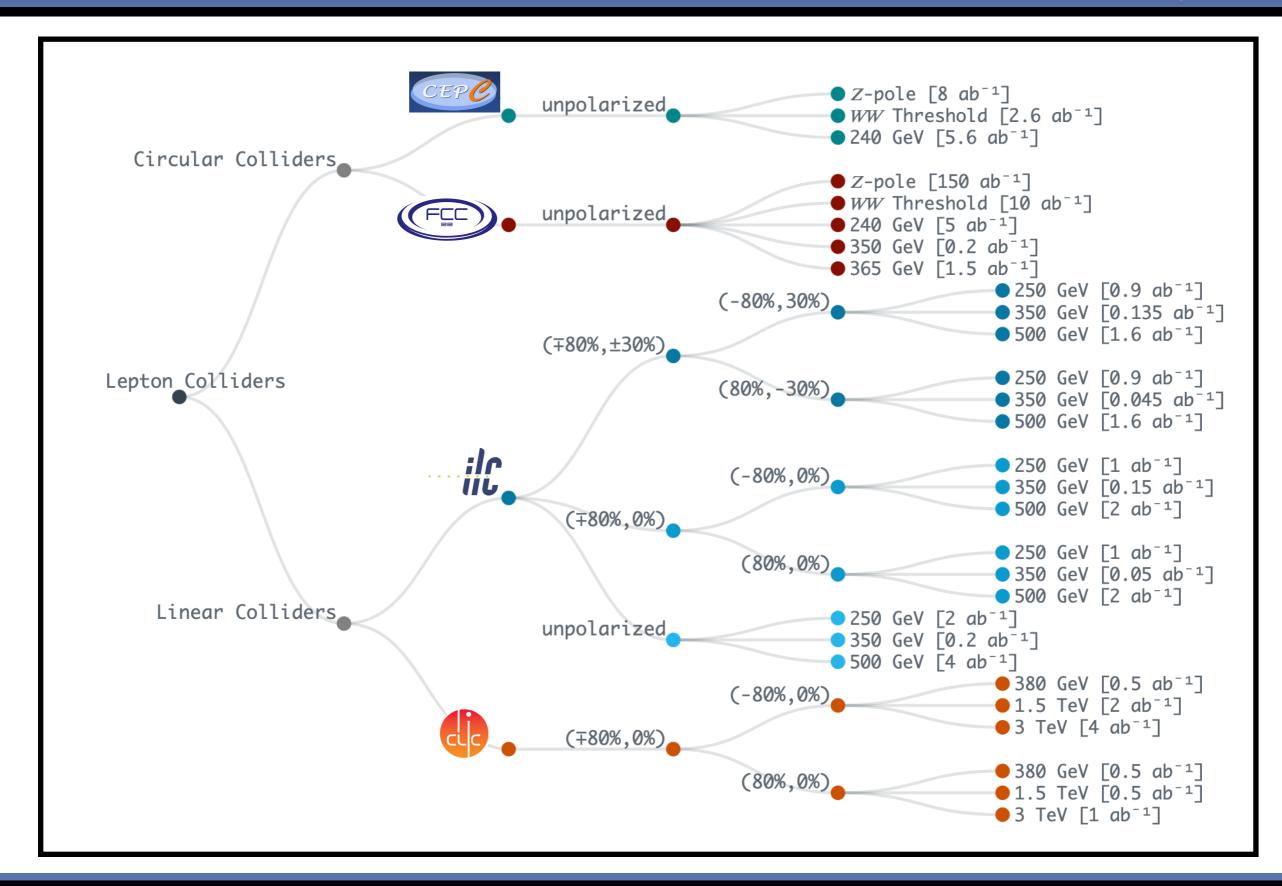
Higgs

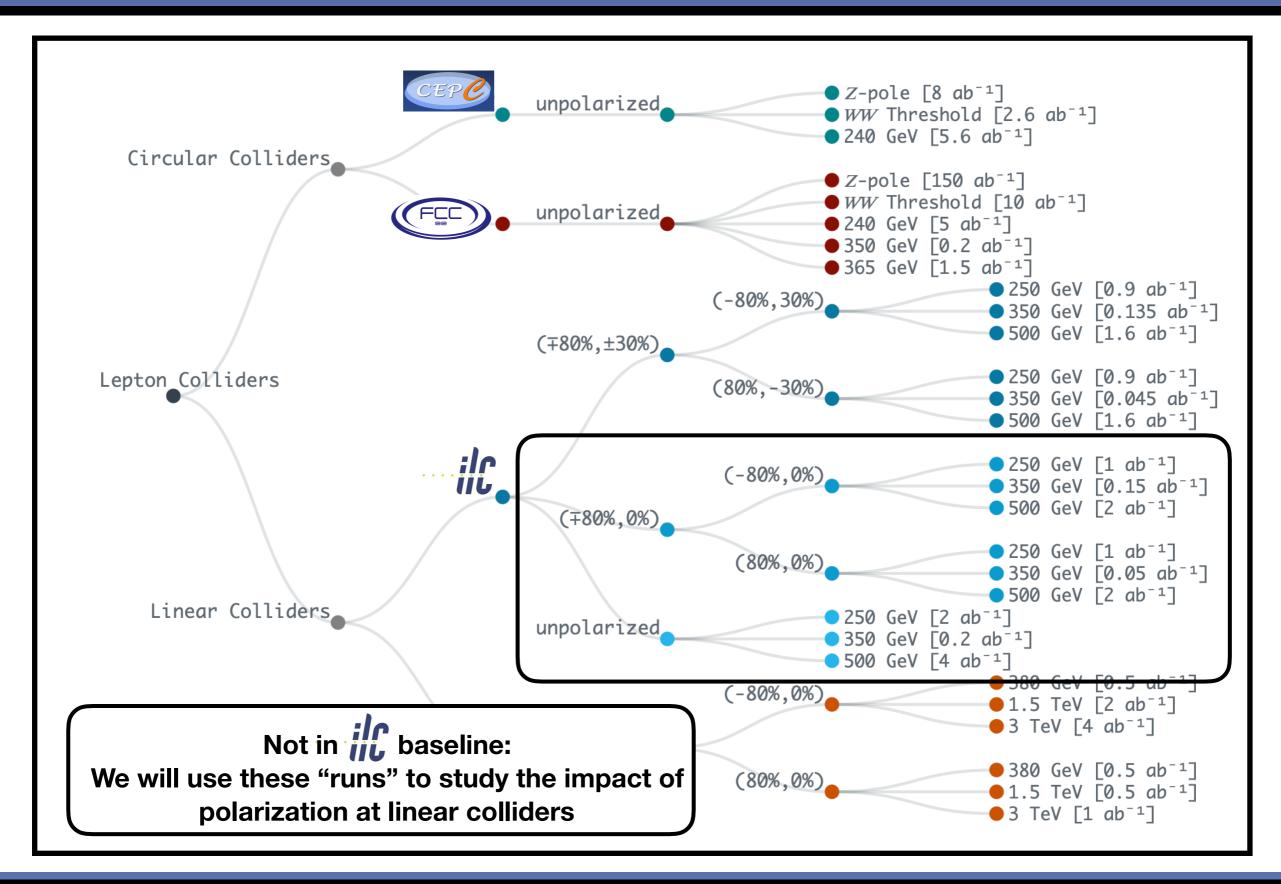
$$g_H^{X~2} \equiv rac{\Gamma_{H o X}}{\Gamma_{H o X}^{ ext{SM}}} \qquad \Gamma_{Z o e^+e^-} = rac{lpha~M_Z}{6\sin^2 heta_w\cos^2 heta_w} (|g_L^{Ze}|^2 + |g_R^{Ze}|^2), \qquad A_e = rac{|g_L^{Ze}|^2 - |g_R^{Ze}|^2}{|g_L^{Ze}|^2 + |g_R^{Ze}|^2}$$

SMEFT-fit Results presented in terms of

$$\left\{\delta g_{H}^{ZZ},\ \delta g_{H}^{WW},\ \delta g_{H}^{\gamma\gamma},\ \delta g_{H}^{Z\gamma},\ \delta g_{H}^{gg},\ \delta g_{H}^{tt},\ \delta g_{H}^{cc},\ \delta g_{H}^{bb},\ \delta g_{H}^{ au au},\ \delta g_{H}^{\mu\mu},\ \delta g_{1z},\delta \kappa_{\gamma},\ \lambda_{z}
ight\} \ rac{15\ Vff/HVff}{\left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{L}^{Zd})_{q_{i}},(\delta g_{L}^{W\ell
u})_{\ell},(\delta g_{L}^{Ze})_{\ell},(\delta g_{R}^{Zu})_{q_{i}},(\delta g_{R}^{Zd})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}
ight\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{L}^{Zd})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{R}^{Zd})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{R}^{Zd})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{R}^{Zd})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{i}},(\delta g_{R}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{1}},(\delta g_{L}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{2}},(\delta g_{L}^{Ze})_{\ell}\right\}_{q_{1}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{2}},(\delta g_{L}^{Ze})_{\ell}\right\}_{q_{2}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{2}},(\delta g_{L}^{Ze})_{q_{2}}\right\}_{q_{2}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{2}},(\delta g_{L}^{Ze})_{q_{2}}\right\}_{q_{2}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q_{2}},(\delta g_{L}^{Zu})_{q_{2}}\right\}_{q_{2}=q_{2},\ \ell=e,\mu,\tau} \ \left\{(\delta g_{L}^{Zu})_{q$$







Official inputs available for Higgs/WW/EW

<u>Higgs</u>

Rates (signal strength)

$$\mu \equiv rac{\sigma \cdot \mathrm{BR}}{\sigma^{\mathrm{SM}} \cdot \mathrm{BR}^{\mathrm{SM}}}$$

(Inclusive) cross section

$$\sigma_{ZH} \equiv \sigma(e^+e^- \to ZH)$$

Only possible at lepton colliders

aTGC

$$\delta g_{1z}, \delta \kappa_{\gamma}, \lambda_{z}$$

EWPO

$$M_Z,\; \Gamma_Z,\; \stackrel{=}{\Gamma_{Z o f}},\; A^f_{FB,LR},\; \ldots$$

$$M_W,\; \Gamma_W,\; \Gamma_{W o f}$$

Z physics via Z-pole:

$$\sqrt{s} = M_Z: e^+e^- \rightarrow Z \rightarrow X$$

or Rad. Return:

$$\sqrt{s} > M_Z: \ e^+e^- o \gamma Z o \gamma X$$

	Higgs	aTGC	EWPO
FCC-ee	Yes (μ, σ _{ZH}) (Complete with HL-LHC)	Yes (aTGC dom.)	Yes (Z-pole run)
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CLIC	Yes (μ, σ _{ZH})	Yes (Full EFT parameterization)	Yes (Rad. Return) (Giga-Z? Not in baseline)

We will always combine with the info expected at the end of the (HL-)LHC era

Official inputs available for Higgs/WW/EW

Higgs

Rates (signal strength)

$$\mu \equiv rac{\sigma \cdot \mathrm{BR}}{\sigma^{\mathrm{SM}} \cdot \mathrm{BR}^{\mathrm{SM}}}$$

(Inclusive) cross section

$$\sigma_{ZH} \equiv \sigma(e^+e^- o ZH)$$

Only possible at lepton colliders

aTGC

$$\delta g_{1z}, \delta \kappa_{\gamma}, \lambda_{z}$$

EWPO

$$M_Z,\; \Gamma_Z,\; \stackrel{=_{III}}{\Gamma_{Z
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or Rad. Return:

 $\sqrt{s} > M_Z: e^+e^- \to \gamma Z \to \gamma X$

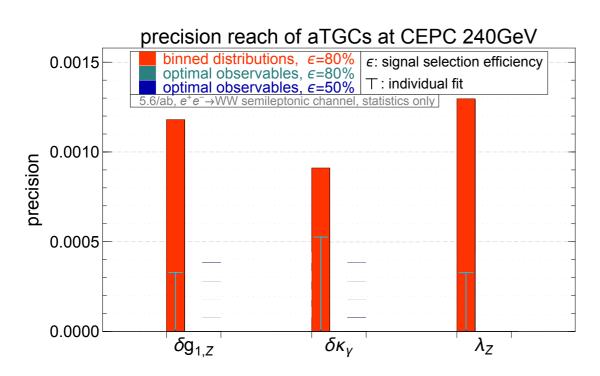
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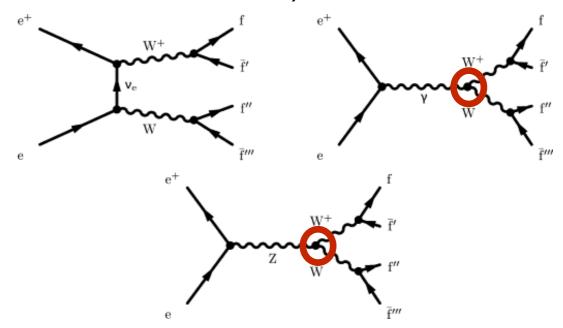
We will always combine with the info expected at the end of the (HL-)LHC era

No full EFT studies available for WW processes at future lepton colliders (Except for CLIC, which uses a preliminary version of this study)

WW production at lepton colliders

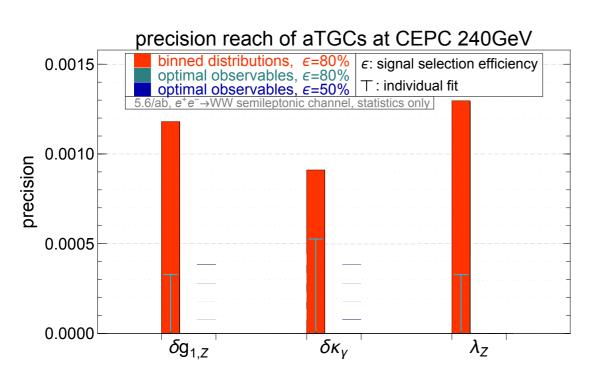
 Current projections based on sensitivity to aTGC ONLY in differential angular distributions (ignoring correlations between bins)

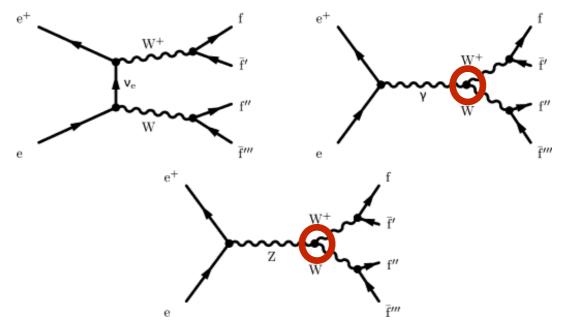




WW production at lepton colliders

 Current projections based on sensitivity to aTGC ONLY in differential angular distributions (ignoring correlations between bins)





 We prepared a new sensitivity study using full information about each event in the formalism of "optimal statistical observables"

Optimal Statistical Observables (OO)

• Consider a Phase-space distribution linear in some coefficients c_i :

$$S(\Phi) = S_0(\Phi) + \sum_i c_i S_i(\Phi)$$

SMEFT: $S(\Phi)=rac{d\sigma}{d\Phi}$ $S_0(\Phi)=rac{d\sigma}{d\Phi}igg|_{ ext{SM}}$ $c_iS_i(\Phi)=rac{d\sigma}{d\Phi}igg|_{ ext{Interf. SM-NP}}$

In the limit of large statistics, the observables

$$O_i = \sum\limits_{k \in ext{events}} rac{S_i(\Phi_k)}{S_0(\Phi_k)}$$
 (See e.g., Z.Phys. C62 (1994) 397-412 Diehl & Nachtmann)

provide the most precise statistical information about the coefficients c_i around the point $c_i=0$, $\forall i$

$$ext{cov}(c_i,c_j) = \left(\mathcal{L} \int d\Phi rac{S_i(\Phi)S_j(\Phi)}{S_0(\Phi)}
ight)^{-1} + \mathcal{O}(c_k)$$

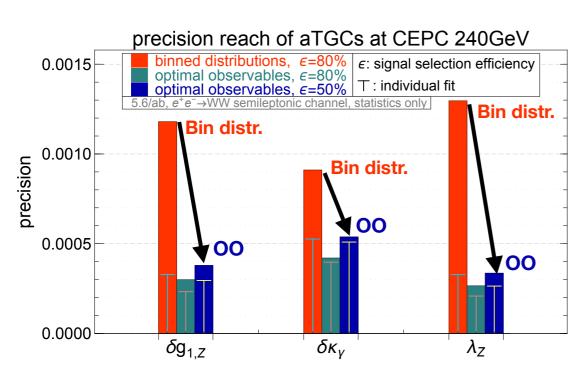


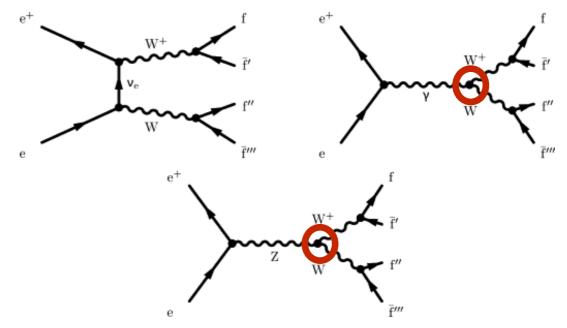
• Idealized (no systematics) \Rightarrow We compensate omission of systematics via conservative selection efficiency ε

$$\mathcal{L}\longrightarrow \varepsilon\mathcal{L}$$

WW production at lepton colliders

 Current projections based on sensitivity to aTGC ONLY in differential angular distributions (ignoring correlations between bins)

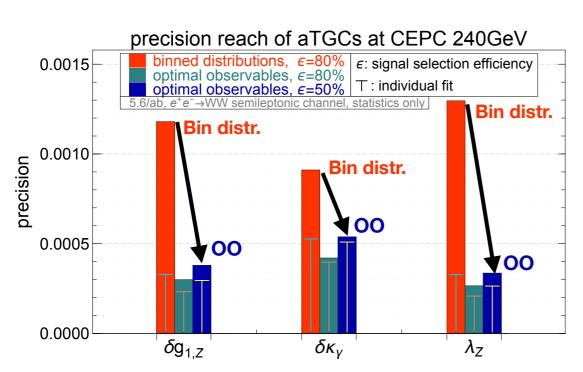


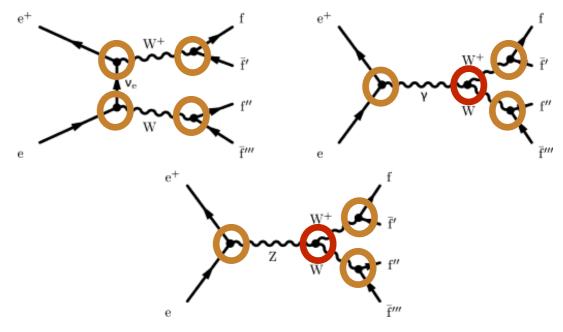


- We prepared a new sensitivity study using full information about each event in the formalism of "optimal statistical observables"
 - Default method only accounts for statistical sensitivity \Rightarrow Compensate omission of systematics via conservative selection efficiency ε

WW production at lepton colliders

 Current projections based on sensitivity to aTGC ONLY in differential angular distributions (ignoring correlations between bins)





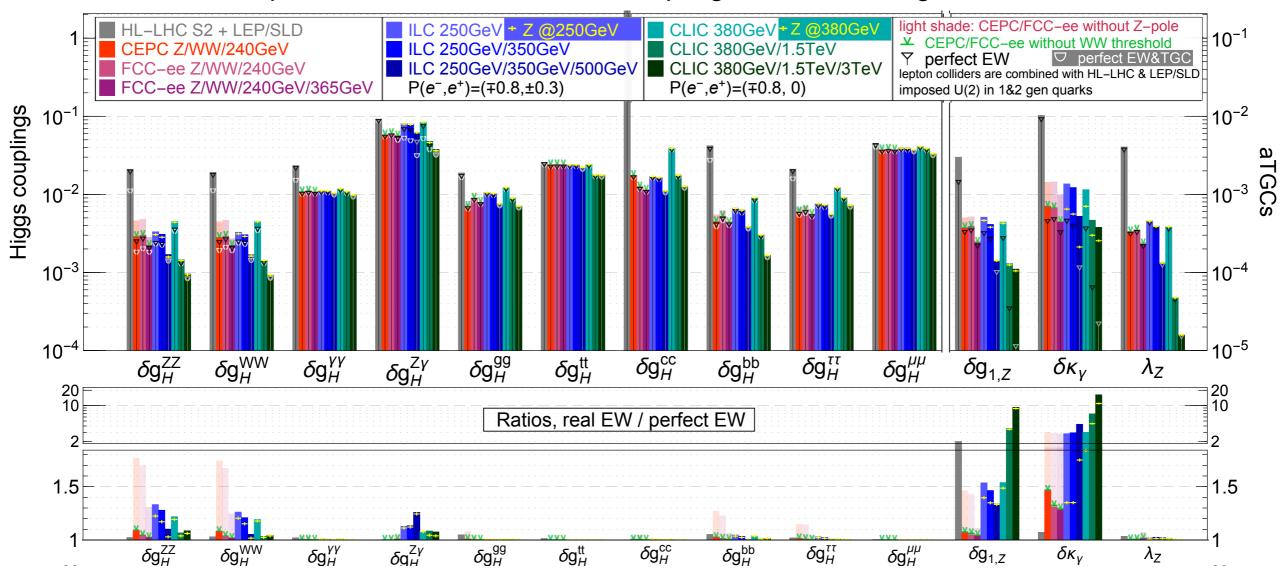
- We prepared a new sensitivity study using full information about each event in the formalism of "optimal statistical observables":
 - We also consider all possible BSM deformations within the dim-6 SMEFT framework



Global fit to EW/Higgs projections

EFT Higgs couplings and aTGC

precision reach on effective couplings from full EFT global fit

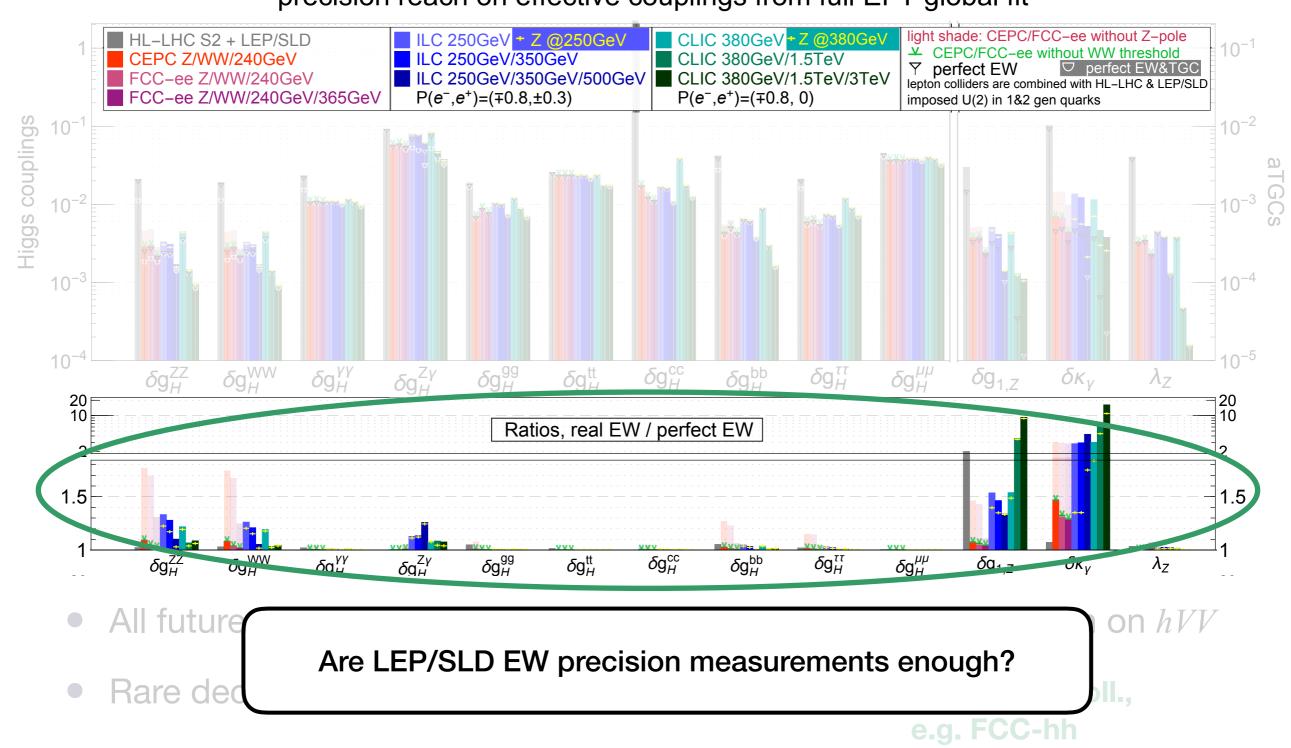


- All future lepton colliders can reach near per-mile level precision on hVV
- Rare decays $(H \rightarrow Z\gamma, H \rightarrow \mu\mu)$ statistically limited \leftarrow Better at pp coll., e.g. FCC-hh

Global fit to EW/Higgs projections

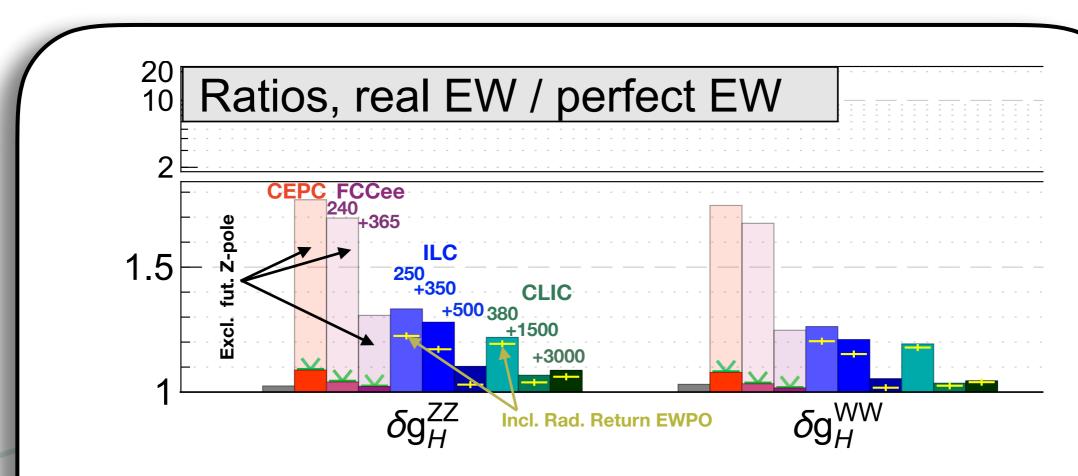
EFT Higgs couplings and aTGC

precision reach on effective couplings from full EFT global fit



Global fit to EW/Higgs projections

EFT Higgs couplings and aTGC



 Note: polarization partially compensates the absence of Z-pole at linear colliders

All future

Are LEP/SLD EW precision measurements enough?

Rare ded

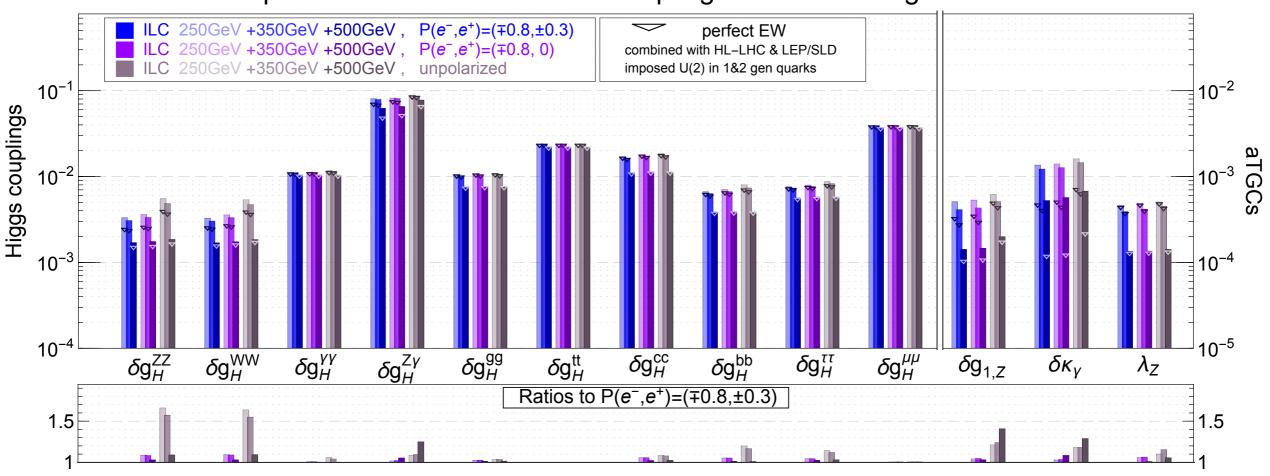
e.g. FCC-hh

on hVV

On the beam polarization at linear colliders

Impact of polarization in Higgs coupling sensitivity

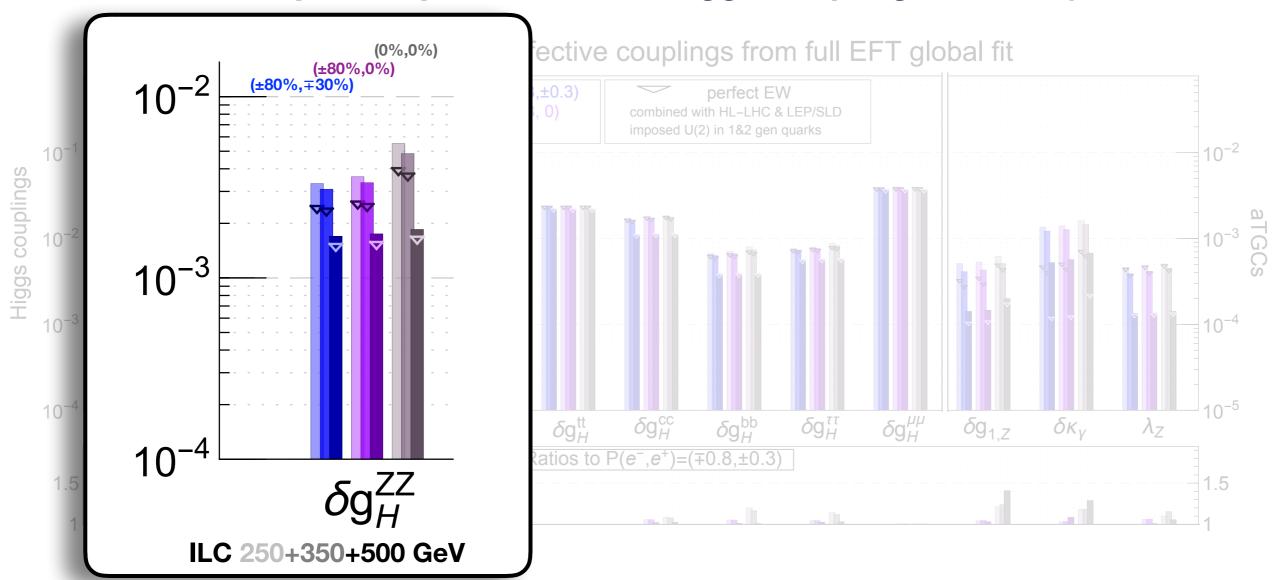
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 Polarization can resolve degeneracies in the ZH rate appearing in the unpolarized case.

On the beam polarization at linear colliders

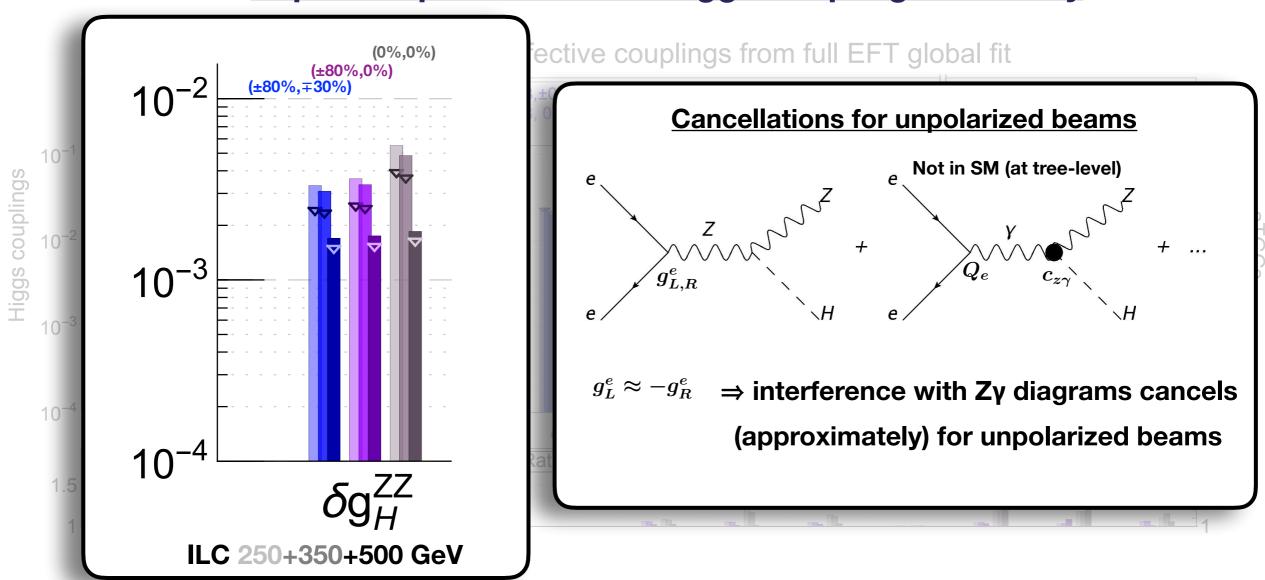
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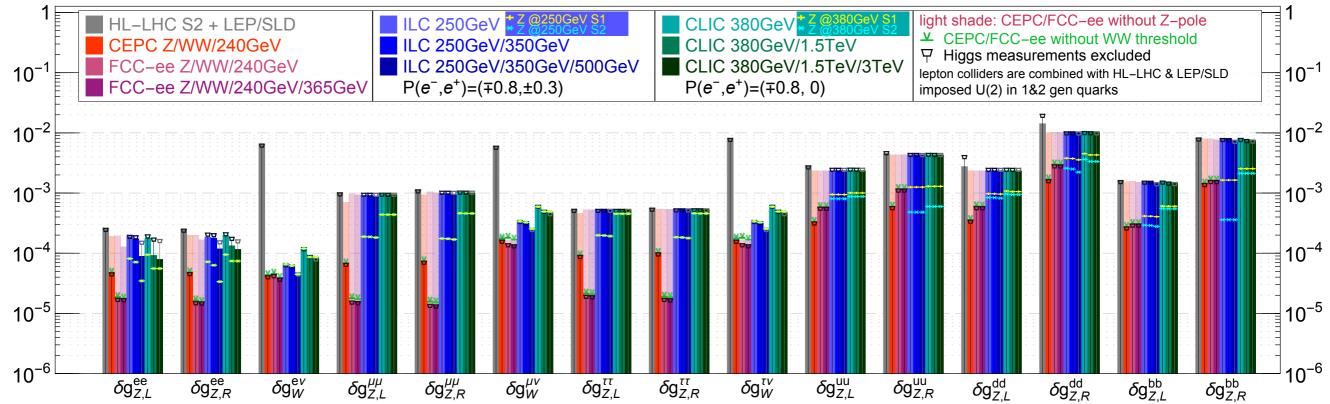
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The other "half" of the EFT fit: EW couplings

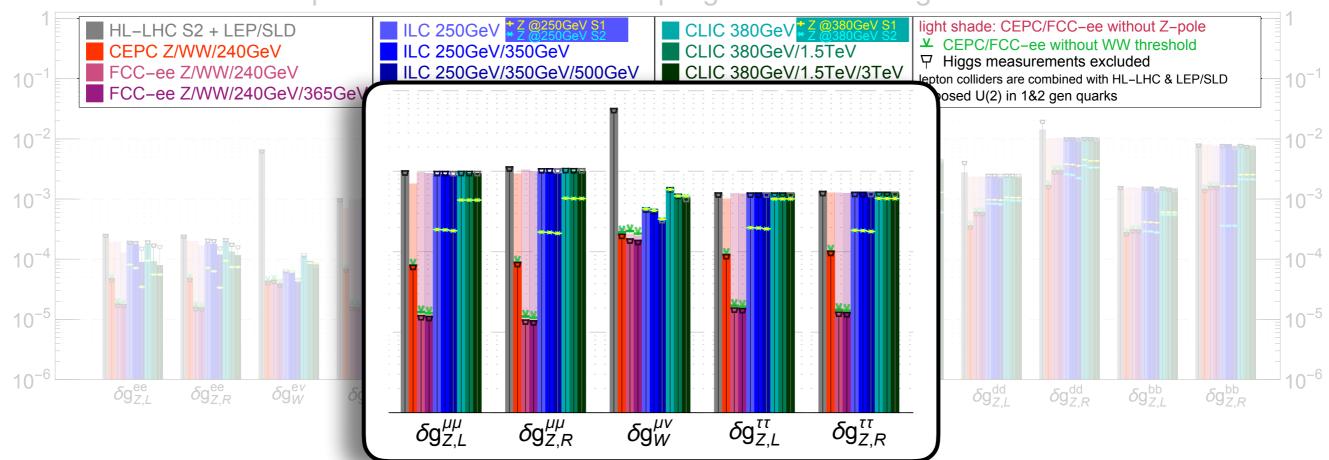
precision reach on EW couplings from full EFT global fit



- EW Vff couplings (and therefore HVff) better determined at circular colliders due to the Tera-Z factory run vs. the use of rad. return measurements at linear colliders
- Conversely, high-energy runs at linear colliders can measure *HZee* (and therefore *Zee*) via Higgs, but only in combination with other *HZZ* interactions

The other "half" of the EFT fit: EW couplings

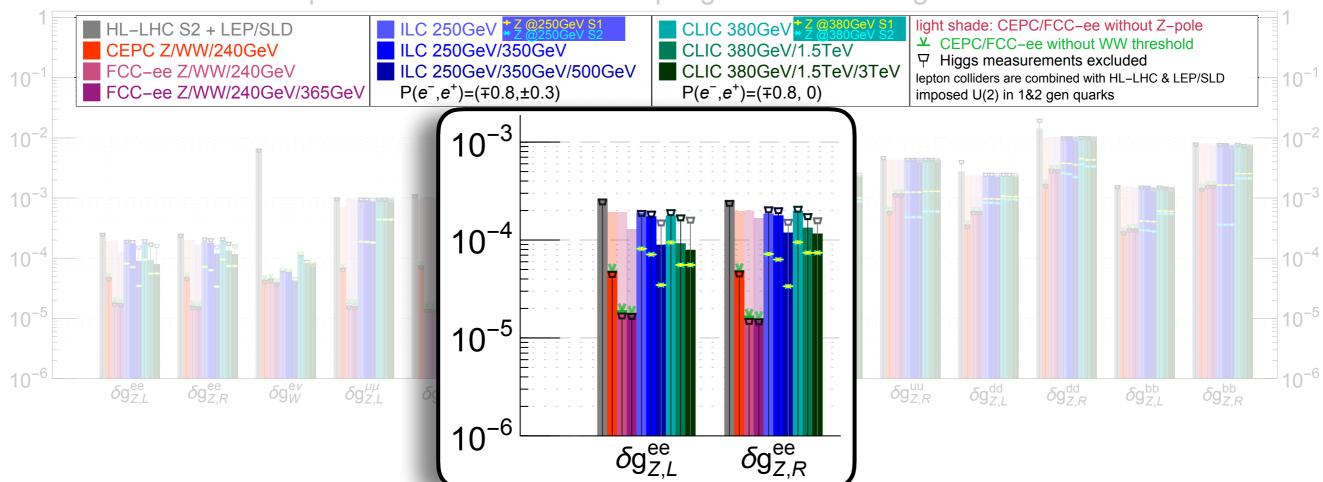
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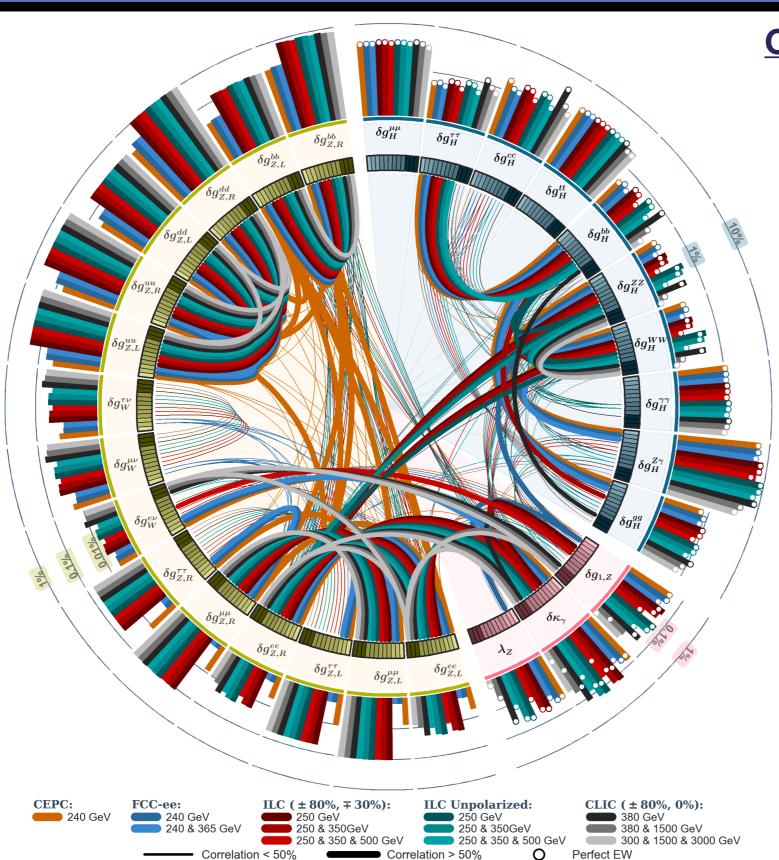
- EW Vff couplings (and therefore HVff) better determined at circular colliders due to the Tera-Z factory run vs. the use of rad. return measurements at linear colliders
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The other "half" of the EFT fit: EW couplings

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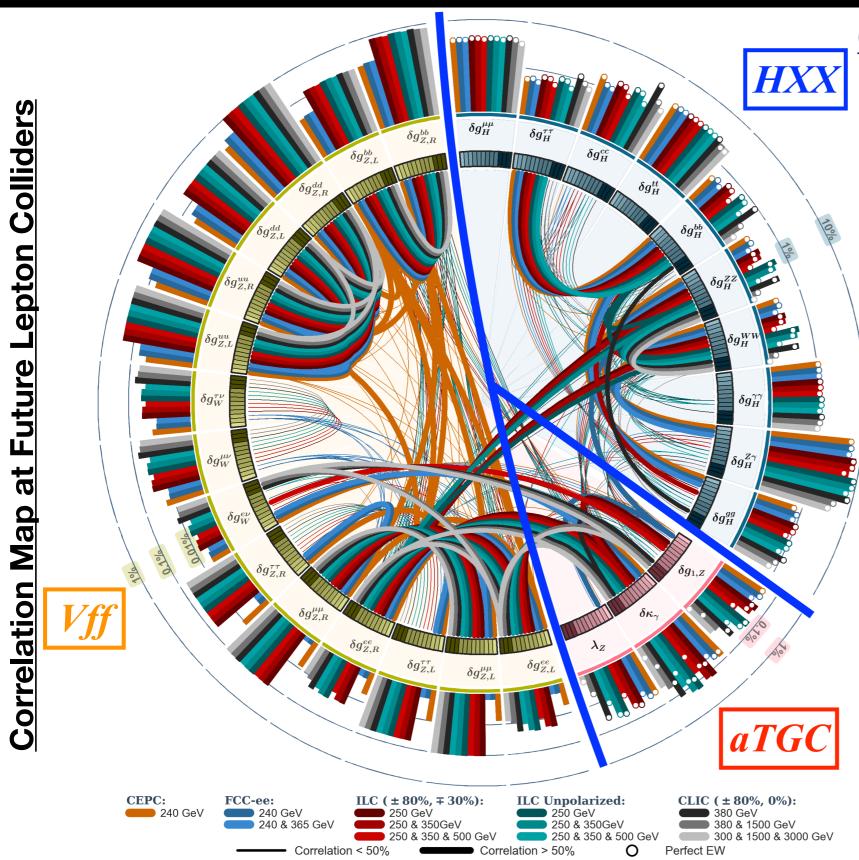


- EW Vff couplings (and therefore HVff) better determined at circular colliders due to the Tera-Z factory run vs. the use of rad. return measurements at linear colliders
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Couplings and correlations

How to read "this"?

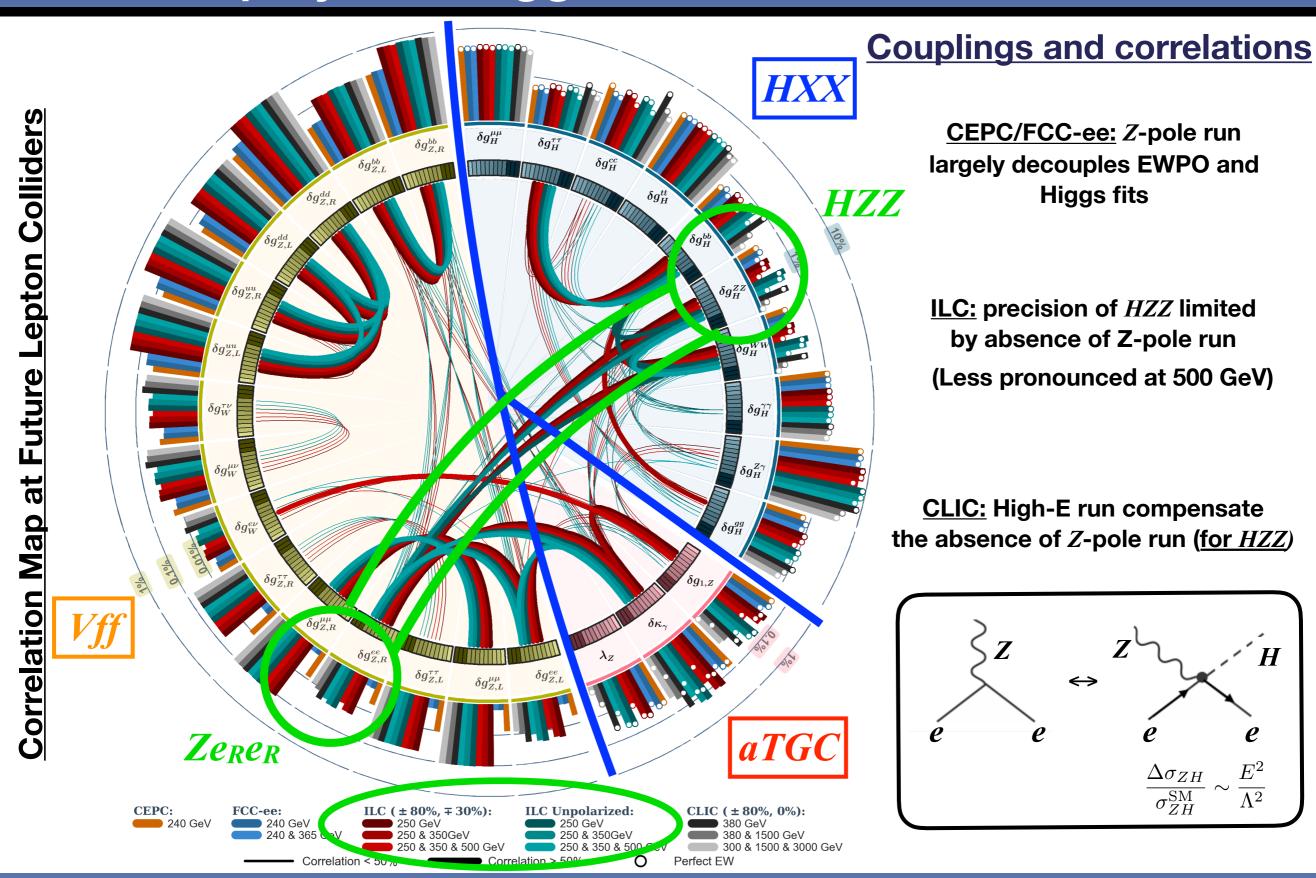


Couplings and correlations

How to read "this"?

On the outside: 1σ uncertainty on the different interactions
Interactions grouped as:
Eff. *H* couplings, *aTGC* and *Vff*

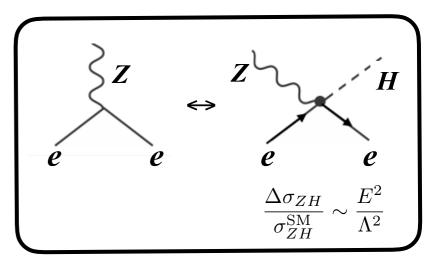
<u>Correlations</u> indicated by lines linking the different couplings



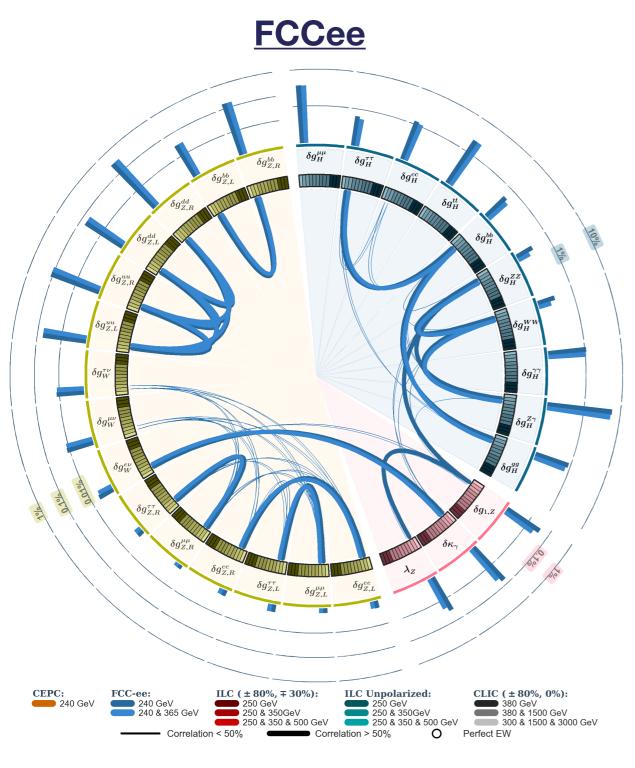
CEPC/FCC-ee: Z-pole run largely decouples EWPO and **Higgs fits**

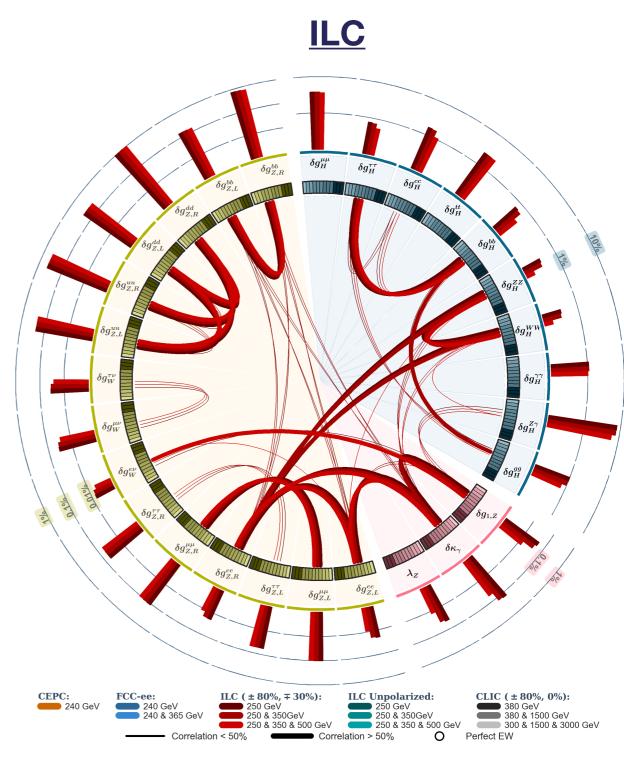
ILC: precision of *HZZ* limited by absence of Z-pole run (Less pronounced at 500 GeV)

CLIC: High-E run compensate the absence of Z-pole run (for HZZ)

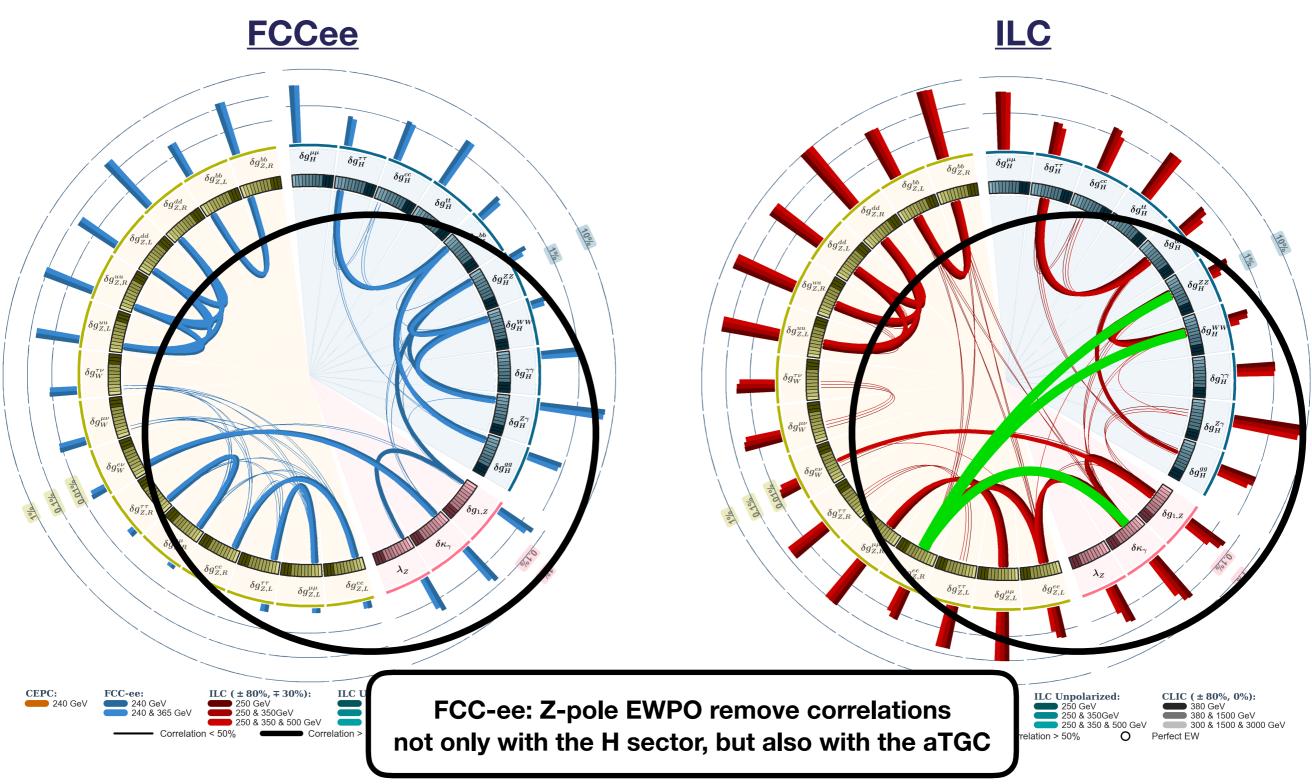


Couplings and correlations





Couplings and correlations





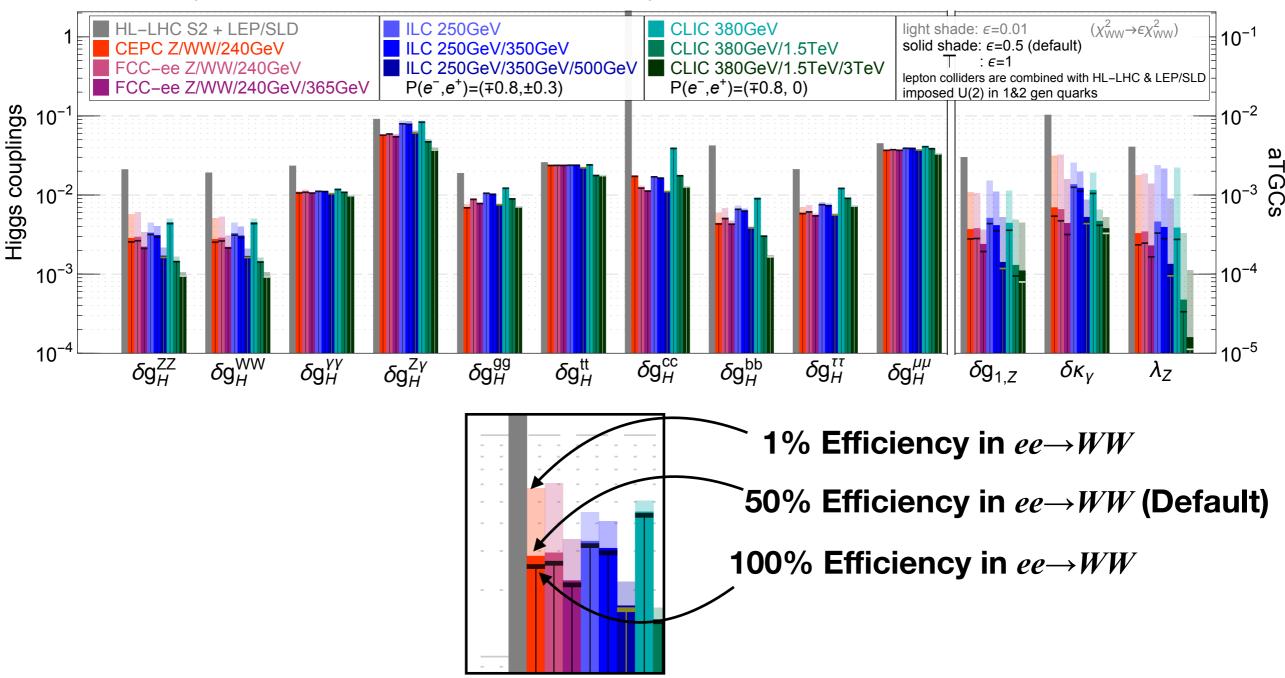
Conclusions

- Motivated by the Higgs factory option, there seems to be a consensus that a future lepton collider must be the next step in particle collider experiments:
 - "Model-independent" determination of H couplings (unlike the HL-LHC)
 - Near per-mille level precision in some H couplings.
 - But rare channels limited by stats ⇒ need Hadron collider afterwards
- But future lepton colliders are more than Higgs factories: possibility of improving the knowledge of ALL EW interactions
- In fact, a precise determination of Higgs properties requires to keep under control uncertainties associated to other EW parameters!
 - We studied the impact of the EW uncertainties adding to the global Higgs
 + EW fit a fully global EFT study of WW at future lepton colliders
- Polarization and higher energies at LC can partially mitigate the impact of the absence of Z-pole run in some couplings (*HZZ*), but cannot the replace the net added value of the EW precision measurements.



EFT Higgs couplings and aTGC: dependence on WW projections

precision reach with different assumptions on $e^+e^- \rightarrow WW$ measurements

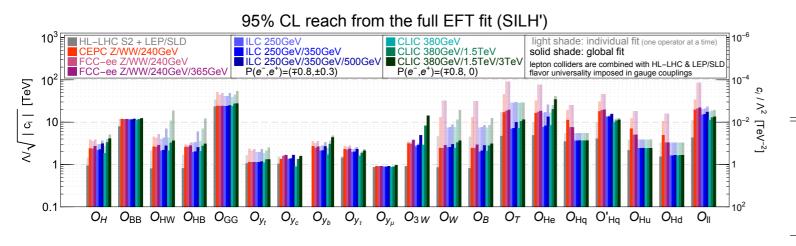


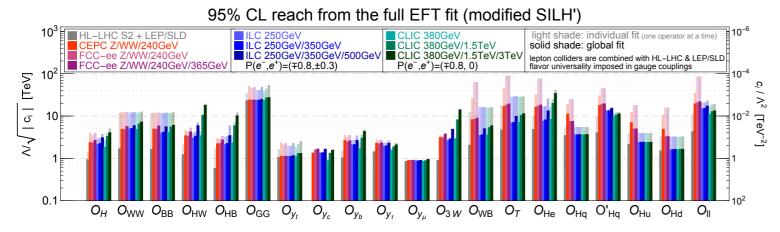
Influence of the assumptions in the 00 study of WW production in the extraction of H couplings

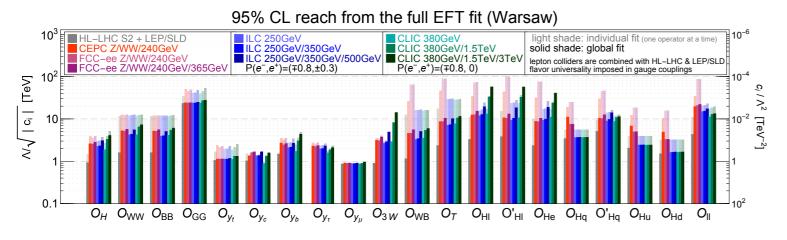


Other EFT results

Results in manifestly gauge-invariant dim-6 bases







Notation

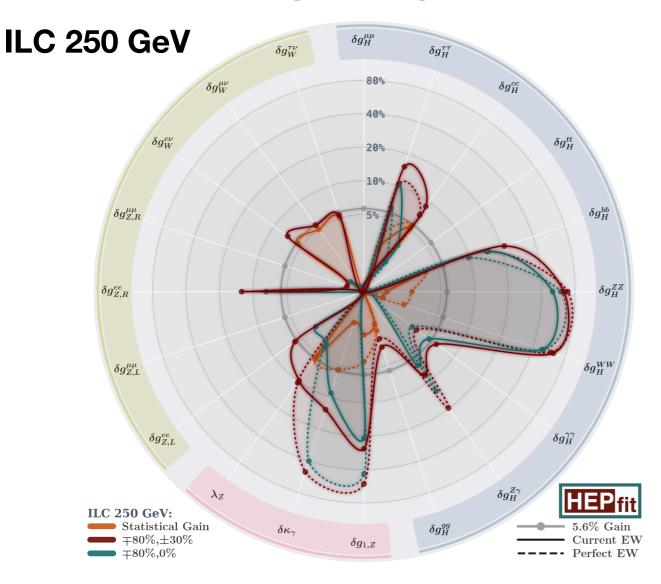
$\mathcal{O}_H = \frac{1}{2} (\partial_\mu H^2)^2$	$\mathcal{O}_{GG} = g_s^2 H ^2 G_{\mu\nu}^A G^{A,\mu\nu}$
$\mathcal{O}_{WW} = g^2 H ^2 W^a_{\mu\nu} W^{a,\mu\nu}$	$\mathcal{O}_{y_u} = y_u H ^2 \bar{q}_L \tilde{H} u_R + \text{h.c.} (u \to t, c)$
$\mathcal{O}_{BB} = g^{\prime 2} H ^2 B_{\mu\nu} B^{\mu\nu}$	$\mathcal{O}_{y_d} = y_d H ^2 \bar{q}_L H d_R + \text{h.c.} (d \rightarrow b)$
$\mathcal{O}_{HW} = ig(D^{\mu}H)^{\dagger}\sigma^{a}(D^{\nu}H)W^{a}_{\mu\nu}$	$\mathcal{O}_{y_e} = y_e H ^2 ar{l}_L H e_R + ext{h.c.} (e ightarrow au, \mu)$
$\mathcal{O}_{HB} = ig'(D^{\mu}H)^{\dagger}(D^{\nu}H)B_{\mu\nu}$	$\mathcal{O}_{3W} = \frac{1}{3!} g \epsilon_{abc} W^{a \nu}_{\mu} W^{b \nu}_{\nu \rho} W^{c \rho \mu}$
$\mathcal{O}_W = \frac{ig}{2} (H^{\dagger} \sigma^a \overleftrightarrow{D_{\mu}} H) D^{\nu} W^a_{\mu\nu}$	$\mathcal{O}_B = \frac{ig'}{2} (H^{\dagger} \overleftrightarrow{D_{\mu}} H) \partial^{\nu} B_{\mu\nu}$
$\mathcal{O}_{WB} = gg'H^{\dagger}\sigma^{a}HW^{a}_{\mu\nu}B^{\mu\nu}$	$\mathcal{O}_{H\ell} = iH^{\dagger} \overleftrightarrow{D}_{\mu} H \bar{\ell}_L \gamma^{\mu} \ell_L$
$\mathcal{O}_T = rac{1}{2} (H^\dagger \overrightarrow{D_\mu} H)^2$	$\mathcal{O}_{H\ell}^{\prime}=iH^{\dagger}\sigma^{a}\overrightarrow{D_{\mu}}Har{\ell}_{L}\sigma^{a}\gamma^{\mu}\ell_{L}$
$\mathcal{O}_{\ell\ell} = (\bar{\ell}_L \gamma^\mu \ell_L)(\bar{\ell}_L \gamma_\mu \ell_L)$	$\mathcal{O}_{He} = iH^{\dagger} \overrightarrow{D_{\mu}} H \bar{e}_R \gamma^{\mu} e_R$
$\mathcal{O}_{Hq} = iH^{\dagger} \overleftrightarrow{D}_{\mu} H \bar{q}_L \gamma^{\mu} q_L$	$\mathcal{O}_{Hu} = iH^{\dagger} \overleftrightarrow{D_{\mu}} H \bar{u}_R \gamma^{\mu} u_R$
$\mathcal{O}'_{Hq} = iH^{\dagger}\sigma^{a} \overleftrightarrow{D_{\mu}} H \bar{q}_{L} \sigma^{a} \gamma^{\mu} q_{L}$	$\mathcal{O}_{Hd} = iH^{\dagger} \stackrel{\longleftrightarrow}{D_{\mu}} H \bar{d}_R \gamma^{\mu} d_R$

Fits assuming flavour universality in O_{Hf} and O'_{Hf}

A bit more on polarization

On the beam polarization at linear colliders

Impact of polarization in Higgs coupling sensitivity



How to read "this"?

 $\delta g_{unpol.}/\delta g_{unpol.(Lx1.12)}$ -1 : Increased stats.

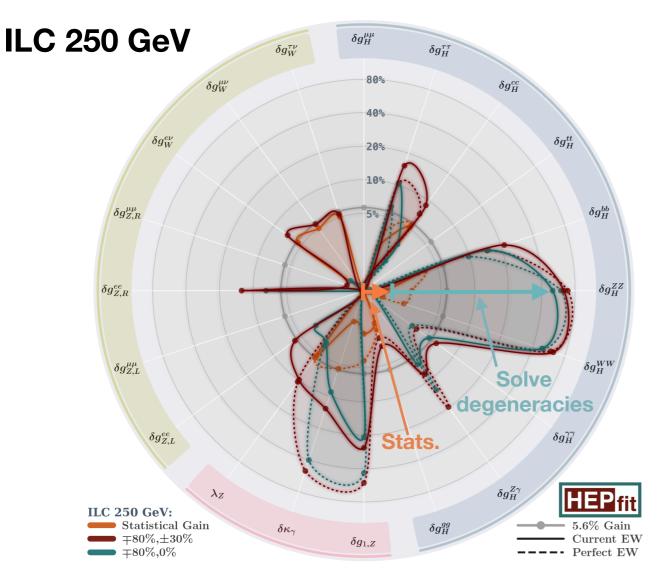
 $\delta g_{unpol.}/\delta g_{(\pm 80\%,0\%)}$ - 1 : Increased stats. + resolving degeneracies

 $\delta g_{unpol.}/\delta g_{(\pm 80\%, \mp 30\%)}$ - 1 : Increased stats. + resolving degeneracies

- Polarization can resolve degeneracies in the ZH rate appearing in the unpolarized case.
- The same can be resolved using data at different energies
 ⇒ negligible influence of polarization in the results at 500 GeV

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Future colliders considered in the study

Official inputs available for Higgs/WW/EW

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(Inclusive) cross section

$$\sigma_{ZH} \equiv \sigma(e^+e^- o ZH)$$

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CEPC	Yes (μ, σ _{ZH}) (Complete with HL-LHC)	Yes (aTGC dom)	Yes	
CLIC	Yes (μ, σ _{ZH})	Yes (Full EFT parameterization)	Yes (Rad. Return) (Giga-Z? Not in baseline)	
FCC-hh	Yes (µ, BR _i /BR _j for rare decays)	From Lepton Coll.	From Lepton Coll.	

We will always combine with the info expected at the end of the (HL-)LHC era

Future colliders considered in the study

Official inputs available for Higgs/WW/EW



Rates (signal strength)

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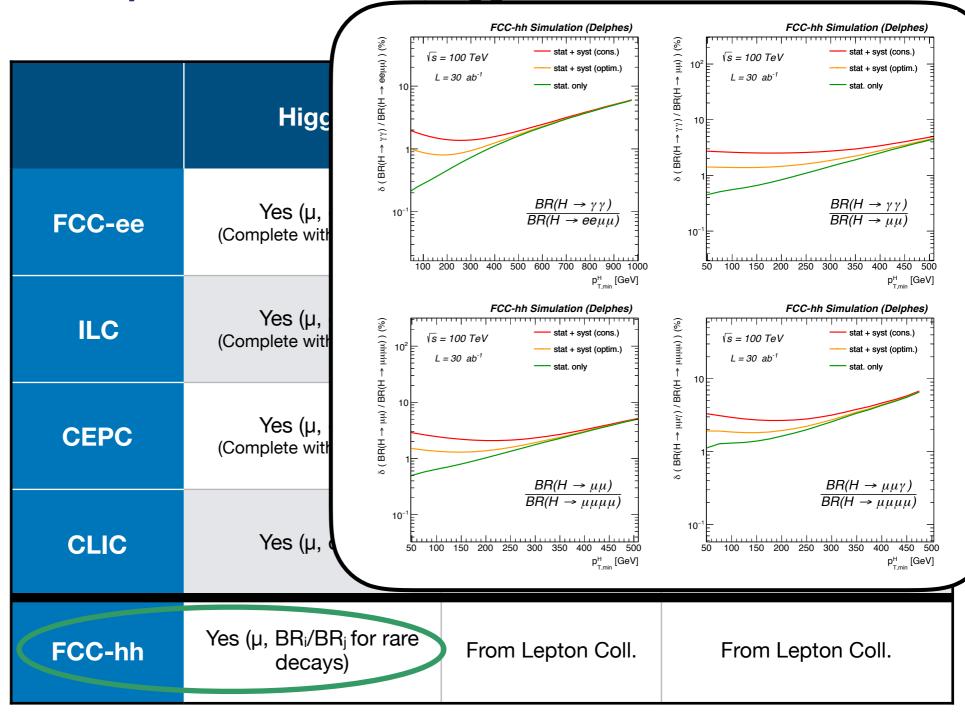
 $M_W,\; \Gamma_W,\; \Gamma_{W o f}$

Z physics via Z-pole:

$$\sqrt{s} = M_Z: e^+e^- \to Z \to X$$

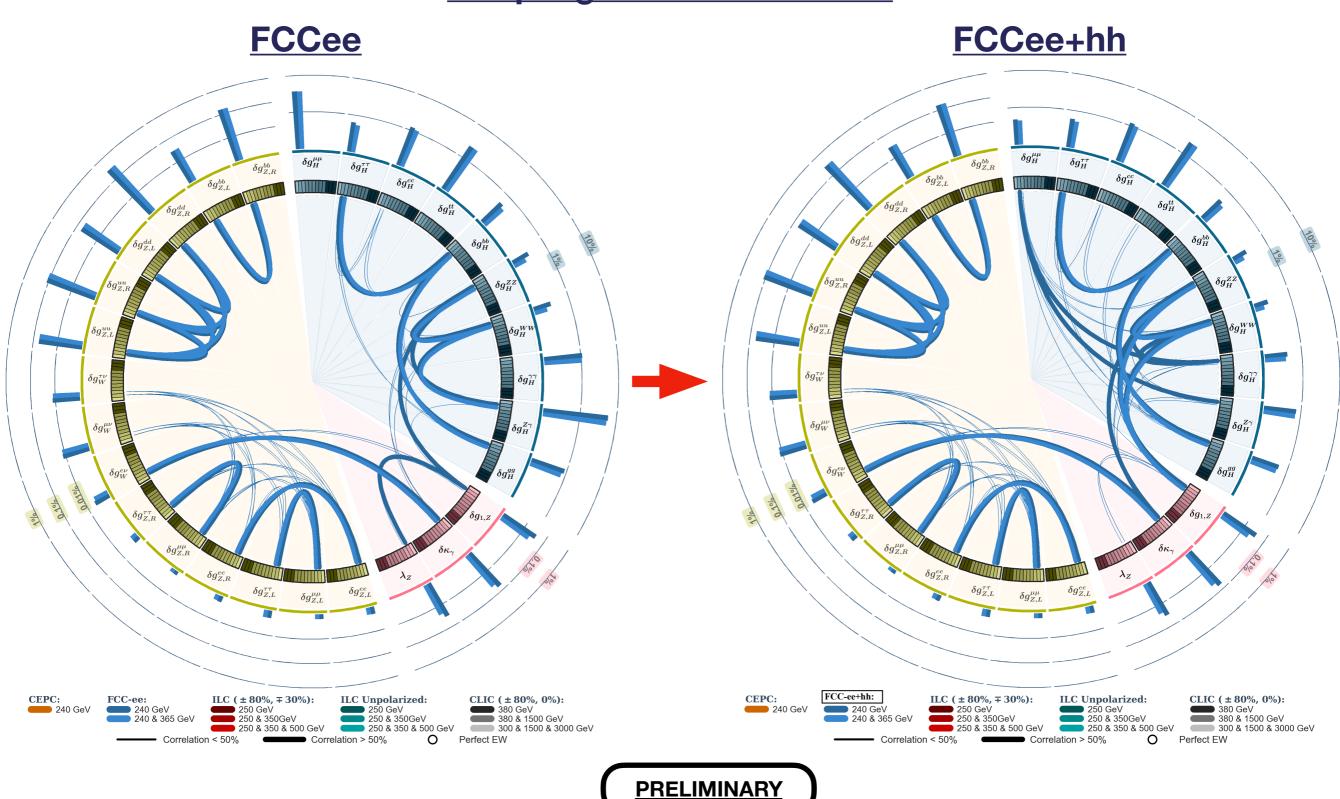
or Rad. Return:

 $\sqrt{s} > M_Z: e^+e^- o \gamma Z o \gamma X$

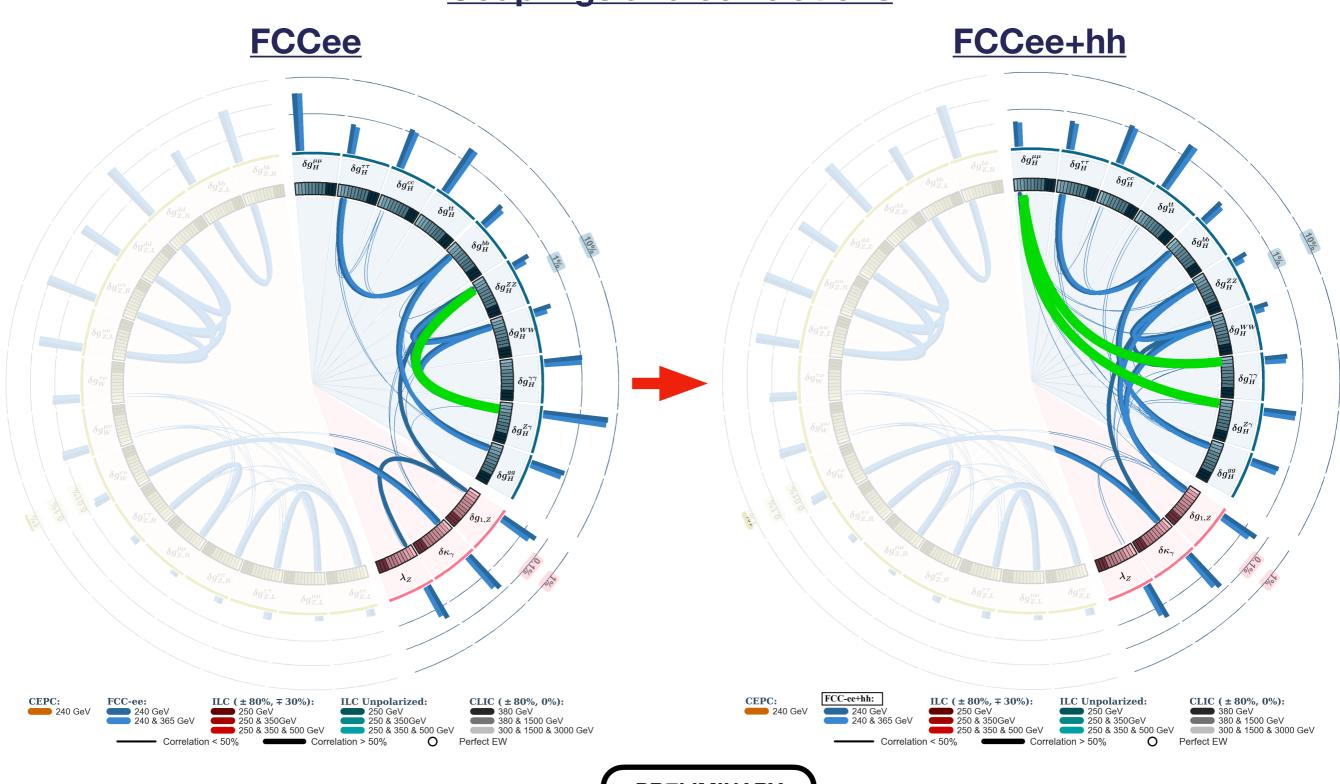


We will always combine with the info expected at the end of the (HL-)LHC era

Couplings and correlations



Couplings and correlations



PRELIMINARY

Future colliders considered in the study

Official inputs available for Higgs/WW/EW

Higgs

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Z physics via Z-pole:

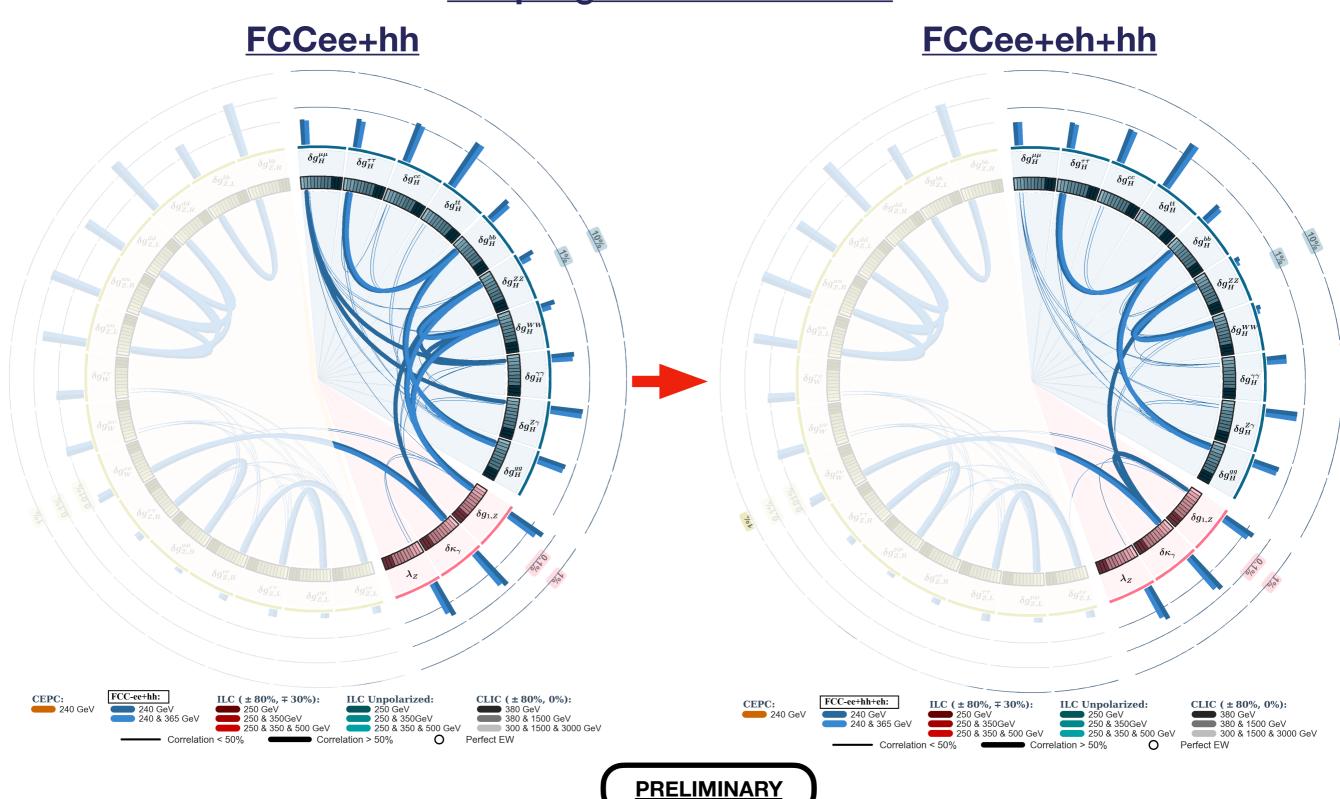
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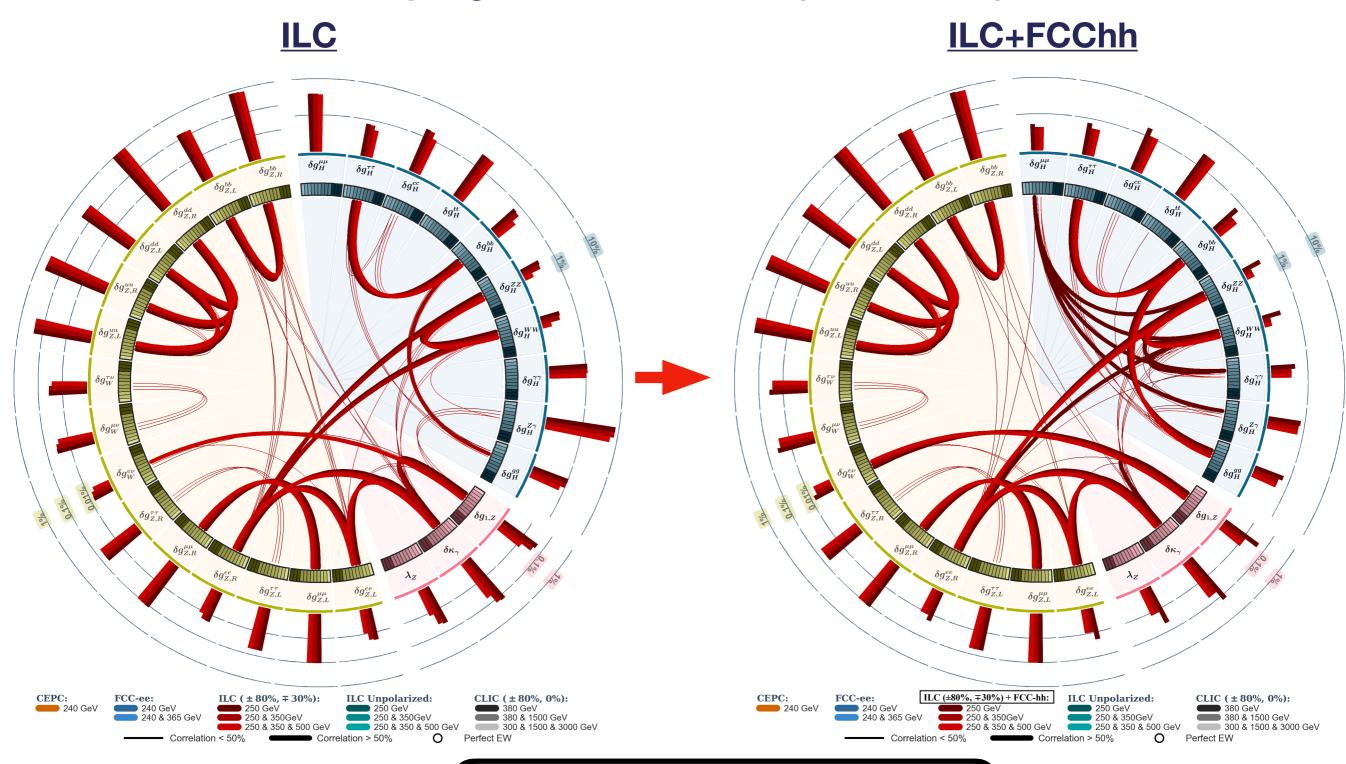
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FCC-hh	Yes (µ, BR _i /BR _j for rare decays)	From Lepton Coll.	From Lepton Coll.	
FCC-eh	Yes (μ: CC & NC DIS → <i>HVV</i>)	From Lepton Coll.	From Lepton Coll.	

We will always combine with the info expected at the end of the (HL-)LHC era

Couplings and correlations



Couplings and correlations (ILC+FCChh)



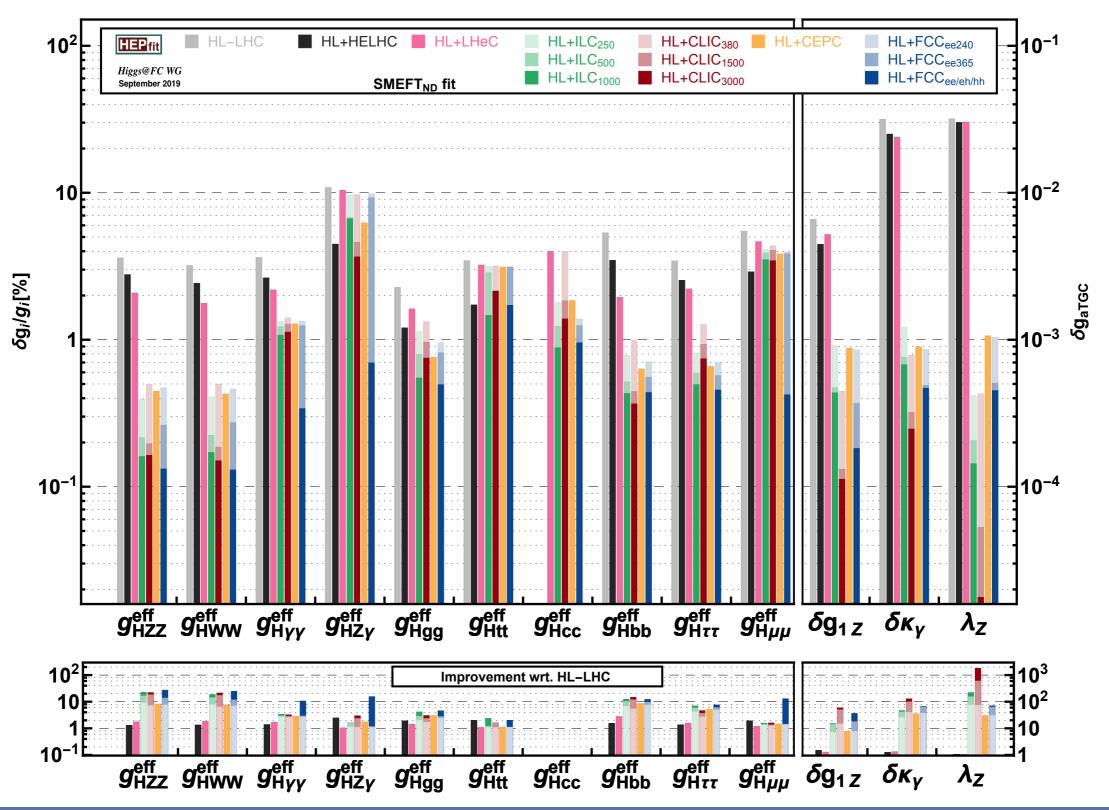
PRELIMINARY: To be updated with FCCeh soon...

Results from Higgs@Future Colliders WG (ESPP)
JB, et al., arXiv:1905.03764 [hep-ph]

Single Higgs couplings

Results in the SMEFT-framework (Higgs/aTGC)

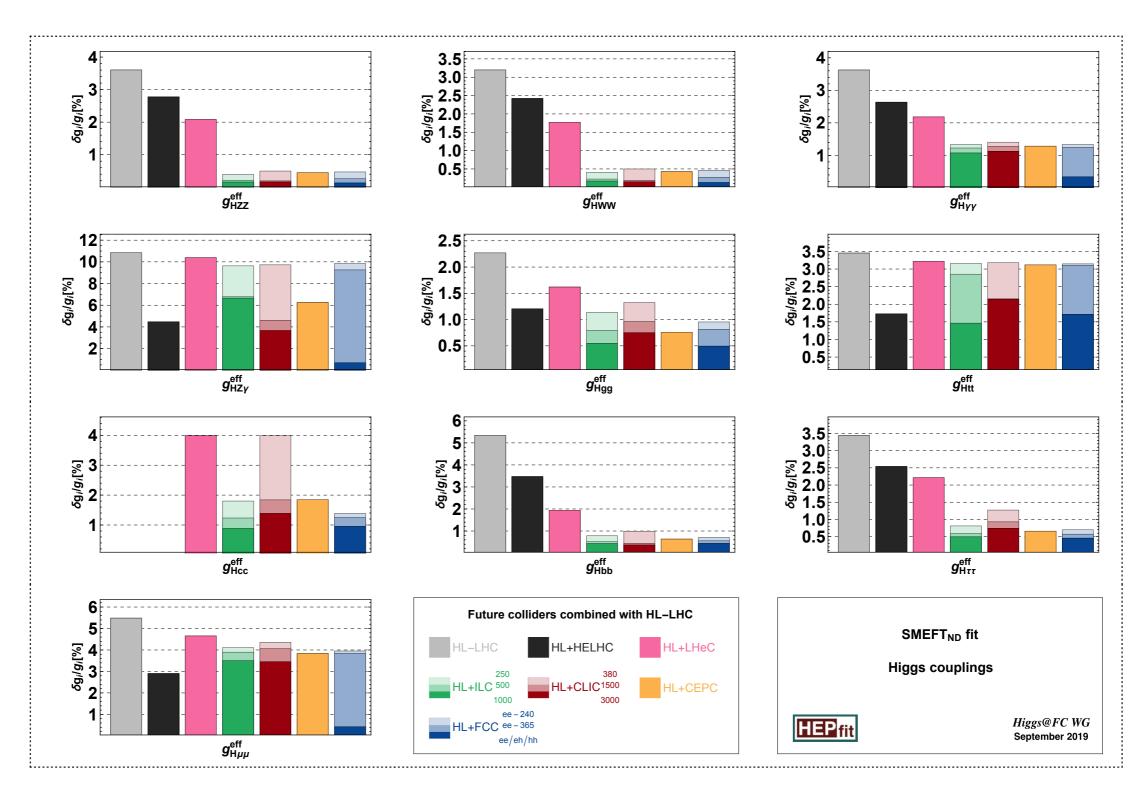




Single Higgs couplings

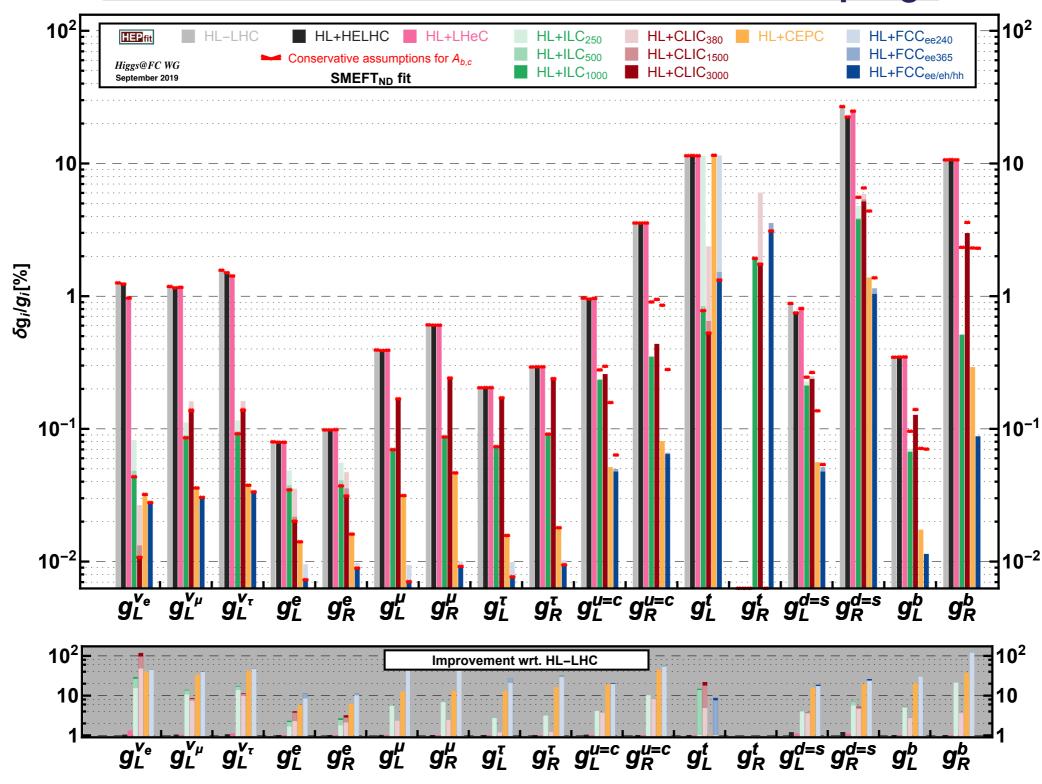
Results in the SMEFT-framework (Higgs)





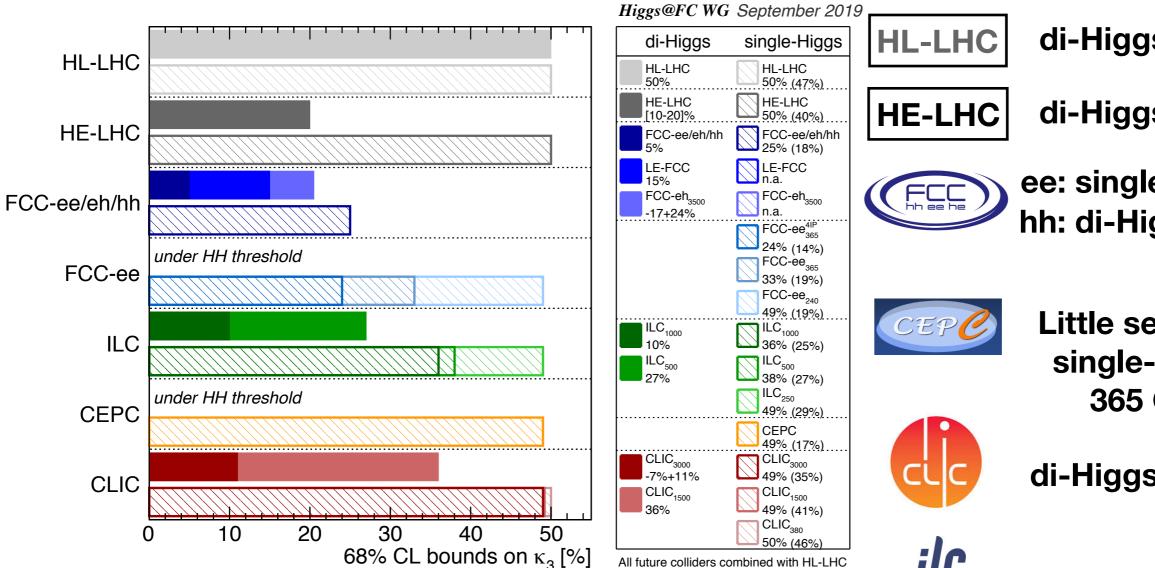
Sensitivity to NP in EW interactions

The other "half" of the SMEFT fit: EW Zff couplings



The Higgs self-coupling

Comparison of capabilities to measure the H³ coupling



di-Higgs ~50%

di-Higgs ~15%

ee: single-Higgs ~34%

hh: di-Higgs ~5-10%

Little sensitivity via single-Higgs w/o 365 GeV run

di-Higgs ~10%

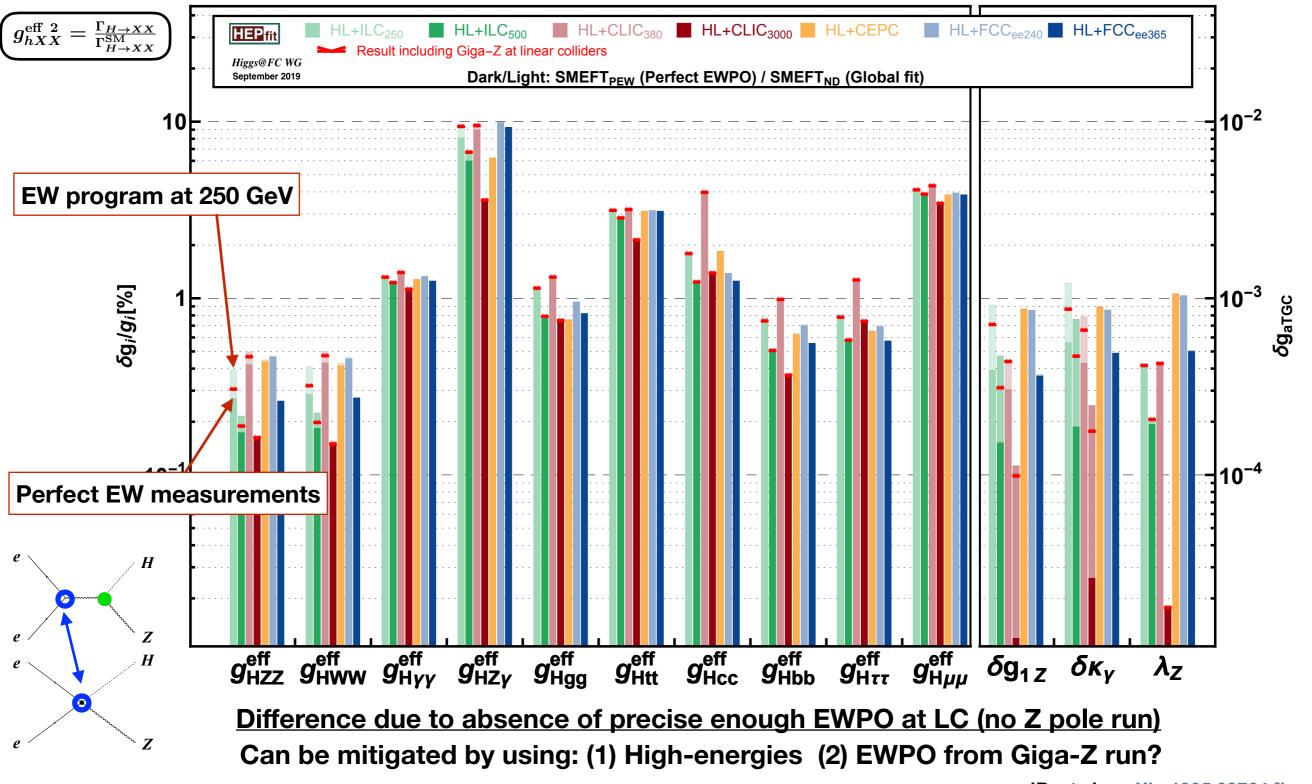


di-Higgs ~27% (10%)

Assuming upgrade to 500 GeV (1000 GeV)

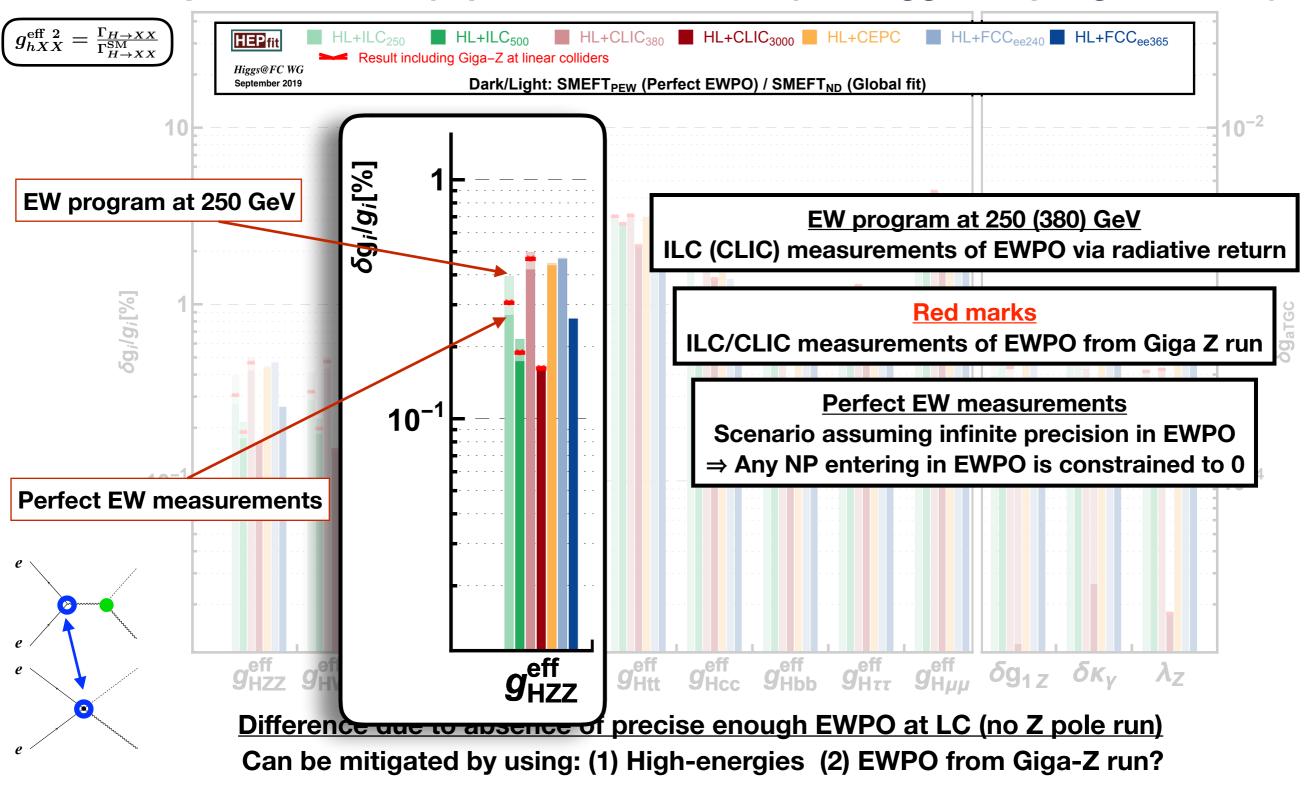
All future colliders combined with HL-LHC

Impact of EWPO (Z pole measurements) in Higgs coupling sensitivity



JB, et al., arXiv:1905.03764 [hep-ph]

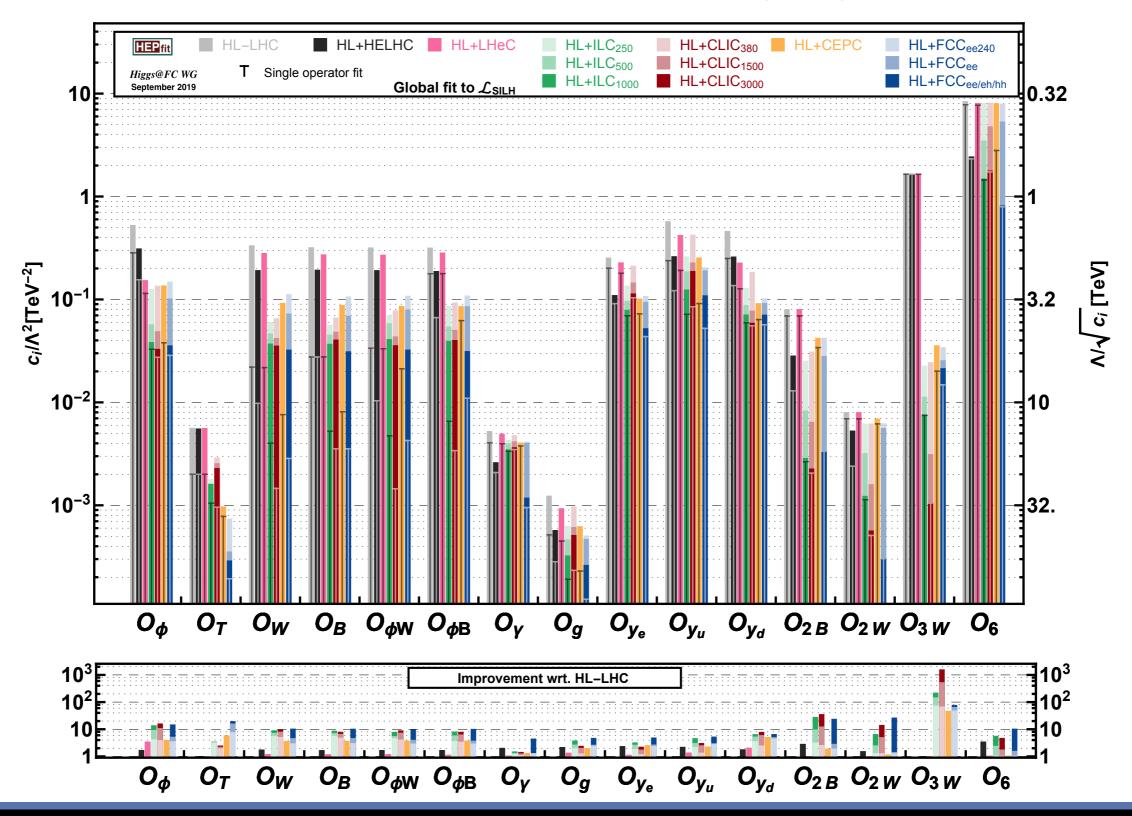
Impact of EWPO (Z pole measurements) in Higgs coupling sensitivity



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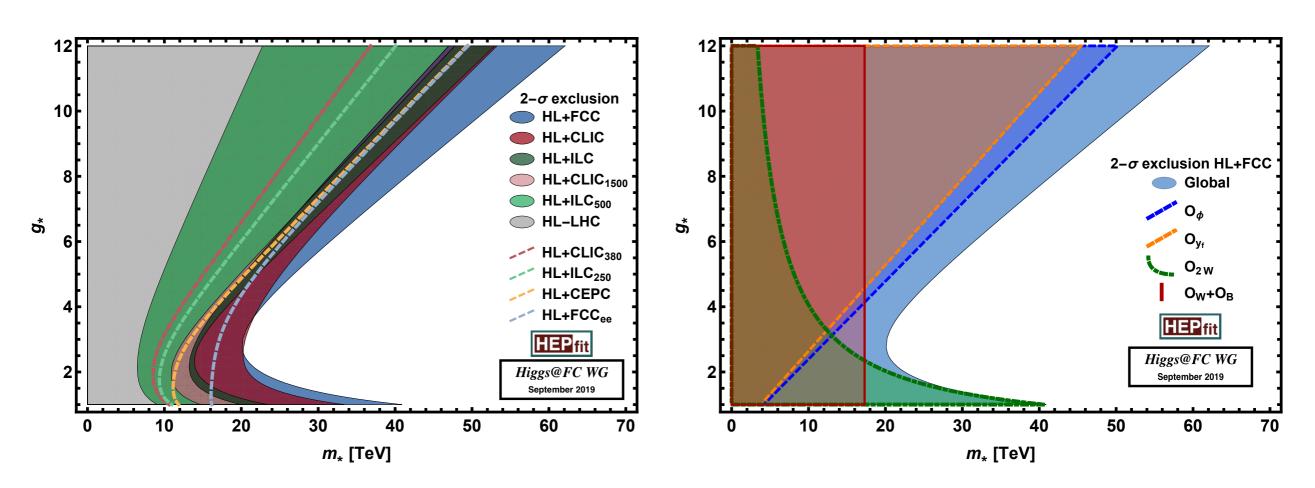
Indirect constraints on Composite Higgs

Constraints on SILH effective Lagrangian



Indirect constraints on Composite Higgs

Projecting into simple Composite Higgs scenarios



Simplified CH benchmark: 1 coupling (g*) - 1 scale (m*)

$$\frac{c_{\phi,6,y_f}}{\Lambda^2} = \frac{g_{\star}^2}{m_{\star}^2}, \qquad \frac{c_{W,B}}{\Lambda^2} = \frac{1}{m_{\star}^2}, \qquad \frac{c_{2W,2B,2G}}{\Lambda^2} = \frac{1}{g_{\star}^2} \frac{1}{m_{\star}^2},$$

$$\frac{c_T}{\Lambda^2} = \frac{y_t^4}{16\pi^2} \frac{1}{m_{\star}^2}, \qquad \frac{c_{\gamma,g}}{\Lambda^2} = \frac{y_t^2}{16\pi^2} \frac{1}{m_{\star}^2}, \qquad \frac{c_{\phi W,\phi B}}{\Lambda^2} = \frac{g_{\star}^2}{16\pi^2} \frac{1}{m_{\star}^2}, \qquad \frac{c_{3W,3G}}{\Lambda^2} = \frac{1}{16\pi^2} \frac{1}{m_{\star}^2}$$

The role of theory

Will SM theory calculations be enough?

Estimates for SM theory uncertainties used in the ESU studies

Decay	Partial width	Projected future unc. $\Delta\Gamma/\Gamma$ [%]				
	$[\mathrm{keV}]$	$ m Th_{Intr}$	$\mathbf{Th}_{\mathrm{Par}}(m_q)$	$\mathbf{Th}_{\mathrm{Par}}(lpha_s)$	$\mathbf{Th}_{\mathrm{Par}}(m_{\mathrm{H}})$	
$H o bar{b}$	2379	0.2	0.6^{\flat}	$< 0.1^{\sharp}$	_	
$H o au^+ au^-$	256	< 0.1	_	_	-	
H o c ar c	118	0.2	1.0^{\flat}	$< 0.1^{\sharp}$	_	
$H o \mu^+\mu^-$	0.89	< 0.1	_	_	_	
$H o WW^*$	883	$\lesssim 0.4$	_	_	0.1^{\ddagger}	
H o gg	335	1.0	_	0.5^{\sharp}	_	
$H o ZZ^*$	108	$\lesssim 0.3^{\dagger}$	_	_	0.1^{\ddagger}	
$H o \gamma\gamma$	_	< 1.0	_	_	_	
$H o Z\gamma$	2.1	1.0	_	_	0.1^{\ddagger}	

 † From $e^+e^- \rightarrow ZH$.

 ‡ For $\delta M_H=10$ MeV. Adjusted for Higgs mass precision at CLIC. $^{\flat}$ For $\delta m_b=13$ MeV, $\delta m_c=7$ MeV. (Lattice projection). $^{\sharp}$ For $\delta \alpha_s=0.0002$. (Lattice projection).

A. Freitas et al., arXiv: 1906.05379 [hep-ph]

Intrinsic TH unc in production

e.g. $e^+e^- \rightarrow ZH$

LO to NLO: 5-10%

Missing 2-loop: O(1%)

Full 2-loop should reduce uncertainty to O(0.1%)

Z width effects relevant at this level of precision?

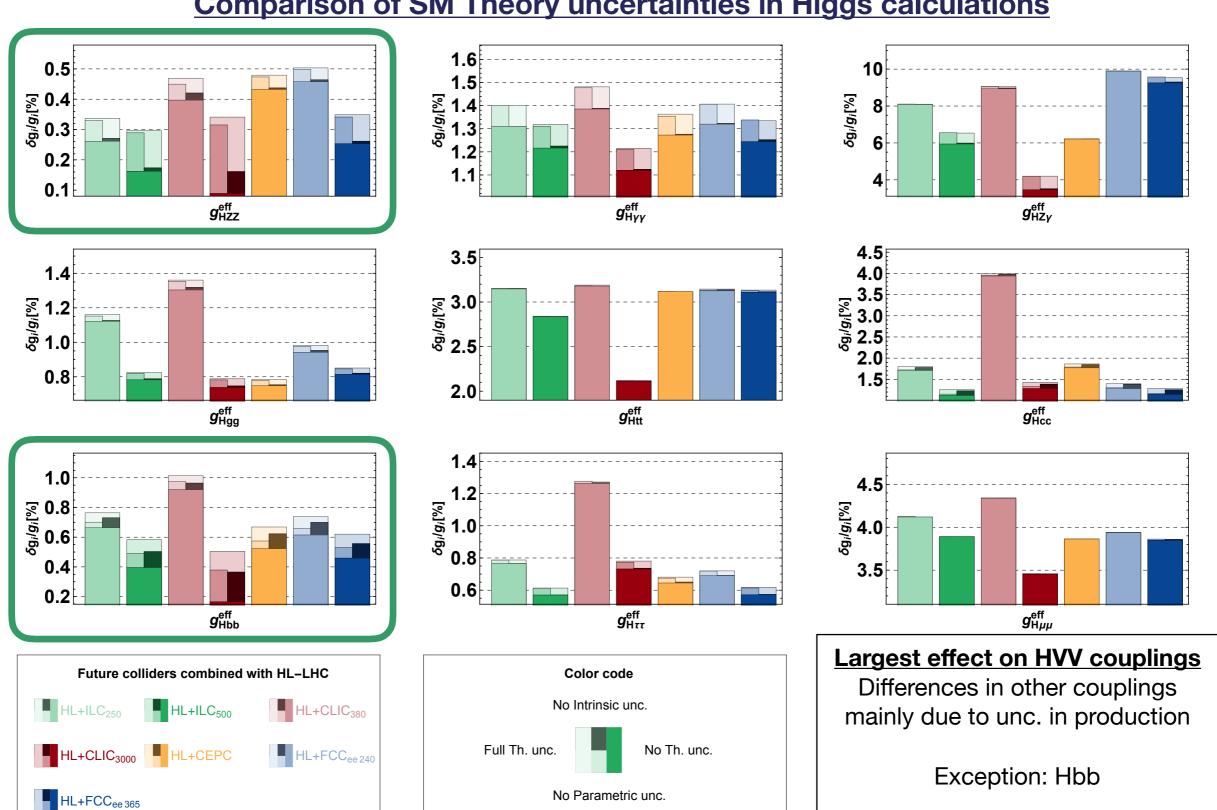
Assessment of TH uncertainty may require full 2->3 NNLO

In any case, <u>reducible</u> with necessary effort from theory side

Hence the choice of presenting main results with parametrics only

Will SM theory calculations be enough?

Comparison of SM Theory uncertainties in Higgs calculations



Will SM theory calculations be enough?

Theory requirements for EWPO

	experimental accuracy			intrinsic theory uncertainty		
	current	ILC	FCC-ee	current	current source	prospect
$\Delta M_{ m Z} [{ m MeV}]$	2.1	_	0.1			
$\Delta\Gamma_{ m Z}[{ m MeV}]$	2.3	1	0.1	0.4	$\alpha^3, \alpha^2\alpha_s, \alpha\alpha_s^2$	0.15
$\Delta \sin^2 \theta_{\rm eff}^{\ell} [10^{-5}]$	23	1.3	0.6	4.5	$\alpha^3,\alpha^2\alpha_{\rm s}$	1.5
$\Delta R_{\rm b}[10^{-5}]$	66	14	6	11	$\alpha^3,\alpha^2\alpha_{\rm s}$	5
$\Delta R_{\ell}[10^{-3}]$	25	3	1	6	$\alpha^3,\alpha^2\alpha_{\rm s}$	1.5
A. Freitas et al., arXiv: 1906.05379 [hep-ph]						

Current: Full 2-loop corrections (Not enough for future Exp. precision)

①

Prospects: Extrapolation assuming EW & QCD 3-loop corrections are known

Technically challenging but feasible (with enough support)

Still a limiting factor... Example: Reach on oblique parameters S & T

Oblique parameters: NP modifying gauge boson self-energies

$$lpha S = 4e^2 \left[\Pi_{33}^{\text{NP}}{}'(0) - \Pi_{3Q}^{\text{NP}}{}'(0) \right]$$
 $lpha T = \frac{e^2}{s_W^2 c_W^2 M_Z^2} \left[\Pi_{11}^{\text{NP}}(0) - \Pi_{33}^{\text{NP}}(0) \right]$

+ W & Y at LO in heavy NP expansion (arXiv: hep-ph/0405040) (Assumed to be ~0 here)

