

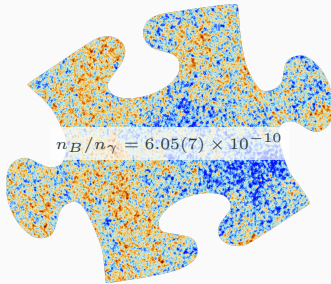
The Parameter Space of Low-scale Leptogenesis

Juraj Klarić

HKUST IAS, Hong Kong, January 20, 2020

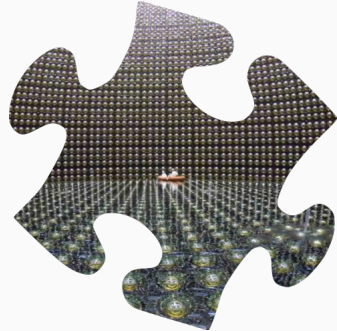
Some of the shortcomings of the standard model:

BAU baryon asymmetry of the universe



[Planck collaboration]

Neutrino masses

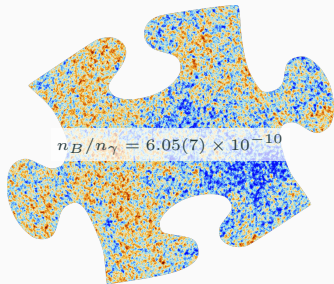


[Super-Kamiokande]

Some of the shortcomings of the standard model:

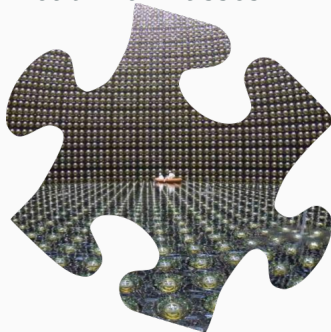
BAU

baryon asymmetry of the universe



[Planck collaboration]

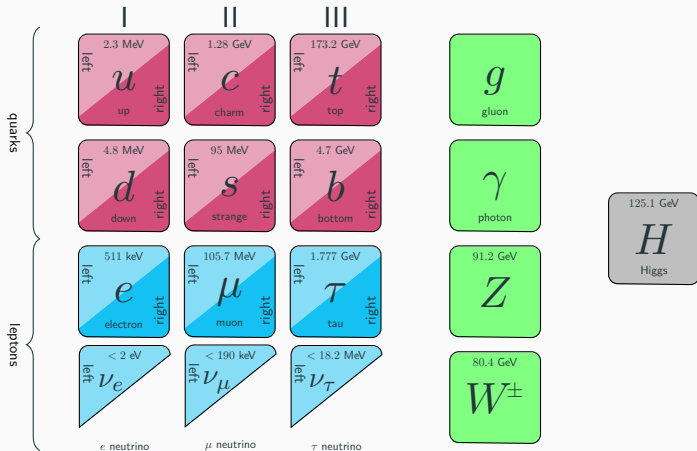
Neutrino masses



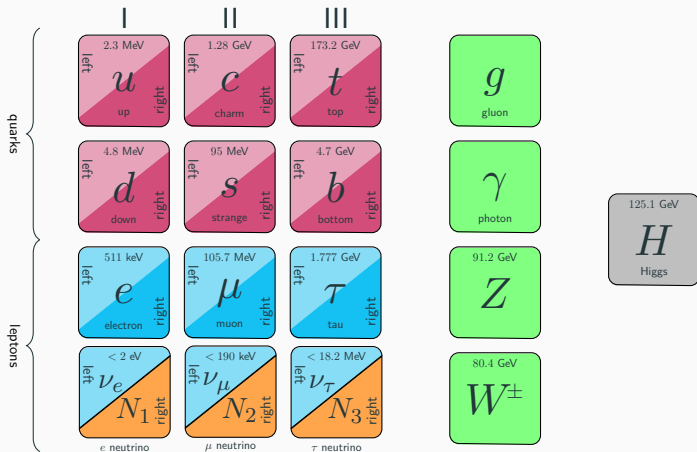
[Super-Kamiokande]

Is there a way to explain both?

Standard Model



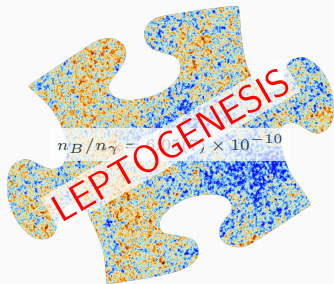
Standard Model



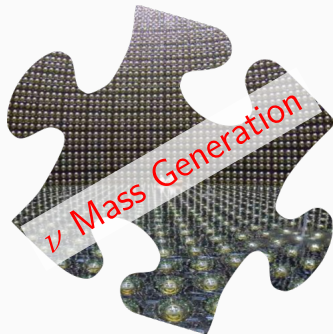
Some of the shortcomings of the standard model:

BAU

baryon asymmetry of the universe



Neutrino masses



The “Seesaw” Relation

- Dirac Mass $m_D = vF$
- Right handed neutrino (RHN) Majorana mass M_M

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu_L} & \overline{N^c} \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L^c \\ N \end{pmatrix}$$

Active neutrino masses

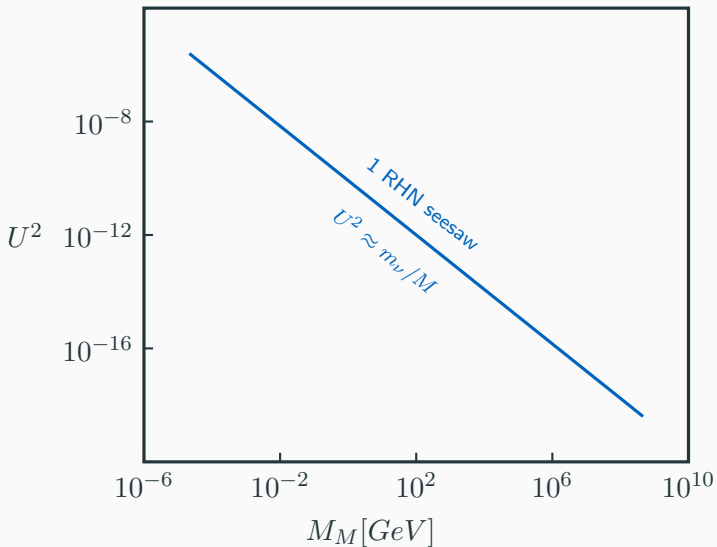
$$m_\nu = -m_D M_M^{-1} m_D^T$$

[Minkowski 1977, Yanagida 1980, . . .]

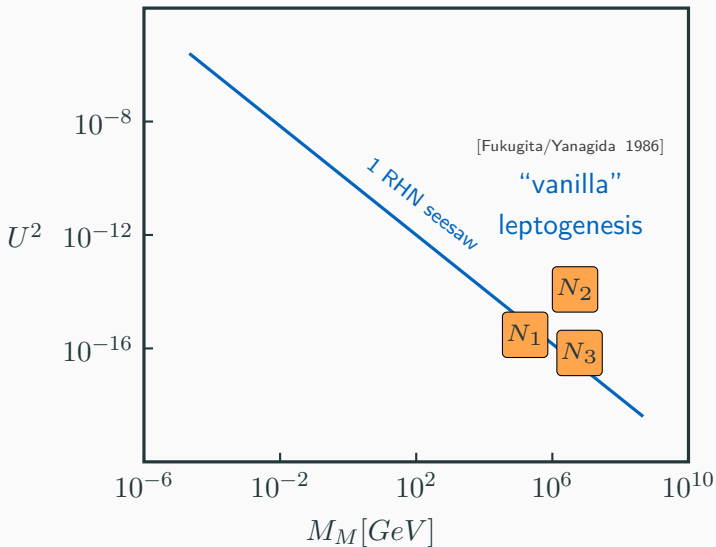
Mixing with RHN

$$U_{ai}^2 \equiv \left| \left(m_D M_M^{-1} \right)_{ai} \right|^2$$
$$U^2 = \sum_{a,i} U_{ai}^2$$

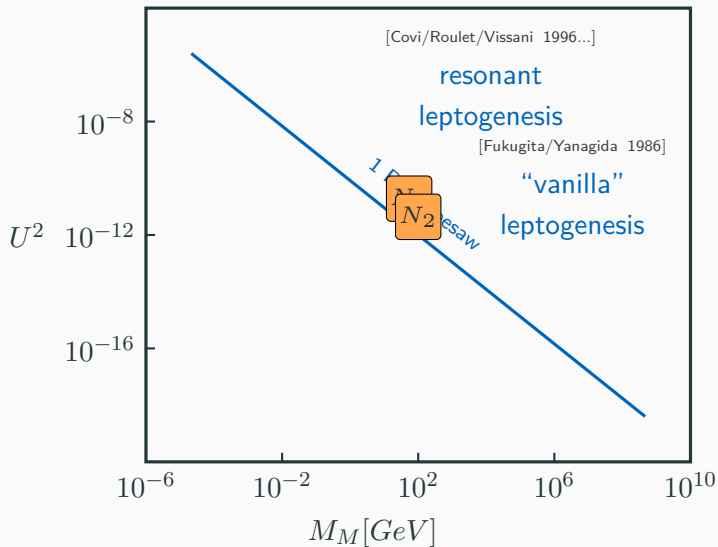
Where could the RHN be hiding?



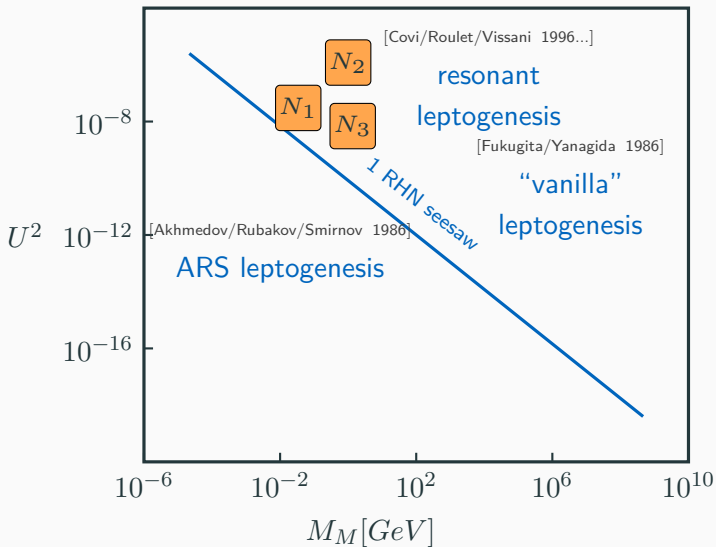
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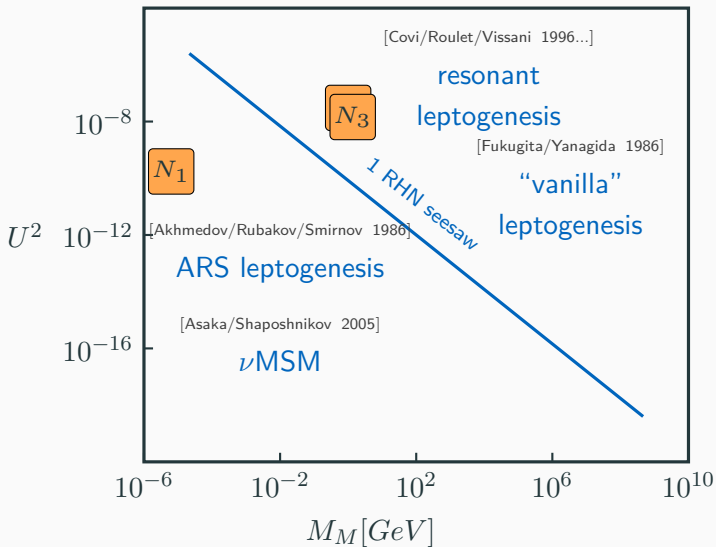
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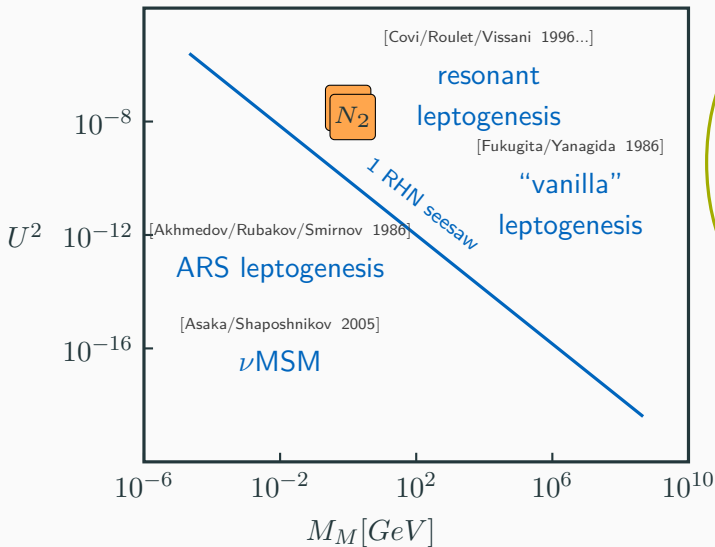
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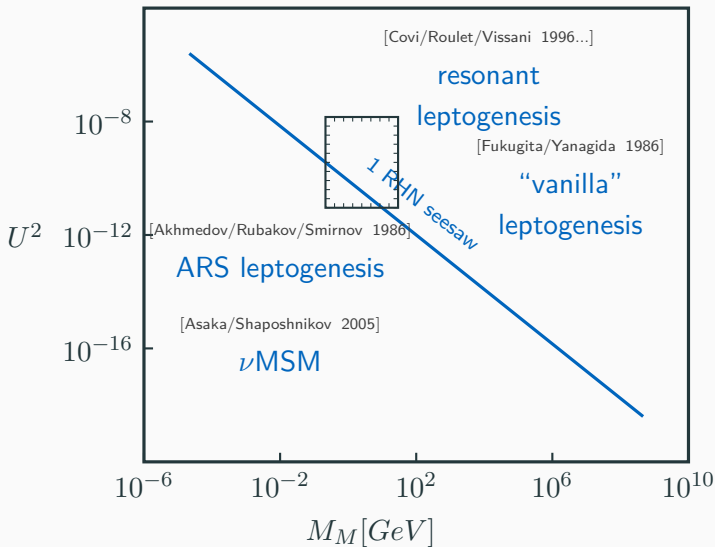
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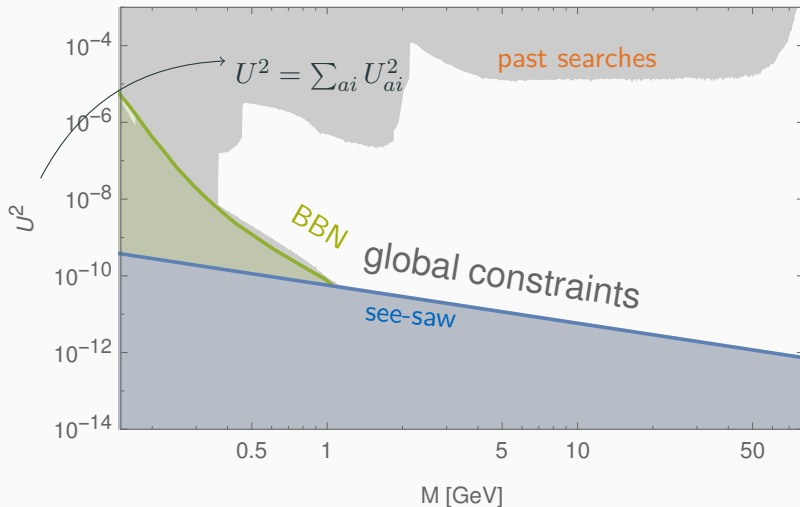
Where could the RHN be hiding?



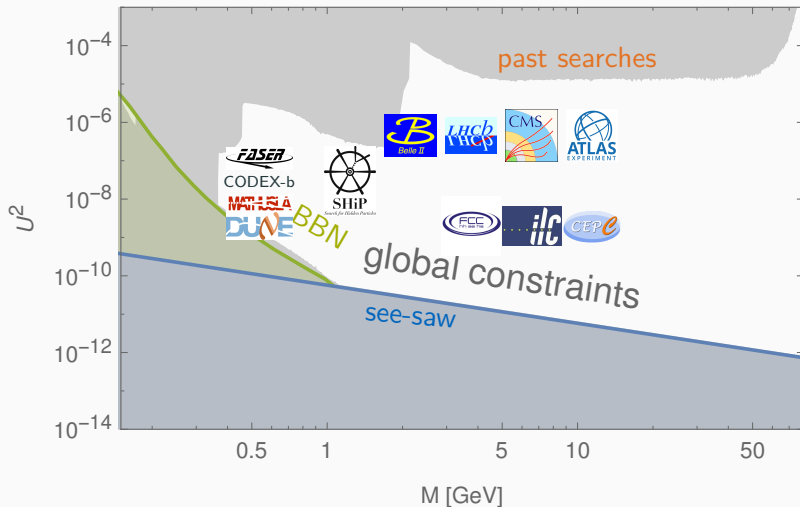
Where could the RHN be hiding?



Constraints on RHN properties: 2 RHN, Normal Ordering

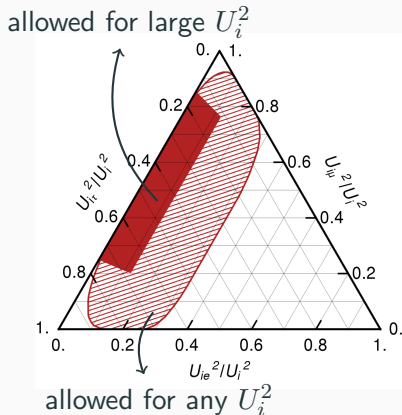


Constraints on RHN properties: 2 RHN, Normal Ordering



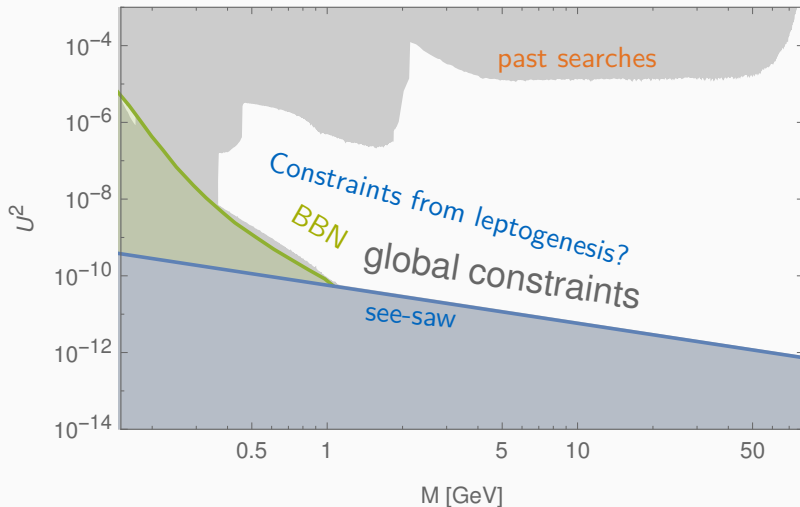
Testing the ν mass generation: Normal Ordering

Low energy parameters fixed to best fit, $\delta_{CP} \in [0, 2\pi]$ and
Majorana phase $\alpha \in [0, 4\pi]$

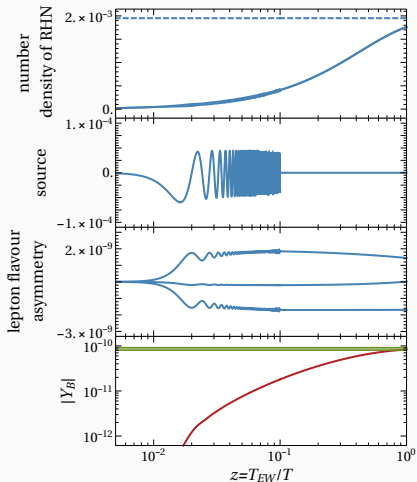
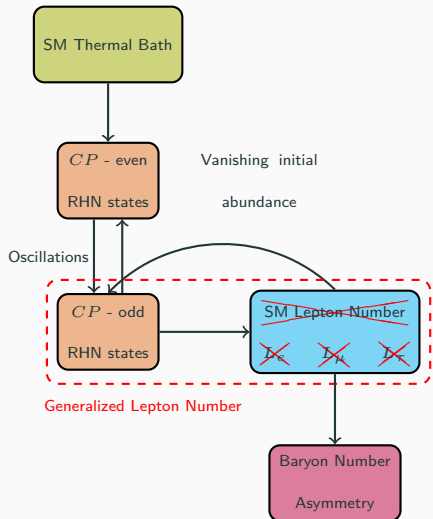


- the mixing angles U_{ai}^2 are related to the measured light neutrino parameters
- this constrains ratios U_{ai}^2/U_i^2
- constraints approach a limit for large U_i^2
- if we can measure individual U_{ai}^2 we have a consistency check for the seesaw mechanism

Constraints on RHN properties: 2 RHN, Normal Ordering



Leptogenesis via Neutrino Oscillations

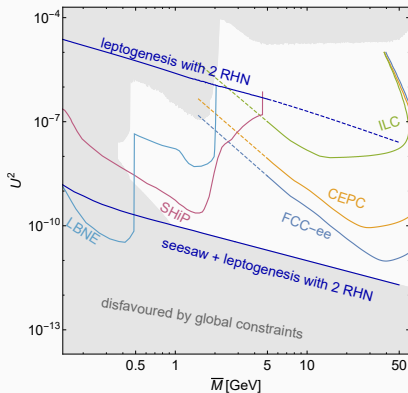


Theoretical challenges and progress:

- derivation of the **density matrix** equations from **first principles**
- **rate calculations** for RHN production at finite T
- **spectator effects**
- freezeout of the baryon number
- **analytical approximations** for different regimes
- violation of generalized lepton number
- momentum dependence
- systematic studies of the parameter space/ phenomenological implications

[A. Abada, S. Antusch, E. K. Akhmedov, G. Arcadi, T. Asaka, S. Blanchet, I. Boiarska, K. Bondarenko, A. Boyarsky, L. Canetti, A. Caputo, E. Cazzato, V. Domcke, M. Drewes, S. Eijima, O. Fischer, T. Frossard, B. Garbrecht, J. Ghiglieri, D. Gueter, T. Hambye, P. Hernandez, H. Ishida, M. Kekic, J. K., M. Laine, J. Lopez-Pavon, M. Lucente, M. Ovchinnikov, J. Racker, N. Rius, V. A. Rubakov, O. Ruchayskiy, J. Salvado, M. Shaposhnikov, B. Shuve, A. Y. Smirnov, D. Teresi, I. Timiryasov, A. Yavin...]

Leptogenesis with 2 RHN

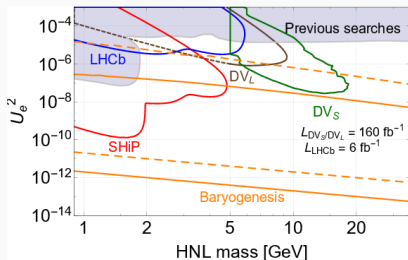


Inverted Ordering

[Drewes/Garbrecht/Gueter/JK 1609.09069]

- large mixing angle $U^2 \rightarrow$ large couplings $\rightarrow$ BAU can be erased before EWPT
- flavour asymmetric washout can help “hide” the BAU in one of the flavours
- **seesaw** allows for U_μ^2/U^2 up to a factor $\mathcal{O}(10^{-3})$ for IO

Leptogenesis with 2 RHN



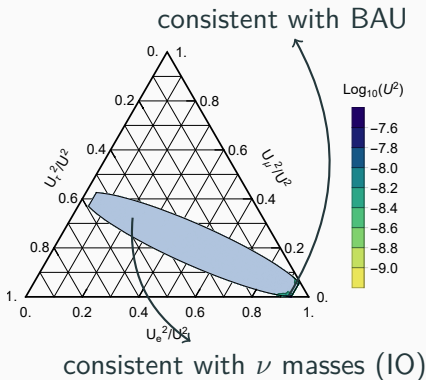
[Eijima/Shaposhnikov/Timiryasov 1808.10833]

[Boiarska/Bondarenko/Boyarsky/Eijima/

Ovchinnikov/Ruchayskiy/Timiryasov 1902.04535]

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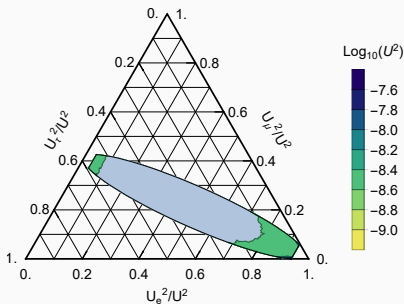


$$\bar{M} = 30 \text{ GeV}$$

$$U^2 = 4 \times 10^{-9}$$

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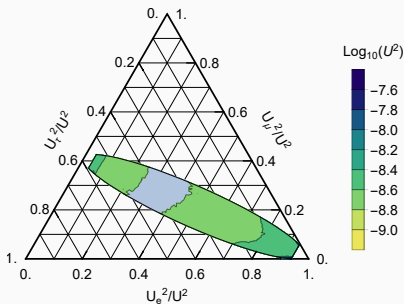


$$\bar{M} = 30 \text{ GeV}$$

$$U^2 = 2.5 \times 10^{-9}$$

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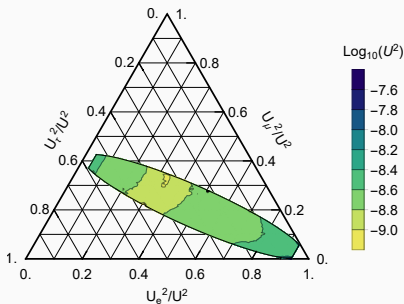
Leptogenesis with 2 RHN



$$\bar{M} = 30 \text{ GeV}$$
$$U^2 = 1.6 \times 10^{-9}$$

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Leptogenesis with 2 RHN

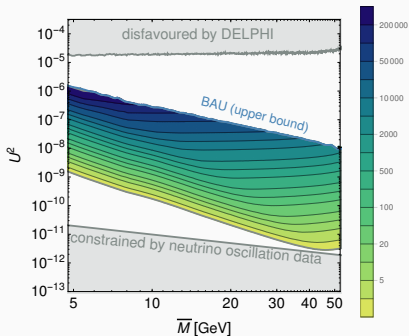


$$\bar{M} = 30 \text{ GeV}$$

$$U^2 < 1 \times 10^{-9}$$

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Measuring Flavour Ratios

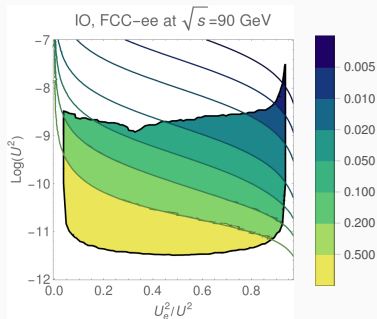


[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

1710.03744]

- large mixing angles - up to $O(10^5)$ RHN can be produced at the FCC
- use this to measure the flavour ratios
- large M - fraction of semileptonic events with flavour a in the final state proportional to U_a^2/U^2

Measuring Flavour Ratios



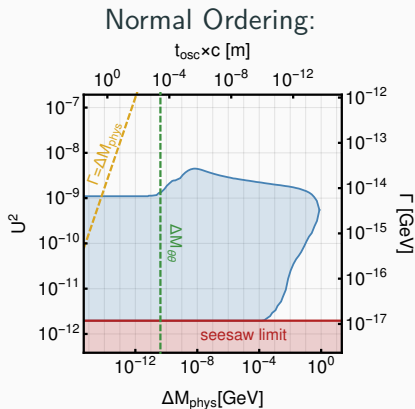
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[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

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Measuring the mass splitting



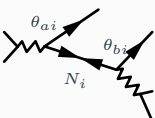
$$M = 30 \text{ GeV}$$

[Antusch/Cazzato/Drewes/Fischer/Garbrecht/Gueter/JK

1710.03744]

- large range of viable ΔM consistent with leptogenesis
- energy resolution of planned experiments -
 $\Delta M/M \sim \mathcal{O}(\text{few}\%)$
- Higgs vev contribution to RHN mass difference $\Delta M_{\theta\theta}$ practically implies lower limit on the mass splitting
- more information possible from LNV observables?

LNV and the RHN mass splitting

$$\Gamma_{X \rightarrow \ell_a \ell_b} \sim \left| \sum_i \text{diagram}_i \right|^2$$
A Feynman diagram representing a lepton number violating (LNV) decay. It shows a vertex on the left from which two lines emerge: a solid line with an arrow pointing right (representing a lepton) and a wavy line (representing a photon or Z boson). The solid line then splits into two solid lines with arrows, labeled with angles θ_{ai} and θ_{bi} . The wavy line connects to a vertex on the right, from which a solid line with an arrow (representing a neutrino N_i) and another wavy line emerge. The entire diagram is enclosed in large vertical bars, with a summation symbol $\sum_i$ to the left of the bars and a superscript 2 to the right of the bars.

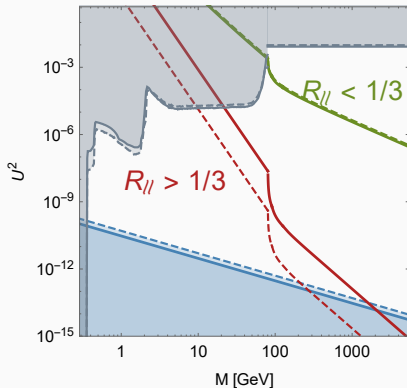
**Same-to-opposite sign
dilepton ratio**

$$R_{\ell\ell} = \frac{\Delta M_{\text{phys}}^2}{\Delta M_{\text{phys}}^2 + \Gamma_N^2}$$

[Anamiati/Hirsch/Nardi 1607.05641]

- ratio between RHN lifetime and mass splitting determines the number of LNV decays
- non-vanishing neutrino masses imply a lower limit on ΔM_{phys}
- radiative stability of ν masses implies an upper limit on ΔM_{phys}
- $R_{\ell\ell}$ can lead to information on the mass splitting

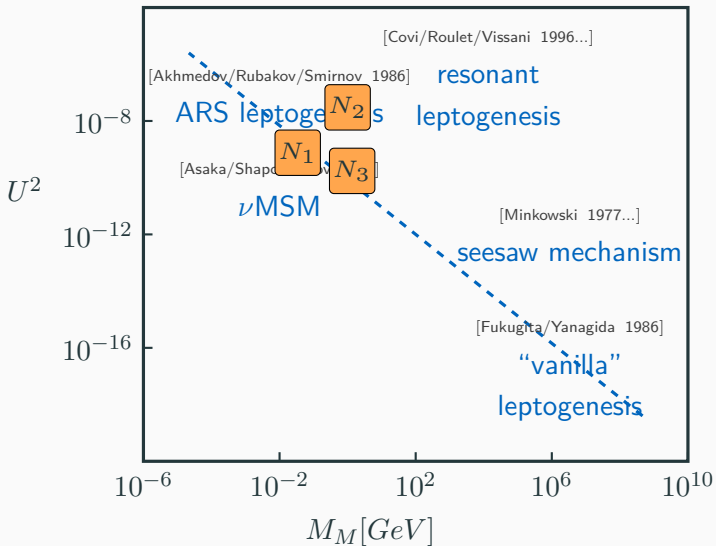
LNv and the RHN mass splitting



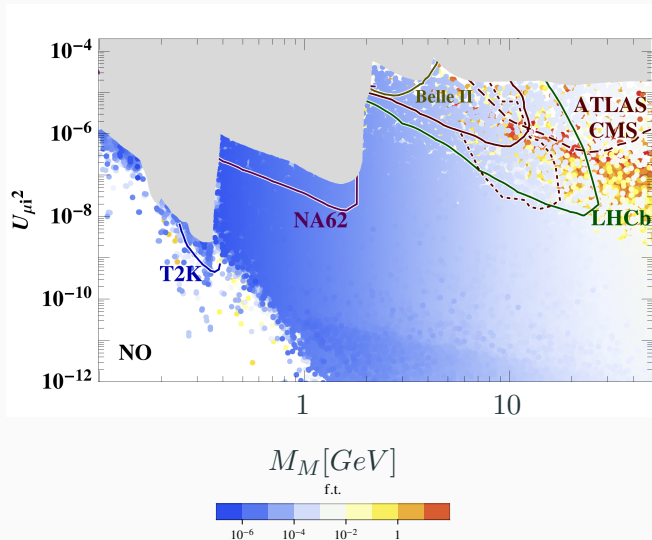
[Drewes/Klose/JK 1907.13034]

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Leptogenesis with 3 RHN



Results: Leptogenesis with 3 RHN (Normal Ordering)



How is $3 \neq 2$?: Leptogenesis

- asymmetry can be generated even without washout

[Akhmedov/Rubakov/Smirnov hep-ph/9803255]

more CP phases than in the case with two RHN

- large hierarchy in the washout is possible

[Canetti/Drewes/Garbrecht 1404.7144]

- level crossing between the heavy neutrinos
new effect!

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Sakharov II: CP

- large hierarchy in the washout is possible

[Canetti/Drewes/Garbrecht 1404.7144]

Sakharov III: non-equilibrium

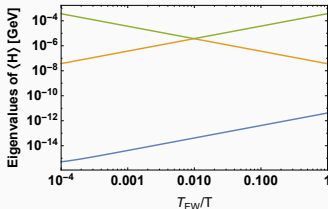
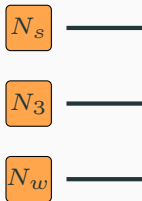
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Sakharov II: CP

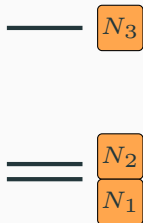
Enhancement due to level crossing

- in the $B - L$ symmetric limit two heavy neutrinos form a pseudo-Dirac pair
- the “3rd” heavy neutrino can be heavier than the pseudo-Dirac pair
- for $T \gg T_{EW}$, the pseudo-Dirac pair also has a thermal mass

$$T \gg T_{EW}$$



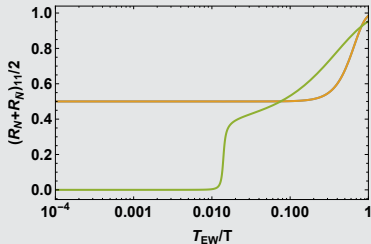
$$T \ll T_{EW}$$



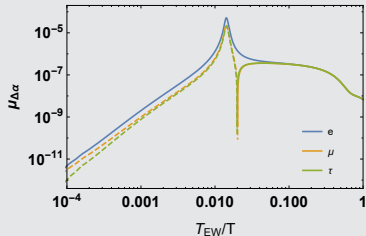
[Abada/Arcadi/Domcke/Drewes/JK/Lucente 1810.12463]

Enhancement due to level crossing

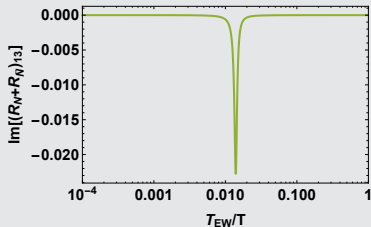
Heavy Neutrino Densities



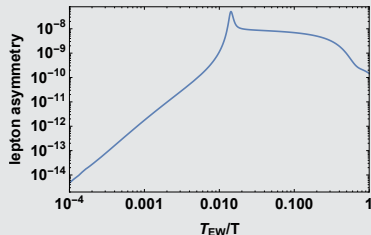
Lepton flavour asymmetries



Heavy Neutrino correlations



Lepton number asymmetry



Conclusions

- GeV-scale RHN are a viable solution to the questions of the BAU and neutrino masses
- the minimal model with 2 RHN imposes strong constraints on the properties of the RHN
- 3 RHN - larger parameter space, weaker predictions
- measurements of the branching ratios $N \rightarrow Xa$ are complementary with the low-energy neutrino oscillation experiments
- leptogenesis is within reach of existing experiments and future lepton colliders
- leptogenesis with large mixing angles implies strong constraints on the RHN branching ratios
- full testability is in principle possible, difficult in practice

Backup

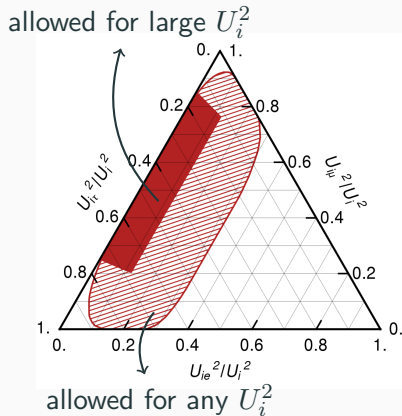
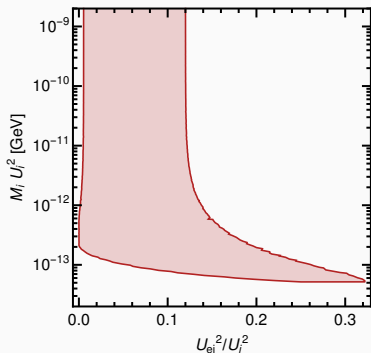
Fine tuning

- if present, symmetries are manifest to all orders in p.t.
- in the case of a large B-L breaking, radiative corrections can cause large neutrino masses
- we can use the size of radiative corrections to the light neutrino masses to quantify tuning

Fine Tuning

$$f.t.(m_\nu) = \sqrt{\sum_{i=1}^3 \left(\frac{m_i^{\text{loop}} - m_i^{\text{tree}}}{m_i^{\text{loop}}} \right)^2}$$

Testing the ν mass generation: Normal Ordering

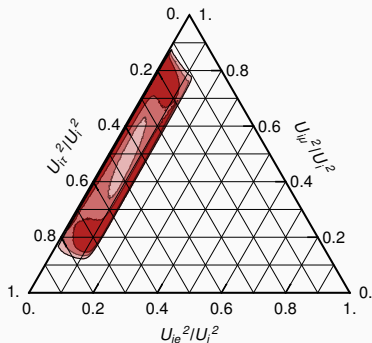
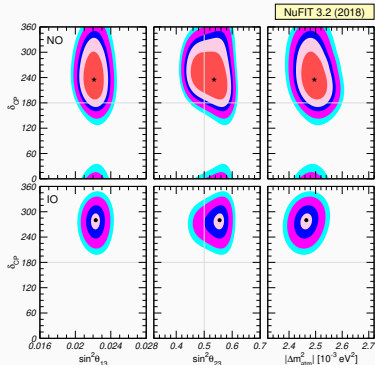


[Drewes/Hajer/JK/Lanfranchi 1801.04207]

[Drewes/Garbrecht/Gueter/JK 1609.09069] [Caputo/Hernandez/Lopez-Pavon/Salvado 1704.08721]

Testing the ν mass generation: Normal Ordering

Measurement of δ_{CP} improves the prediction:



[Esteban/Gonzalez-Garcia/Maltoni/Martinez-

Soler/Schwetz 1611.01514]

[Drewes/Hajer/JK/Lanfranchi 1801.04207]