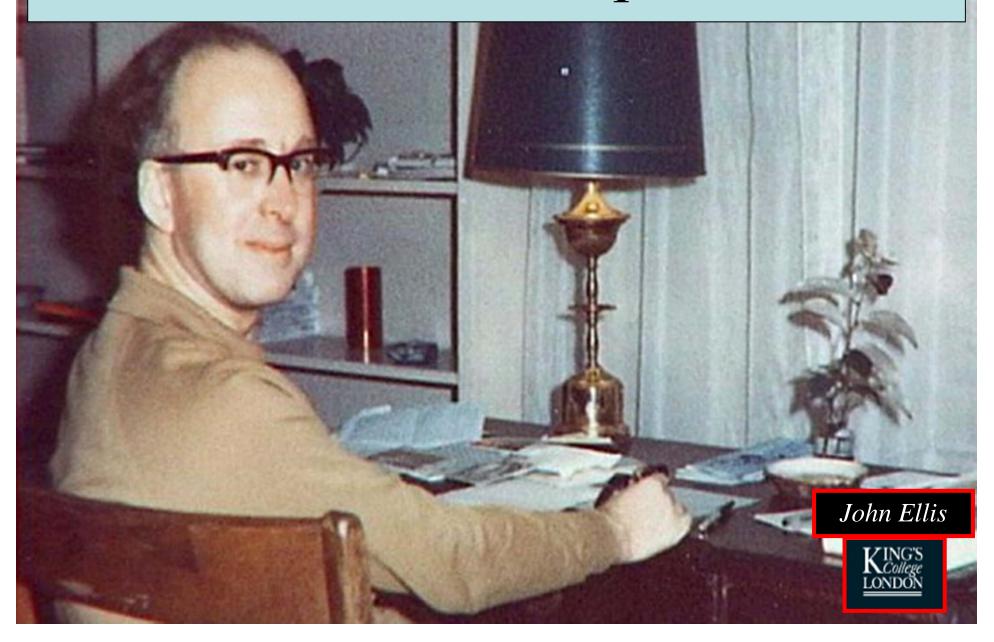
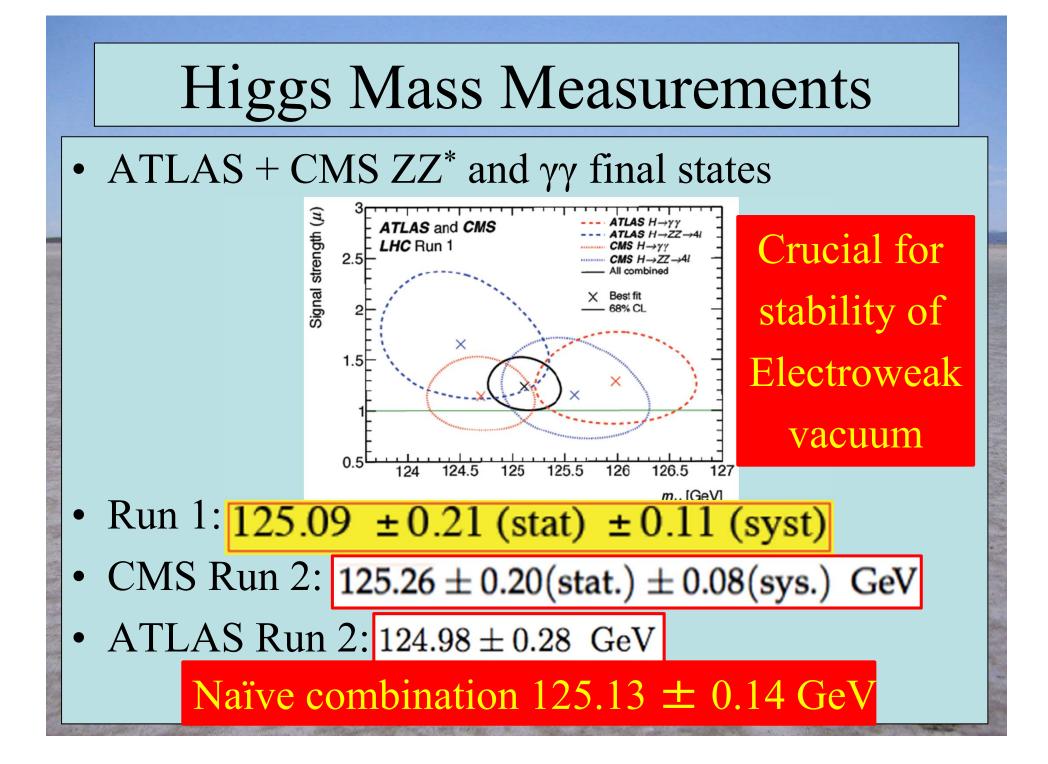
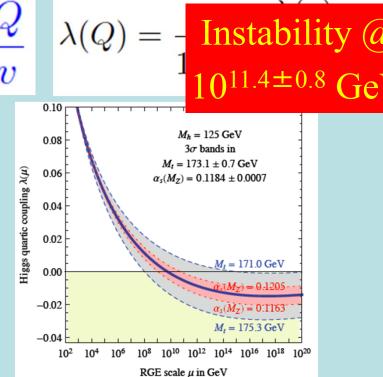
Conclusions/Perspectives





Theoretical Constraints on Higgs Mass

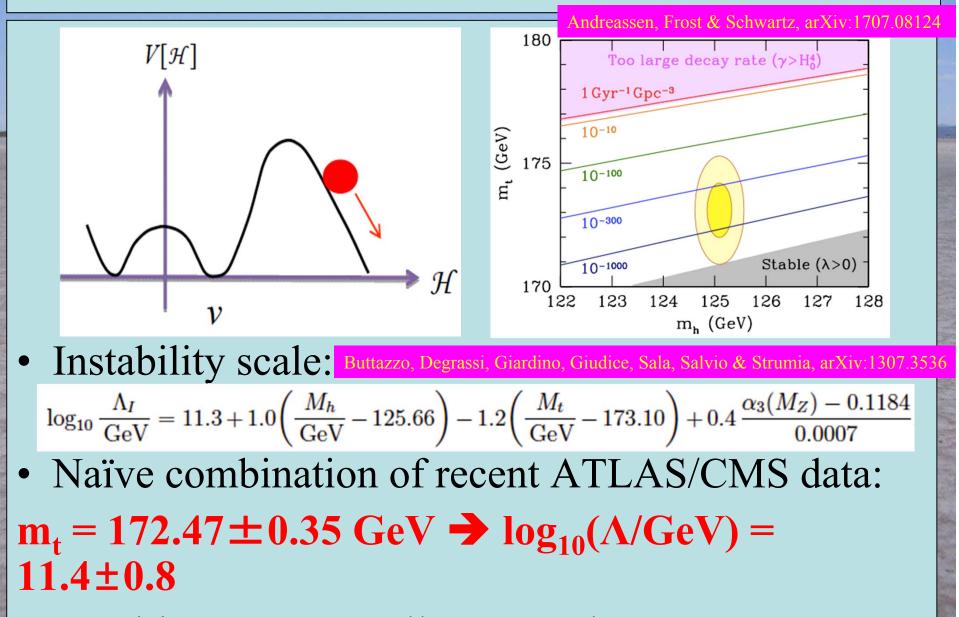
- Large $M_h \rightarrow$ large self-coupling \rightarrow blow up at $\lambda(Q) = \lambda(v) - \frac{3m_t^4}{2\pi^2 v^4} \log \frac{Q}{v} \quad \lambda(Q) = -\frac{1}{2\pi^2 v^4} \ln \frac{Q}{v}$
- Small: renormalization due to t quark drives quartic coupling < 0at some scale Λ \rightarrow vacuum unstable

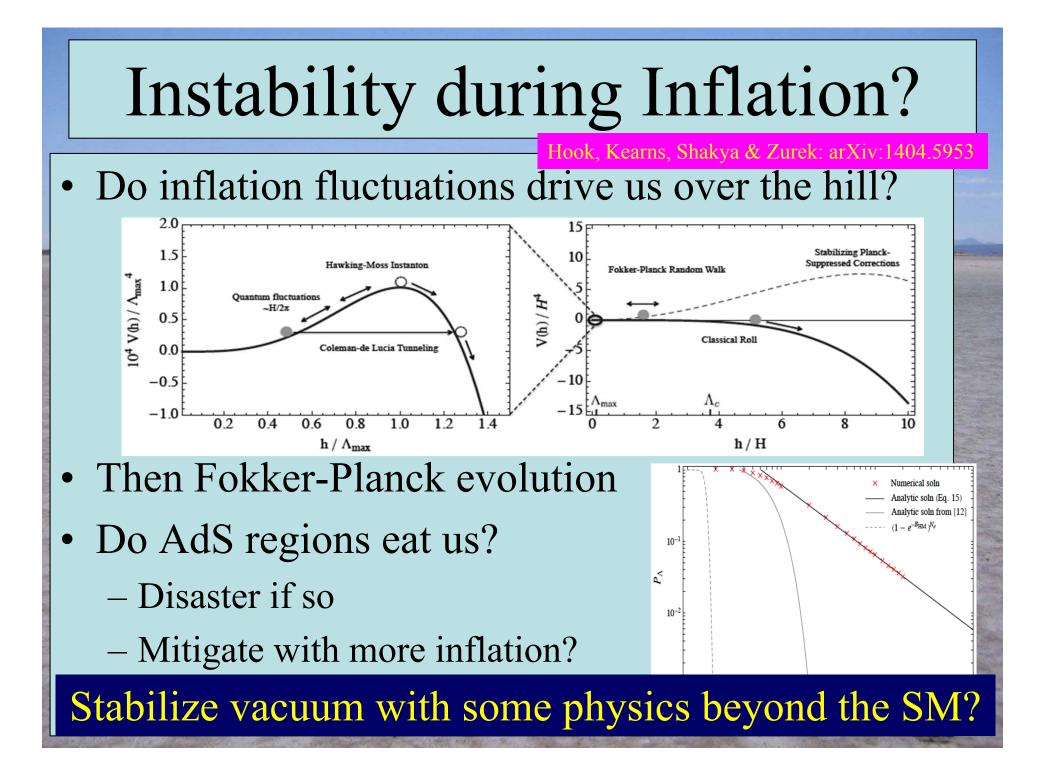


• Vacuum could be stabilized by **Supersymmetry**

Degrassi, Di Vita, Elias-Miro, Giudice, Isodori & Strumia, arXiv:1205.6497

Vacuum Instability in the Standard Model





Assuming H(125) is SM-like: Model-independent search for new physics Standard Model Effective Field Theory

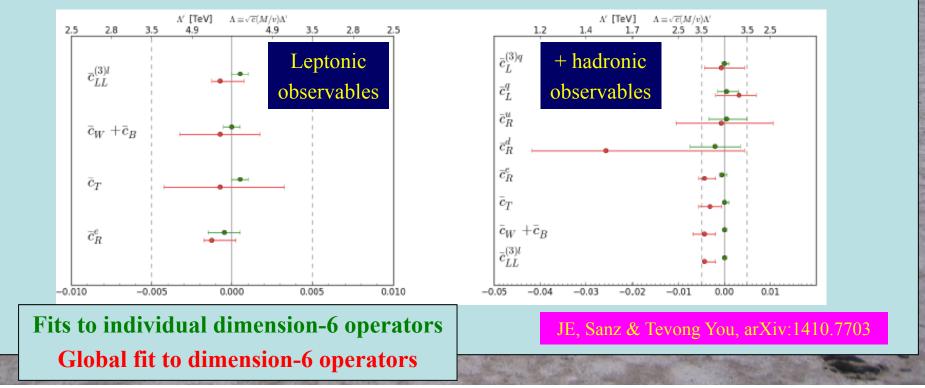
- Higher-dimensional operators as relics of higherenergy physics, e.g., dimension 6: $\mathcal{L}_{\text{eff}} = \sum \frac{f_n}{\Lambda^2} \mathcal{O}_n$
- Operators constrained by SU(2) \times U(1) symmetry:
 - $\mathcal{L} \supset \frac{\bar{c}_H}{2v^2} \partial^{\mu} [\Phi^{\dagger} \Phi] \partial_{\mu} [\Phi^{\dagger} \Phi] + \frac{g^{\prime 2} \bar{c}_{\gamma}}{m_W^2} \Phi^{\dagger} \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g_s^2 \bar{c}_g}{m_W^2} \Phi^{\dagger} \Phi G^a_{\mu\nu} G^{\mu\nu}_a$
 - $+ \frac{2ig\ \bar{c}_{HW}}{m_W^2} \left[D^{\mu} \Phi^{\dagger} T_{2k} D^{\nu} \Phi \right] W_{\mu\nu}^k + \frac{ig'\ \bar{c}_{HB}}{m_W^2} \left[D^{\mu} \Phi^{\dagger} D^{\nu} \Phi \right] B_{\mu\nu}$ $+ \frac{ig\ \bar{c}_W}{m_W^2} \left[\Phi^{\dagger} T_{2k} \overleftrightarrow{D}^{\mu} \Phi \right] D^{\nu} W_{\mu\nu}^k + \frac{ig'\ \bar{c}_B}{2m_W^2} \left[\Phi^{\dagger} \overleftrightarrow{D}^{\mu} \Phi \right] \partial^{\nu} B_{\mu\nu}$ $= \bar{a} \qquad \bar{a}$
 - $+ \quad \frac{\bar{c}_t}{v^2} y_t \Phi^{\dagger} \Phi \ \Phi^{\dagger} \cdot \bar{Q}_L t_R + \frac{\bar{c}_b}{v^2} y_b \Phi^{\dagger} \Phi \ \Phi \cdot \bar{Q}_L b_R + \frac{\bar{c}_\tau}{v^2} y_\tau \ \Phi^{\dagger} \Phi \ \Phi \cdot \bar{L}_L \tau_R$
- Constrain with precision EW, Higgs data, TGCs ...

Electroweak Precision Data

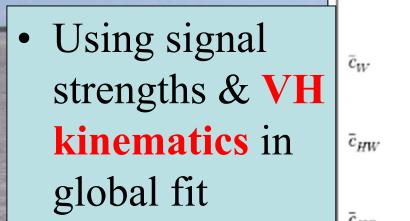
• Operators affecting oblique parameters

$$\mathcal{L}_{\text{dim-6}} \subset \frac{\overline{c}_{WB}}{m_W^2} \mathcal{O}_{WB} + \frac{\overline{c}_W}{m_W^2} \mathcal{O}_W + \frac{\overline{c}_B}{m_W^2} \mathcal{O}_B + \frac{\overline{c}_T}{v^2} \mathcal{O}_T + \frac{\overline{c}_{2W}}{m_W^2} \mathcal{O}_{2W} + \frac{\overline{c}_{2B}}{m_W^2} \mathcal{O}_{2E}$$

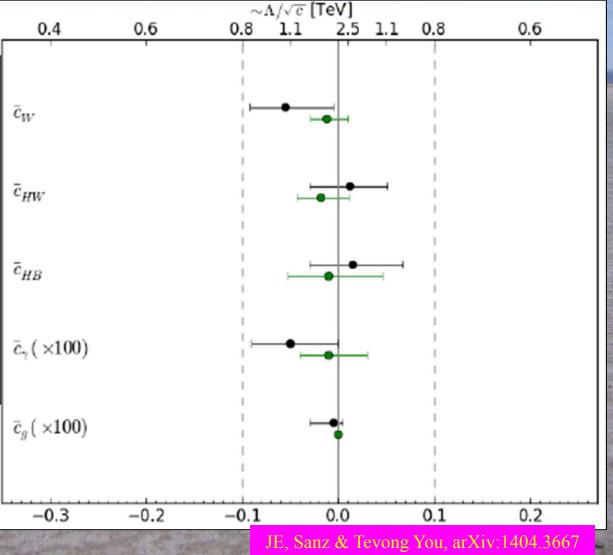
- Also other electroweak tests
- Constraints from LEP et al. data

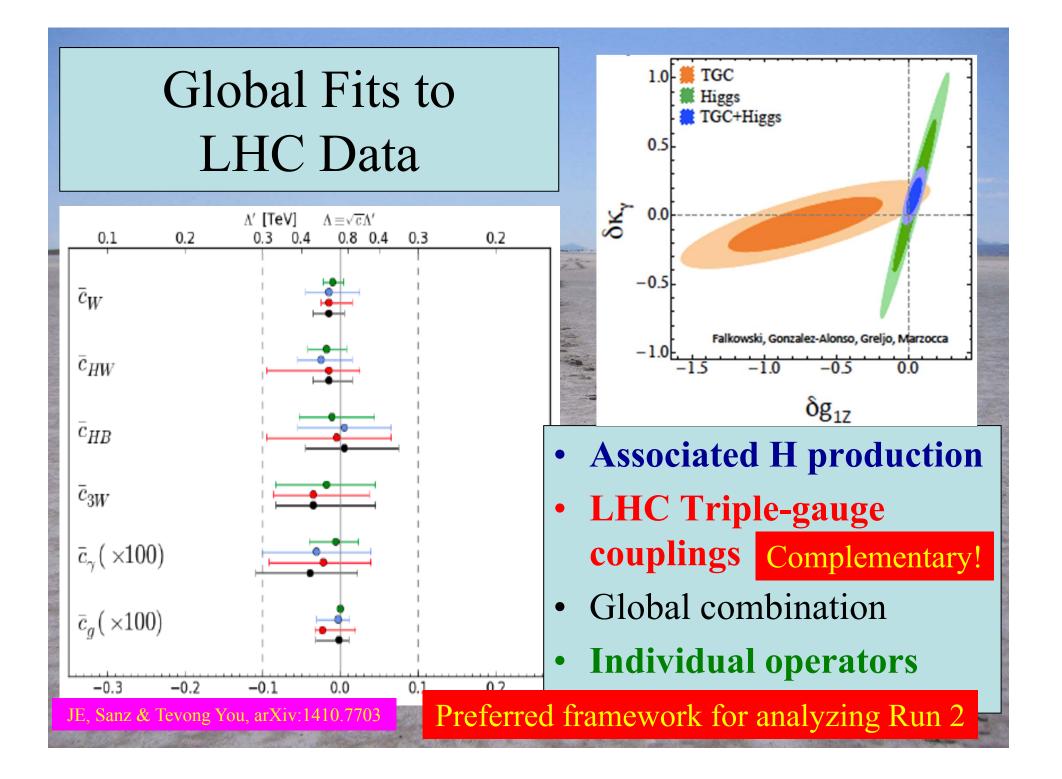


Fits to LHC Higgs Production Data

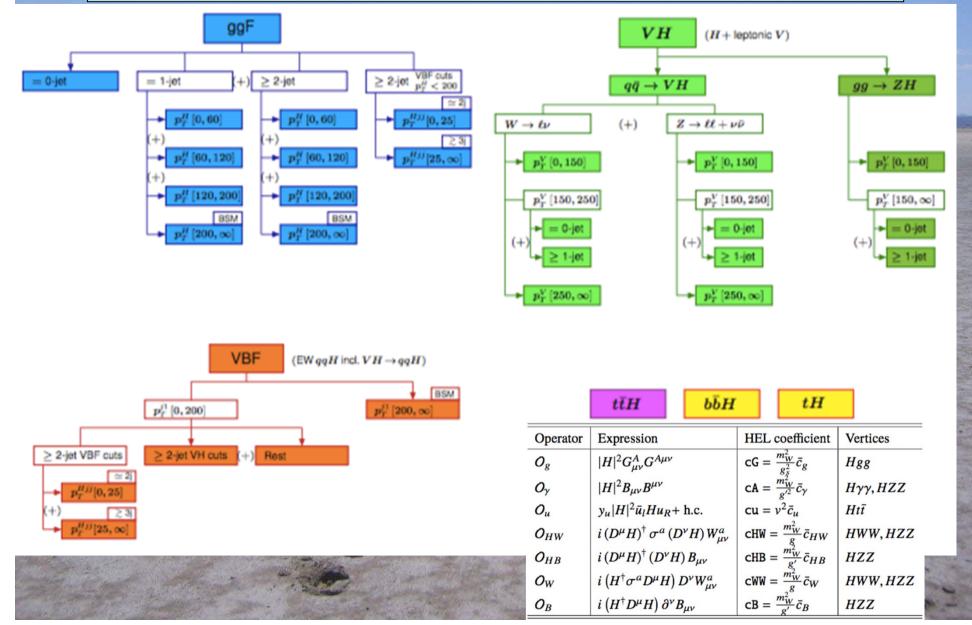


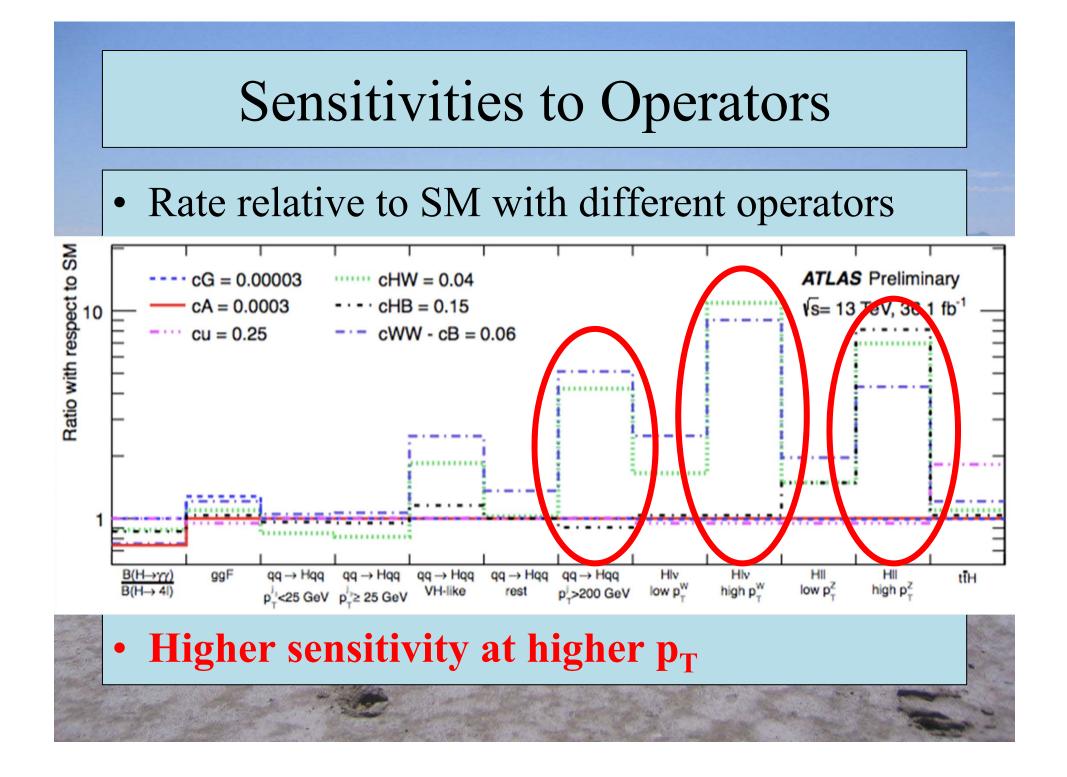
• Singleparameter fits



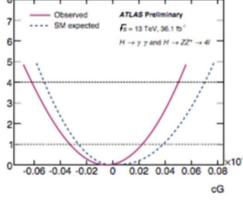


ATLAS Higgs EFT Analysis

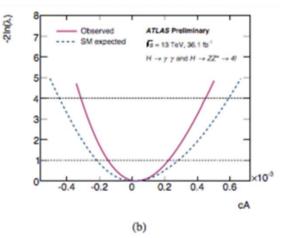




Δχ² Distributions for Higgs EFT Coefficients

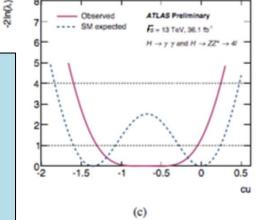


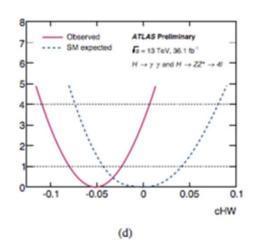
(a)



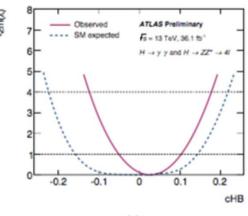
• SM (coefficient = 0) always allowed at $\Delta \chi^2 < 4$ level (< 2 σ)

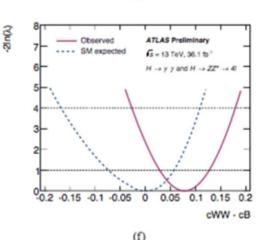




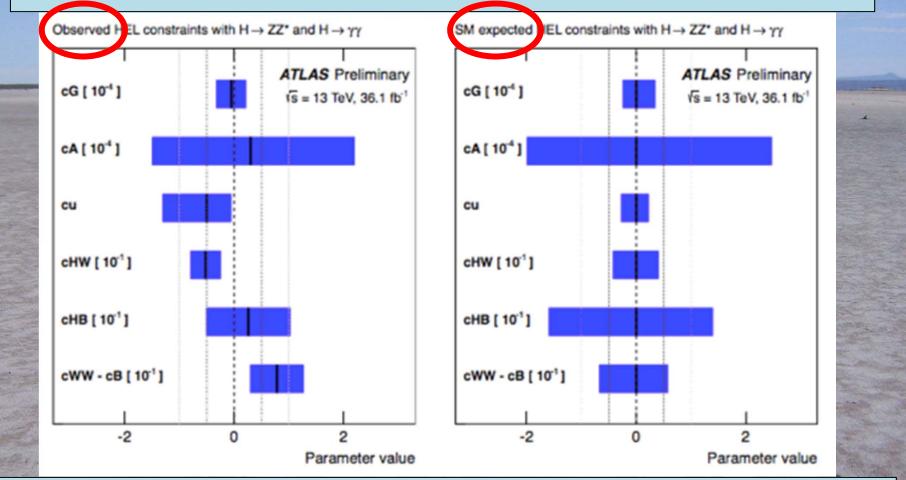


2In(J.)

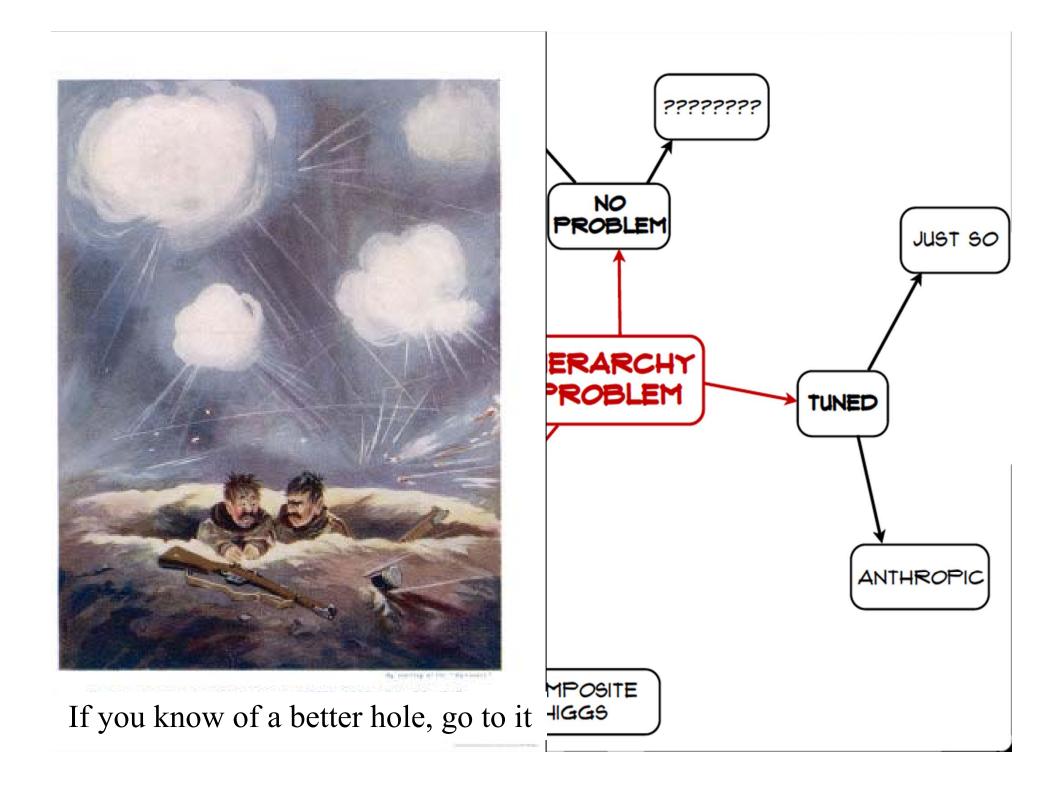




Observed vs Expected Constraints



- No significant deviations from SM
- Future: combine with precision electroweak and TGC data



What lies beyond the Standard Model?

Supersymmetry

New motivations

From LHC Run 1

- Stabilize electroweak vacuum
- Successful prediction for Higgs mass
 Should be < 130 GeV in simple models
 - Should be < 130 GeV in simple models</p>
- Successful predictions for couplings

 Should be within few % of SM values
- Naturalness, GUTs, string, ..., dark matter

Inputs to Global Fits for New Physics



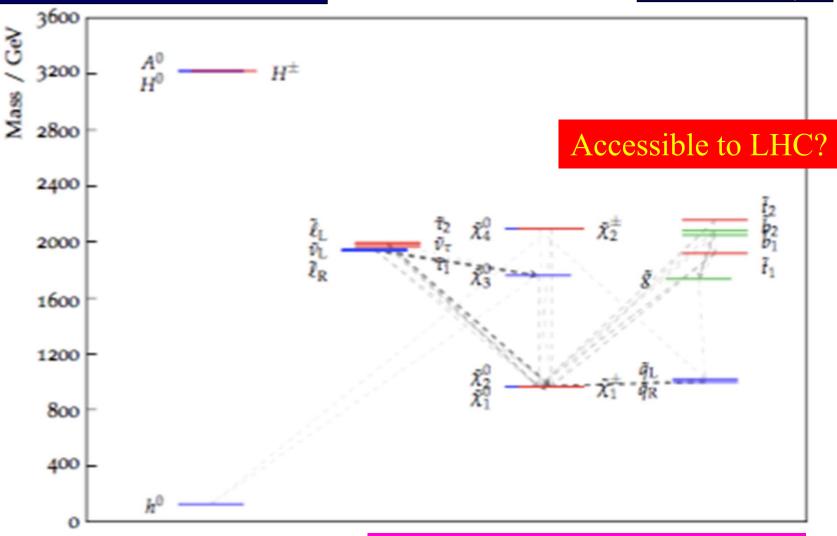
Electroweak	Observable	Source Th./Ex.		
observables	M _W [GeV]	(00) / (57 50) (22) / (27,22)	c) <u>\$0.270 + 0.012 + 0.010_{MSSM}</u>	
00501740105	$a_{\mu}^{\rm EXP} - a_{\mu}^{\rm SM}$	[59] / [60]	$(30.2 \pm 8.8 \pm 2.0_{\text{MSSM}}) \times 10^{-10}$	
Flavour	$R_{\mu\mu}$	[01-00]	2D likelihood, MFV	
Tavoui	$ au(B_s \to \mu^+ \mu^-)$	[63]	2.04 ± 0.44 (stat.) ± 0.05 (syst.) ps	
observables:	$BR_{b \rightarrow s\gamma}^{EXP/SM}$	[65]/ [66]	$0.988 \pm 0.045_{\rm EXP} \pm 0.068_{\rm TH,SM} \pm 0.050_{\rm TH,SUSY}$	
observables.	BR ^{EXP/SM}	[00,07]	$0.882 \pm 0.58_{\rm EXP} \pm 0.096_{\rm SM}$	
Interpretation	$BR_{B \to X_{a}\ell\ell}^{\text{EXP/SM}}$	[68]/ [66]	$0.966 \pm 0.278_{\rm EXP} \pm 0.037_{\rm SM}$	
	$\Delta M_{B_{\theta}}$	[01, 60] / [66]	$0.000 \pm 0.001_{\rm EXP} \pm 0.078_{\rm SM}$	
requires	$\frac{\Delta M_{B_s}^{\rm EXP/SM}}{\Delta M_{B_d}^{\rm EXP/SM}}$	[34,69] / [66]	$1.007\pm 0.004_{\rm EXP}\pm 0.116_{\rm SM}$	
-	$BR_{K \to \mu\nu}^{\text{EXP/SM}}$	[34,70] / [71]	$1.0005\pm0.0017_{\rm EXP}\pm0.0093_{\rm TH}$	
lattice inputs	$BR_{K \to \pi \nu \bar{\nu}}^{\text{EXP/SM}}$	[72]/ [73]	$2.01 \pm 1.30_{\rm EXP} \pm 0.18_{\rm SM}$	
Dark Matter	0p	[5,5,6]	Combined intermode in the $(m_{\tilde{\chi}_1^0}, \sigma_p)$ plane	
Dark Watter	σ_{e}^{SD}	[4]	Likelihood in the (m_{z^0}, σ_z^{SD}) plane	
LHC	$g ightarrow qq \chi_1^\circ, bb \chi_1^\circ, tt \chi_1^\circ$	[16,17]	Combined likelihood in the $(m_{ ilde{g}}, m_{ ilde{\chi}_1^0})$ plane	
	$ ilde{q} ightarrow q ilde{\chi}_1^0$	[16]	Likelihood in the $(m_{ar{q}}, m_{ar{\chi}_1^0})$ plane	
observables	$ ilde{b} o b ilde{\chi}_1^0$	[16]	Likelihood in the $(m_{\tilde{b}}, m_{\tilde{\chi}^0_1})$, plane	
	$ ilde{t}_1 ightarrow t ilde{\chi}_1^0, c ilde{\chi}_1^0, b ilde{\chi}_1^\pm$	[16]	Likelihood in the $(m_{\tilde{t}_1},m_{\tilde{\chi}^0_1}),$ plane	
	$ ilde{\chi}_1^\pm o u \ell^\pm ilde{\chi}_1^0, u au^\pm ilde{\chi}_1^0, W^\pm ilde{\chi}_1^0$	[18]	Likelihood in the $(m_{ ilde{\chi}_1^\pm},m_{ ilde{\chi}_1^0})$ plane	
	$ ilde{\chi}^0_2 ightarrow \ell^+ \ell^- ilde{\chi}^0_1, au^+ au^- ilde{\chi}^0_1, Z ilde{\chi}^0_1$	[18]	Likelihood in the $(m_{ ilde{\chi}^0_2},m_{ ilde{\chi}^0_1})$ plane	
	Heavy stable charged particles	[74]	Fast simulation based on [74, 75]	
	$H/A \rightarrow \tau^+ \tau^-$	[28, 29, 76, 77]	Likelihood in the $(M_A, \tan\beta)$ plane	

Best-Fit Sparticle Spectrum



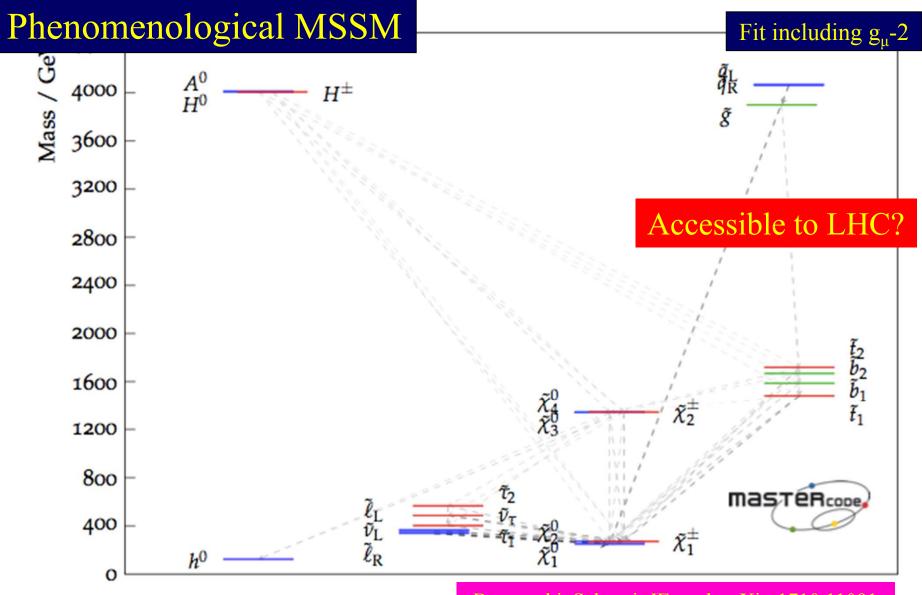
Fit excluding g_{μ} -2

Phenomenological MSSM



Best-Fit Sparticle Spectrum

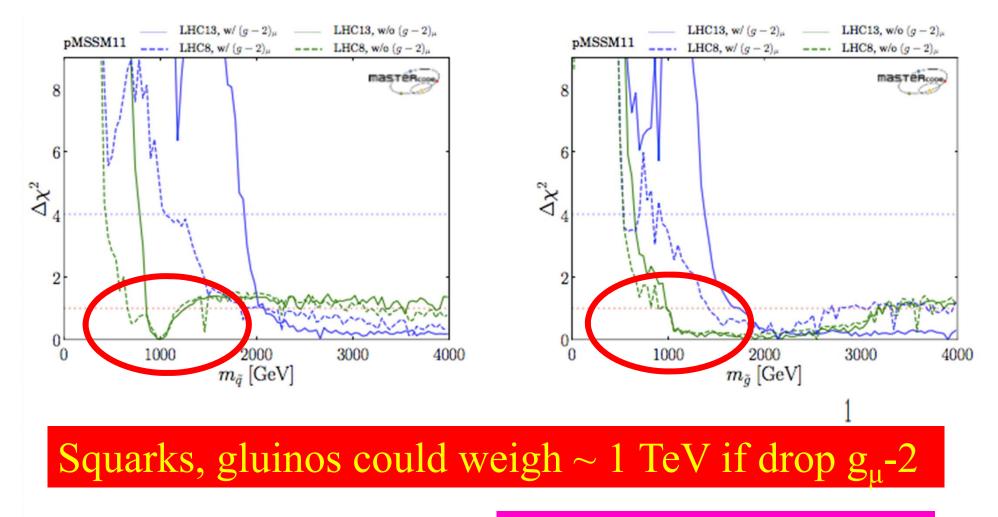




How Light can Squarks & Gluinos be?



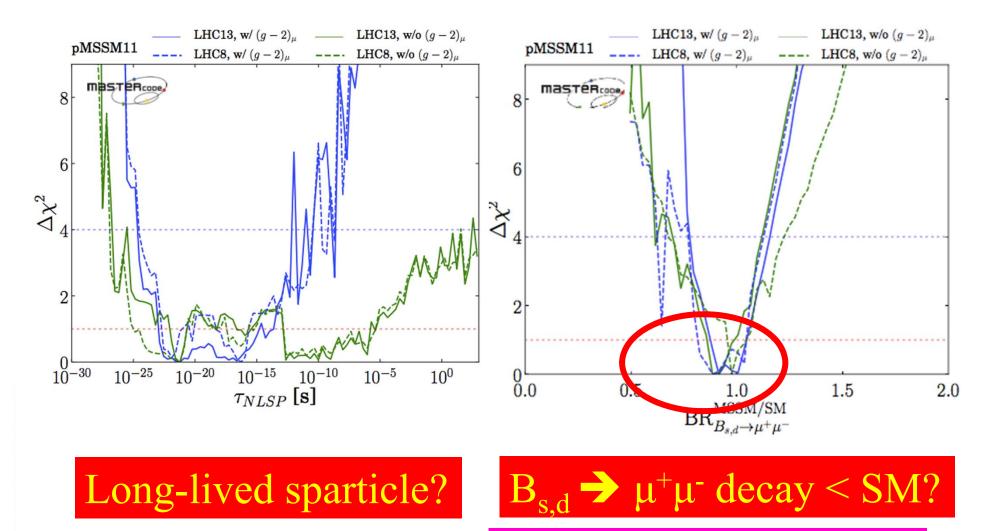
Phenomenological MSSM



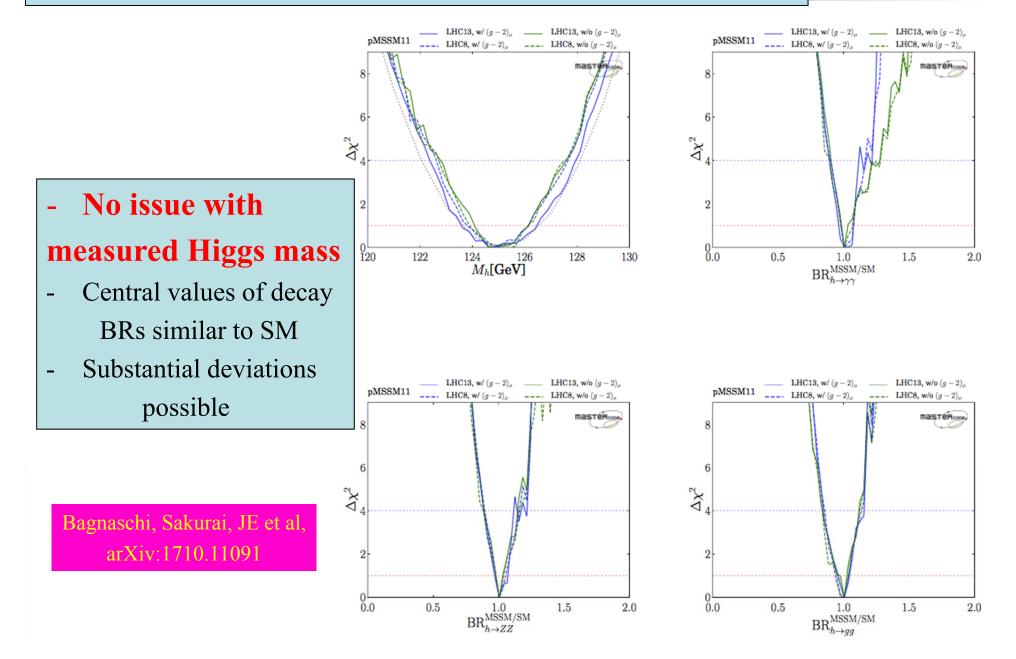
Other Possible LHC Signatures



Phenomenological MSSM



Higgs properties in the pMSSM



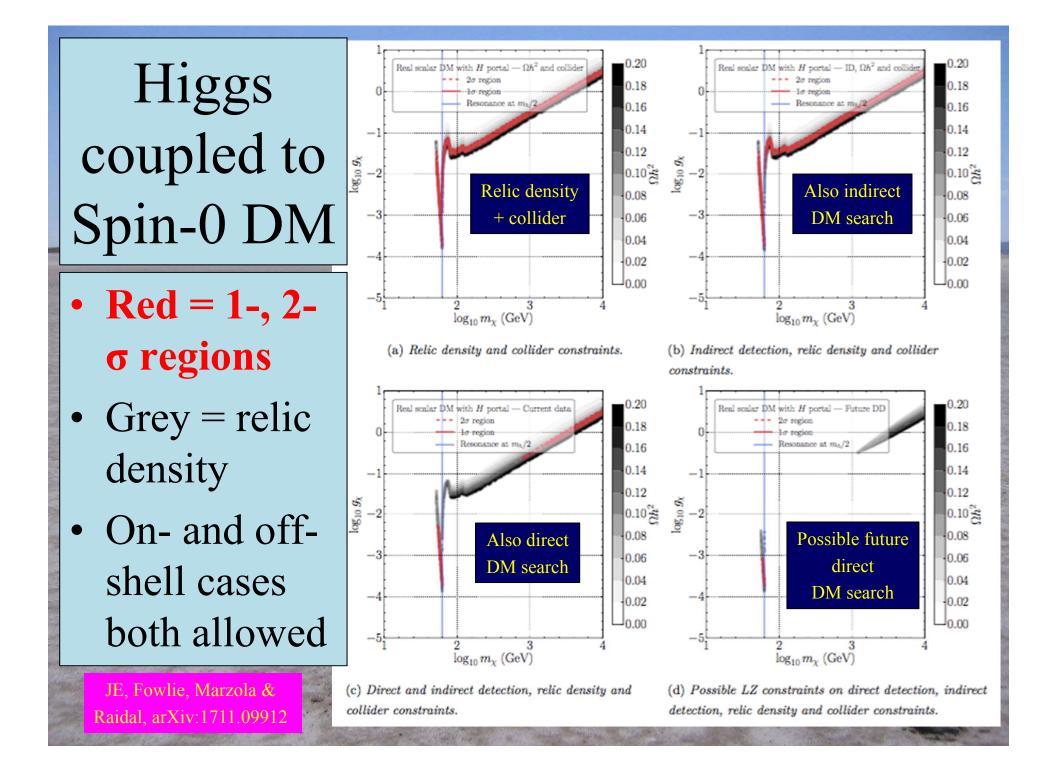
Mas TeRcope

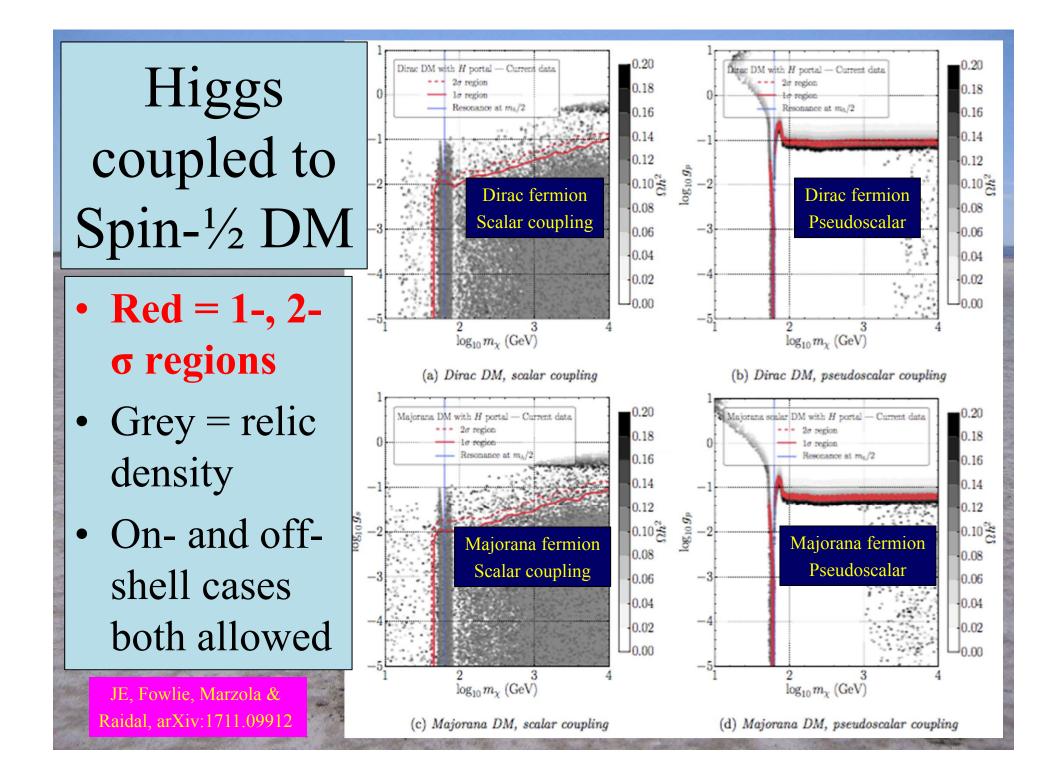


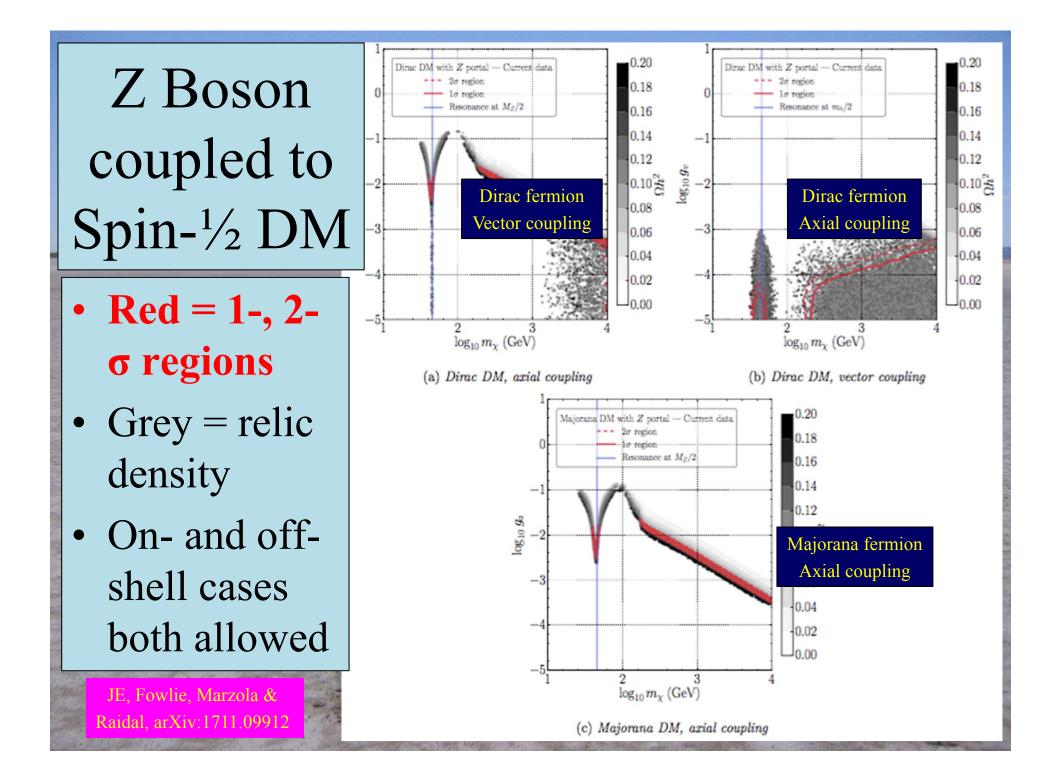
H- and Z-Portal Models are not dead yet

Consider spin-0, -1/2, -1 DM coupled to Standard Model via Higgs or Z boson All available collider, DM search constraints Bayesian & frequentist statistical analyses

JE, Fowlie, Marzola & Raidal, arXiv:1711.09912





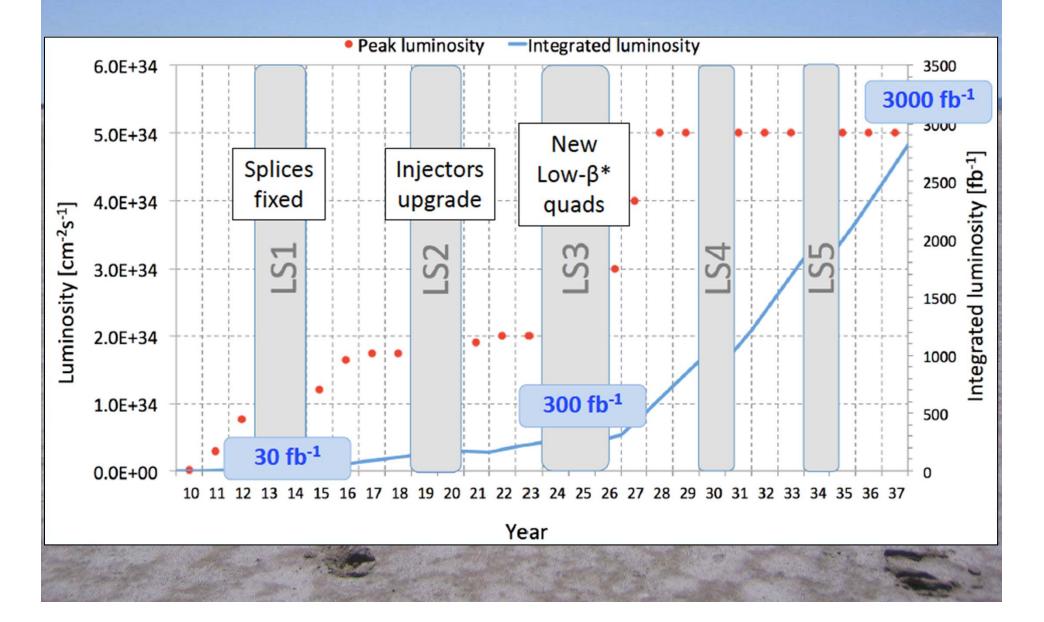


Summary of Results

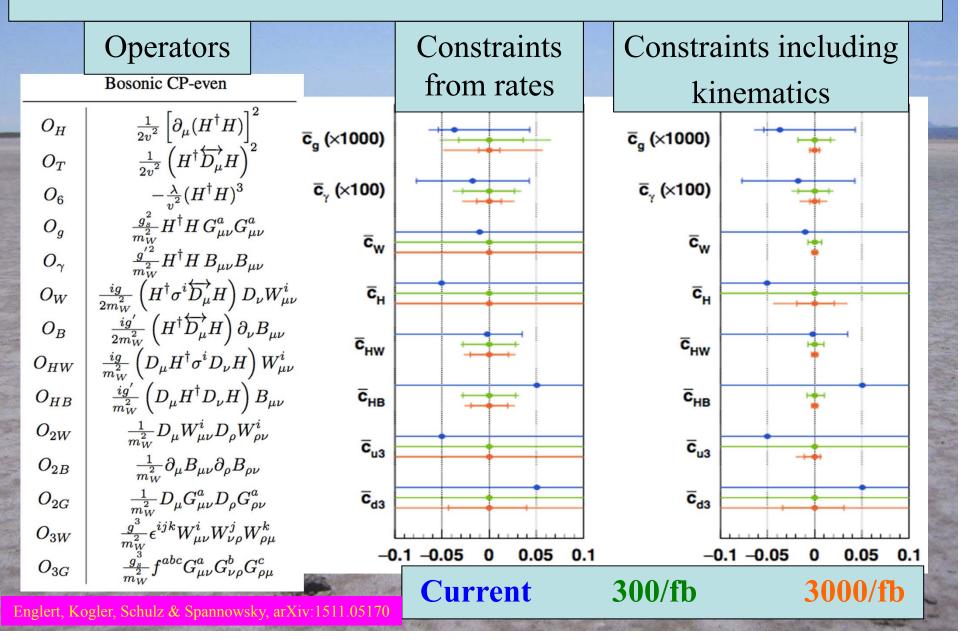
Model	Bayes factor	$\min \chi^2$	<i>p</i> -value
Real scalar h -portal	0.55	2.6	0.27
Complex scalar h -portal	0.28	2.6	0.27
Real vector h -portal	0.23	2.6	0.27
Complex vector h -portal	0.059	2.6	0.27
Majorana h -portal	0.59	2.6	0.27
Dirac <i>h</i> -portal	0.71	2.6	0.27
Scalar Z-portal Strong	$y > 10^{-14}$	55	1.4×10^{-12}
Vector Z-portal disfavour	red 8×10^{-10}	35	$2.2 imes 10^{-8}$
Majorana Z-portal	1	2.6	0.27
Dirac Z-portal	0.24	2.6	0.27

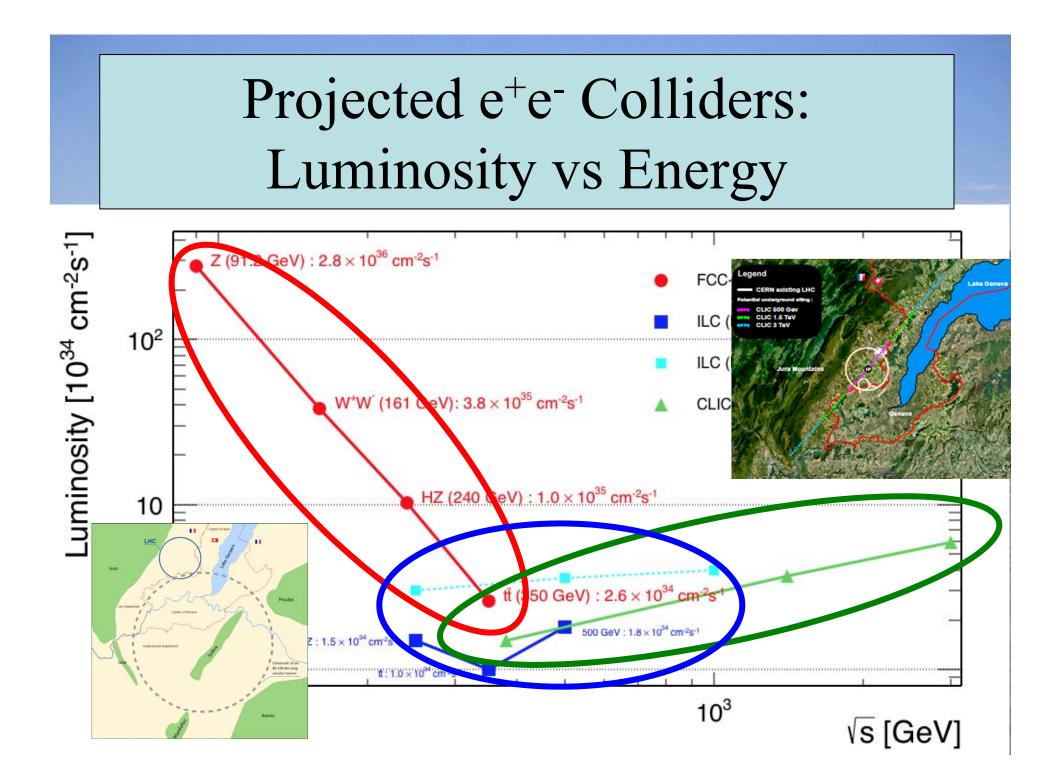
JE, Fowlie, Marzola & Raidal, arXiv:1711.09912

The LHC in Future Years



Present & Future Constraints on D=6 Operators



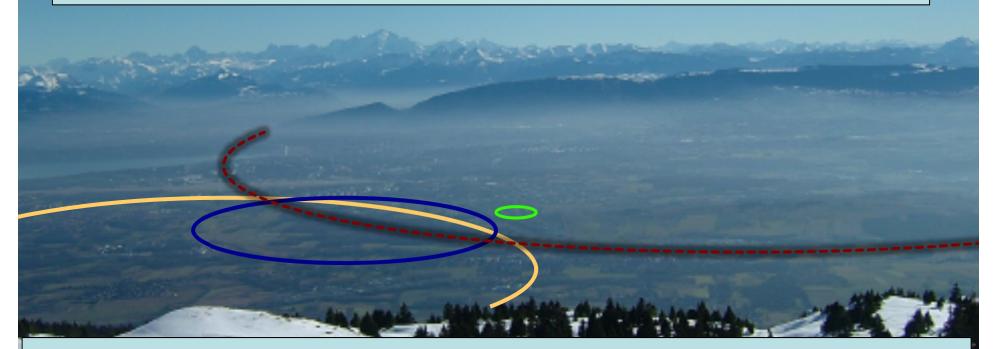




CEPC-SPPC

Preliminary Conceptual Design Report

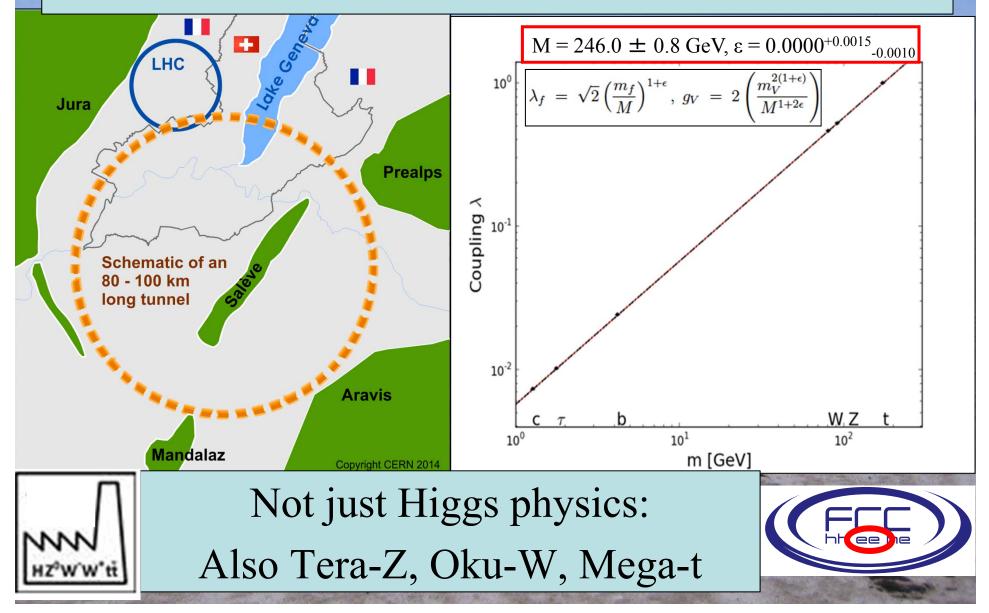
Future Circular Colliders

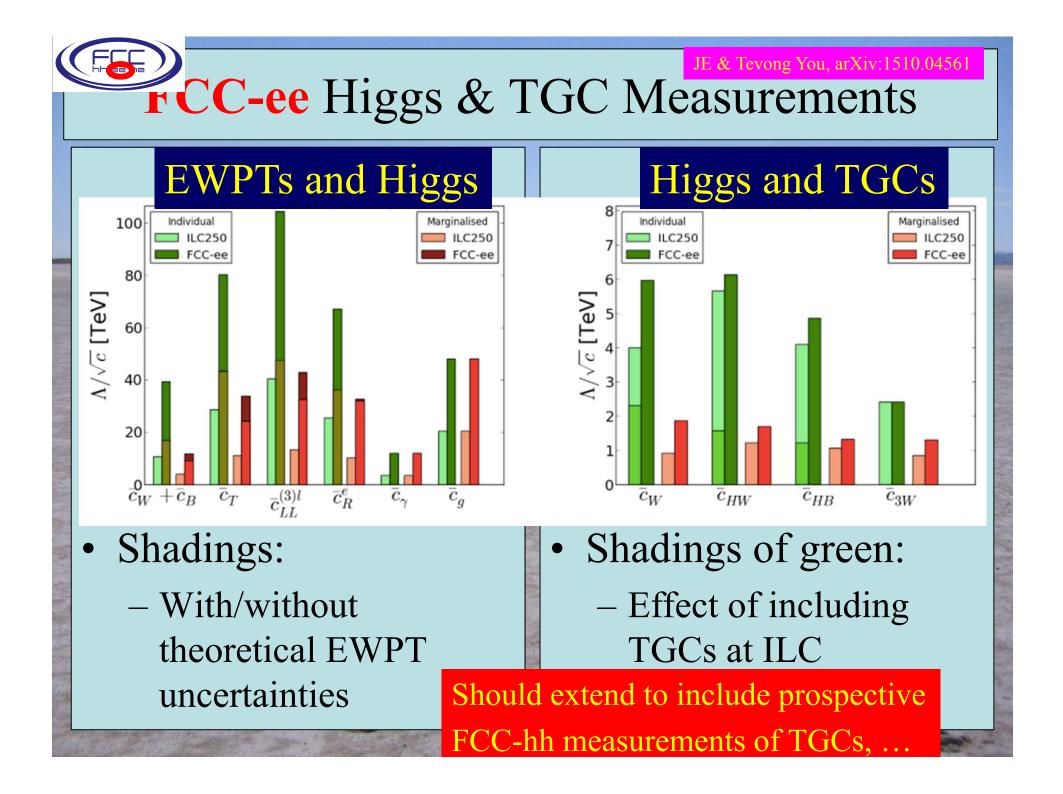


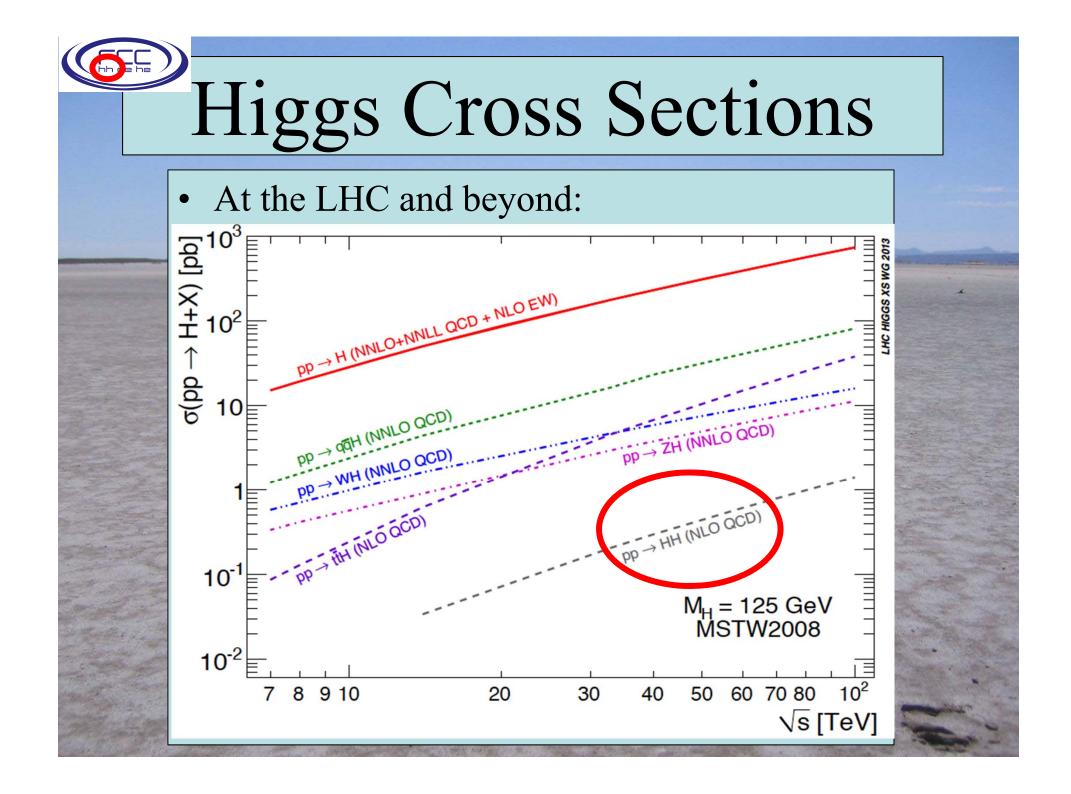
The vision:

explore 10 TeV scale directly (100 TeV pp) + indirectly (e^+e^-)

Future Circular e⁺e⁻ Collider?







Summary

- The discovery of the Higgs boson at the LHC is a big challenge for theoretical physics!
- The LHC may yet discover physics beyond the SM at ~ 13 TeV
- If it **does**, priority will be to study it
- If it does **not**, natural to focus on the Higgs
- In either case, a large circular collider offers the best prospects for future discoveries