

# Global constraints on the top-quark EFT at lepton colliders

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(DESY)

preliminary results  
based on work with Cen Zhang (IHEP),  
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*Next steps for particle physics: LHC and beyond*  
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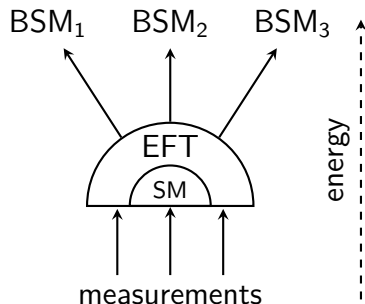


# Top EFT for lepton colliders

# The standard model effective field theory

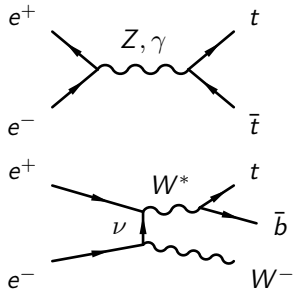
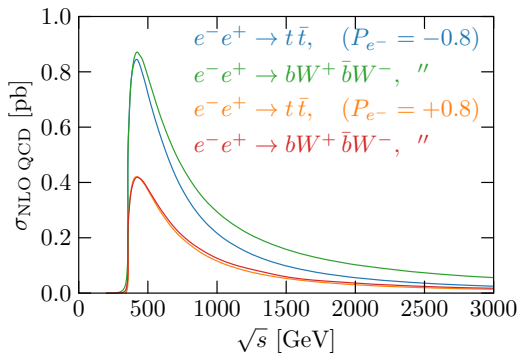
systematically parametrizes the theory space  
in direct vicinity of the SM

- ▶ in a low-energy limit,
- ▶ through a proper QFT,
- ▶ when global.



# Aiming at a global EFT analysis of $e^+e^- \rightarrow bW^+ \bar{b}W^-$

- including notably four-fermion operators
- exploiting statistically optimal observables
- allowing for NLO QCD corrections



# Up-sector SMEFT

[Grzadkowski et al '10]

Two-quark operators:

$$\mathcal{L}_{\text{EFT}} = \sum_i \frac{C_i}{\Lambda^2} O_i$$

Scalar:  $O_{u\varphi} \equiv \bar{q} u \tilde{\varphi} \varphi^\dagger \varphi,$

Vector:  $O_{\varphi q}^1 \equiv \bar{q} \gamma^\mu q \varphi^\dagger \overleftrightarrow{D}_\mu \varphi \equiv O_{\varphi q}^+ + O_{\varphi q}^V - O_{\varphi q}^A,$

$O_{\varphi q}^3 \equiv \bar{q} \gamma^\mu \tau^I q \varphi^\dagger \overleftrightarrow{D}_\mu^I \varphi \equiv O_{\varphi q}^+ - O_{\varphi q}^V + O_{\varphi q}^A$  (CC also)

$O_{\varphi u} \equiv \bar{u} \gamma^\mu u \varphi^\dagger \overleftrightarrow{D}_\mu \varphi \equiv O_{\varphi q}^V + O_{\varphi q}^A$

$O_{\varphi ud} \equiv \bar{u} \gamma^\mu d \tilde{\varphi}^\dagger \overleftrightarrow{D}_\mu \varphi,$  (CC only,  $m_b$  int.)

Tensor:  $O_{uB} \equiv \bar{q} \sigma^{\mu\nu} u \tilde{\varphi} g_Y B_{\mu\nu}, \equiv O_{uA} - \tan \theta_W O_{uZ}$

$O_{uW} \equiv \bar{q} \sigma^{\mu\nu} \tau^I u \tilde{\varphi} g_W W_{\mu\nu}^I, \equiv O_{uA} + \cotan \theta_W O_{uZ}$  (CC also)

$O_{dW} \equiv \bar{q} \sigma^{\mu\nu} \tau^I d \tilde{\varphi} g_W W_{\mu\nu}^I,$  (CC only,  $m_b$  int.)

$O_{uG} \equiv \bar{q} \sigma^{\mu\nu} T^A u \tilde{\varphi} g_s G_{\mu\nu}^A.$  (NLO only)

Two-quark–two-lepton operators:

Scalar:  $O_{1equ}^S \equiv \bar{l} e \varepsilon \bar{q} u,$  (CC also,  $m_e$  int.)

$O_{1edq} \equiv \bar{l} e \bar{d} q,$  (CC only,  $m_e$  int.)

Vector:  $O_{1q}^1 \equiv \bar{l} \gamma_\mu l \bar{q} \gamma^\mu q \equiv O_{1q}^+ + O_{1q}^V - O_{1q}^A,$

$O_{1q}^3 \equiv \bar{l} \gamma_\mu \tau^I l \bar{q} \gamma^\mu \tau^I q \equiv O_{1q}^+ - O_{1q}^V + O_{1q}^A,$  (CC also)

$O_{1u} \equiv \bar{l} \gamma_\mu l \bar{u} \gamma^\mu u \equiv O_{1q}^V + O_{1q}^A,$

$O_{eq} \equiv \bar{e} \gamma^\mu e \bar{q} \gamma_\mu q \equiv O_{eq}^V - O_{eq}^A,$

$O_{eu} \equiv \bar{e} \gamma_\mu e \bar{u} \gamma^\mu u \equiv O_{eq}^V + O_{eq}^A,$

Tensor:  $O_{1equ}^T \equiv \bar{l} \sigma_{\mu\nu} e \varepsilon \bar{q} \sigma^{\mu\nu} u.$  (CC also,  $m_e$  int.)

# Anomalous vertices

$$\begin{aligned}
 t\bar{t}\gamma : & \quad \gamma_\mu \overbrace{(F_{1V}^\gamma + \gamma_5 F_{1A}^\gamma)}^{\sim \phi} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^\gamma + i\gamma_5 F_{2A}^\gamma)}^{\sim \text{Re,Im}\{C_{uA}\}} \\
 t\bar{t}Z : & \quad \gamma_\mu \overbrace{(F_{1V}^Z + \gamma_5 F_{1A}^Z)}^{\sim C_{\varphi q}^V, C_{\varphi q}^A} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^Z + i\gamma_5 F_{2A}^Z)}^{\sim \text{Re,Im}\{C_{uZ}\}} \\
 t\bar{t}W : & \quad \gamma_\mu \overbrace{(F_{1V}^W + \gamma_5 F_{1A}^W)}^{\sim C_{\varphi q}^+ - \frac{1}{2}(C_{\varphi q}^V - C_{\varphi q}^A) \pm C_{\varphi ud}} + \frac{\sigma_{\mu\nu} i q^\nu}{2m_t} \overbrace{(F_{2V}^W + i\gamma_5 F_{2A}^W)}^{\sim s_W^2 C_{uA} + s_W c_W C_{uZ} \pm C_{dW}^*}
 \end{aligned}$$

## Insufficiencies:

- miss four-fermion operators,
- conflict with gauge invariance,  
do not allow for radiative corrections to be computed,
- complex couplings where the tree-level EFT prescribes real ones,
- hide correlations induced by gauge invariance,  
preclude the combination of measurements in various sectors

Two ideas  
for global EFT analyses

# Statistically optimal observables

[Atwood, Soni '92]

[Davier et al '93]

[Diehl, Nachtmann '94]

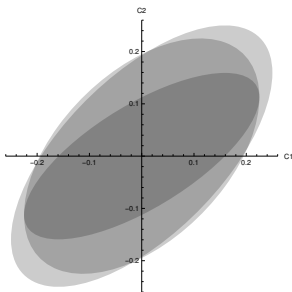
minimize the one-sigma ellipsoid in EFT parameter space

(joint efficient set of estimators,

saturating the Rao-Cramér-Fréchet bound:  $V^{-1} = I$ ,

just like MEM)

For small  $C_i$ , with a phase-space distribution  $\sigma(\Phi) = \sigma_0(\Phi) + \sum_i C_i \sigma_i(\Phi)$ ,  
the stat. opt. obs. are the average values of  $O_i(\Phi) = \sigma_i(\Phi)/\sigma_0(\Phi)$ .



e.g.  $\sigma(\phi) = 1 + \cos(\phi) + C_1 \sin(\phi) + C_2 \sin(2\phi)$

1. asymmetries:  $O_i \sim \text{sign}\{\sin(i\phi)\}$

2. moments:  $O_i \sim \sin(i\phi)$

3. statistically optimal:  $O_i \sim \frac{\sin(i\phi)}{1 + \cos \phi}$

$\Rightarrow$  area ratios 1.9 : 1.7 : 1

Previous applications in  $e^+e^- \rightarrow t\bar{t}$ , on different distributions:

[Grzadkowski, Hioki '00]

[Janot '15]

[Khiem et al '15]



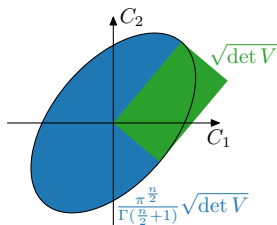
# Global determinant parameter

[GD, Grojean, Gu, Wang, '17]

In a  $n$ -dimensional Gaussian fit,  
with covariance matrix  $V$ ,

$$\text{GDP} \equiv \sqrt[n]{\det V}$$

provides a geometric average  
of the constraints strengths.



Interestingly, GDP ratios are operator-basis independent!

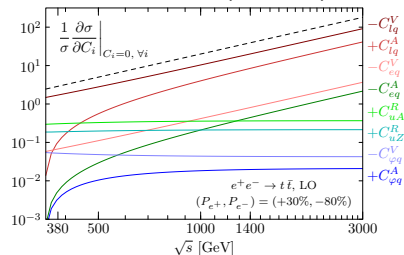
- as the volume scales linearly with coefficient normalization
- as the volume is invariant under rotations

$\implies$  conveniently assess constraint strengthening.

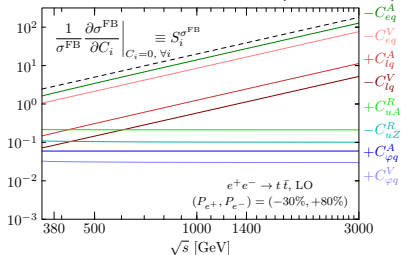
# Operator sensitivities

# $\sigma$ and $A^{\text{FB}}$ sensitivities

Total cross section (left pol.):



FB-integrated cross section (right pol.):



Few features:

- quadratic energy growth for four-fermion operators
- no growth for two-fermion operators (dipoles included)
- $p$ -wave  $\beta = \sqrt{1 - 4m_t^2/s}$  suppression of axial vectors at threshold
- enhanced sensitivity of axial vector operators in  $\sigma^{\text{FB}}$
- sensitivity sign flip for  $C_{\phi q}^V$  and  $C_{uZ}^R$  when polarization is reversed
- etc.

# Helicity amplitude decomposition in $bW^+\bar{b}W^-$

[Jacob,Wick '59]

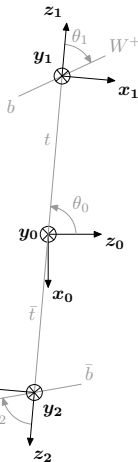
Production amplitudes:  $++ : A_1 \sim \frac{2m_t}{\sqrt{s}} V + \sqrt{s} (D - \beta\tilde{D})$   
 $-- : A_2 \sim \frac{2m_t}{\sqrt{s}} V + \sqrt{s} (D + \beta\tilde{D})$   
 $+- : A_3 \sim (V + \beta A) + 2m_t D$   
 $-+ : A_4 \sim (V - \beta A) + 2m_t D$

[Schmidt '95]

In terms of  $\Omega = \{\theta_0, \theta_1, \phi_1, \theta_2, \phi_2\}$  helicity angles:

	+3/4	$( A_3 ^2 +  A_4 ^2)$	$ a_2 ^2 +  a_4 ^2$	$ b_1 ^2 +  b_3 ^2$	$(1 + \cos^2 \theta_0)$			
	+3/4	$( A_3 ^2 -  A_4 ^2)$	$ a_2 ^2 +  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$(1 + \cos^2 \theta_0)$		$\cos \theta_2$	
	+3/4	$( A_3 ^2 -  A_4 ^2)$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 +  b_3 ^2$	$(1 + \cos^2 \theta_0)$	$\cos \theta_1$		
	+3/4	$( A_3 ^2 +  A_4 ^2)$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$(1 + \cos^2 \theta_0)$	$\cos \theta_1$		$\cos \theta_2$
	-3/2	$( A_3 ^2 -  A_4 ^2)$	$ a_2 ^2 +  a_4 ^2$	$ b_1 ^2 +  b_3 ^2$	$\cos \theta_0$			
	-3/2	$( A_3 ^2 +  A_4 ^2)$	$ a_2 ^2 +  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\cos \theta_0$		$\cos \theta_2$	
	-3/2	$( A_3 ^2 +  A_4 ^2)$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 +  b_3 ^2$	$\cos \theta_0$	$\cos \theta_1$		
	-3/2	$( A_3 ^2 -  A_4 ^2)$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\cos \theta_0$	$\cos \theta_1$		$\cos \theta_2$
	+3/2	$( A_1 ^2 +  A_2 ^2)$	$ a_2 ^2 +  a_4 ^2$	$ b_1 ^2 +  b_3 ^2$	$\sin^2 \theta_0$			
	-3/2	$( A_1 ^2 -  A_2 ^2)$	$ a_2 ^2 +  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin^2 \theta_0$		$\cos \theta_2$	
	+3/2	$( A_1 ^2 -  A_2 ^2)$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 +  b_3 ^2$	$\sin^2 \theta_0$	$\cos \theta_1$		
	-3/2	$( A_1 ^2 +  A_2 ^2)$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin^2 \theta_0$	$\cos \theta_1$		$\cos \theta_2$
	+3/2	$\sqrt{2} \operatorname{Re}\{A_1^* A_4\}$	$ a_2 ^2 -  a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\cos \phi_1$
	+3/2	$\sqrt{2} \operatorname{Re}\{A_1^* A_4\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\cos \phi_1$
	+3/2	$\sqrt{2} \operatorname{Re}\{A_2^* A_3\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\cos \phi_1$
	+3/2	$\sqrt{2} \operatorname{Re}\{A_2^* A_3\}$	$ a_2 ^2 -  a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\cos \phi_1$
	-3/2	$\sqrt{2} \operatorname{Re}\{A_2^* A_4\}$	$ a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3/2	$\sqrt{2} \operatorname{Re}\{A_2^* A_4\}$	$ a_2 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3/2	$\sqrt{2} \operatorname{Re}\{A_1^* A_3\}$	$ a_2 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3/2	$\sqrt{2} \operatorname{Re}\{A_1^* A_3\}$	$ a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\cos \phi_2$
	-3	$\operatorname{Re}\{A_1^* A_2\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin^2 \theta_0$	$\sin \theta_1$	$\sin \theta_2$	$\cos(\phi_1 + \phi_2)$
	-3/2	$\operatorname{Re}\{A_3^* A_4\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin^2 \theta_0$	$\sin \theta_1$	$\sin \theta_2$	$\cos(\phi_1 - \phi_2)$
	+3/2	$\sqrt{2} \operatorname{Im}\{A_1^* A_4\}$	$ a_2 ^2 -  a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$
	+3/2	$\sqrt{2} \operatorname{Im}\{A_1^* A_4\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$
	-3/2	$\sqrt{2} \operatorname{Im}\{A_2^* A_3\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$
	-3/2	$\sqrt{2} \operatorname{Im}\{A_2^* A_3\}$	$ a_2 ^2 -  a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$
	+3/2	$\sqrt{2} \operatorname{Im}\{A_2^* A_4\}$	$ a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
	+3/2	$\sqrt{2} \operatorname{Im}\{A_2^* A_4\}$	$ a_2 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
	-3/2	$\sqrt{2} \operatorname{Im}\{A_1^* A_3\}$	$ a_2 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
	-3/2	$\sqrt{2} \operatorname{Im}\{A_1^* A_3\}$	$ a_2 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$

$$\frac{d\sigma}{d\Omega} \propto$$



# Helicity amplitude decomposition in $bW^+\bar{b}W^-$

[Jacob,Wick '59]

Production amplitudes:  $++ : A_1 \sim \frac{2m_t}{\sqrt{s}} V + \sqrt{s} (D - \beta\tilde{D})$

[Schmidt '95]

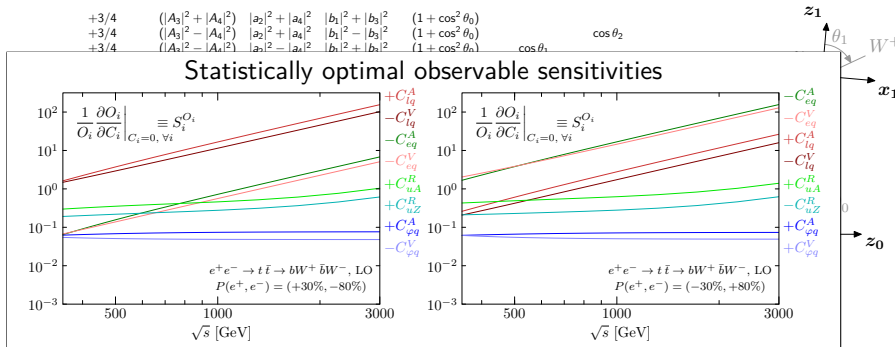
$-- : A_2 \sim \frac{2m_t}{\sqrt{s}} V + \sqrt{s} (D + \beta\tilde{D})$

$+ - : A_3 \sim (V + \beta A) + 2m_t D$

$- + : A_4 \sim (V - \beta A) + 2m_t D$

In terms of  $\Omega = \{\theta_0, \theta_1, \phi_1, \theta_2, \phi_2\}$  helicity angles:

$$\begin{array}{ccccccc} +3/4 & (|A_3|^2 + |A_4|^2) & |a_2|^2 + |a_4|^2 & |b_1|^2 + |b_3|^2 & (1 + \cos^2 \theta_0) & & \\ +3/4 & (|A_3|^2 - |A_4|^2) & |a_2|^2 + |a_4|^2 & |b_1|^2 - |b_3|^2 & (1 + \cos^2 \theta_0) & \cos \theta_2 & \\ +3/4 & (|A_3|^2 - |A_4|^2) & |a_2|^2 - |a_4|^2 & |b_1|^2 + |b_3|^2 & (1 + \cos^2 \theta_0) & & \end{array}$$



$+3/2$	$\sqrt{2}$	$\text{Im}\{A_1^* A_4\}$	$ a_2 ^2 -  a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$
$+3/2$	$\sqrt{2}$	$\text{Im}\{A_1^* A_4\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$
$-3/2$	$\sqrt{2}$	$\text{Im}\{A_2^* A_3\}$	$ a_2 ^2 -  a_4 ^2$	$ b_1 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 + \cos \theta_2)$	$\sin \phi_1$
$-3/2$	$\sqrt{2}$	$\text{Im}\{A_2^* A_3\}$	$ a_2 ^2 -  a_4 ^2$	$ b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$\sin \theta_1$	$(1 - \cos \theta_2)$	$\sin \phi_1$
$+3/2$	$\sqrt{2}$	$\text{Im}\{A_2^* A_4\}$	$ a_4 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
$+3/2$	$\sqrt{2}$	$\text{Im}\{A_2^* A_4\}$	$ a_2 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 + \cos \theta_0)$	$(1 - \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$
$-3/2$	$\sqrt{2}$	$\text{Im}\{A_1^* A_3\}$	$ a_2 ^2$	$ b_1 ^2 -  b_3 ^2$	$\sin \theta_0 (1 - \cos \theta_0)$	$(1 + \cos \theta_1)$	$\sin \theta_2$	$\sin \phi_2$

# Benchmark analysis

resonant  $e^+e^- \rightarrow t\bar{t} \rightarrow bW^+ \bar{b}W^-$   
 $m_b/m_t \rightarrow 0$

analytically at LO

with perfect detector,  $\text{Br} \times \epsilon = 20\%$   
and statistical uncertainties only

$500 \text{ fb}^{-1}$  at  $\sqrt{s} = 500 \text{ GeV}$

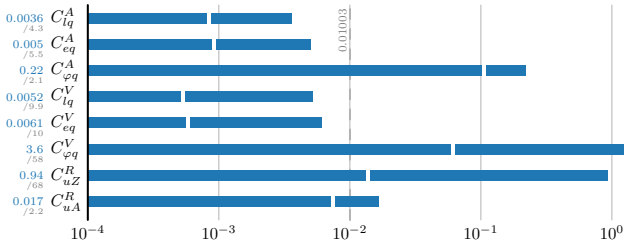
$1 \text{ ab}^{-1}$  at  $\sqrt{s} = 1 \text{ TeV}$

50% with  $P(e^+, e^-) = (+30\%, -80\%)$

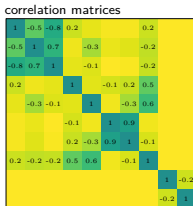
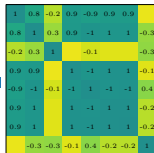
50% with  $P(e^+, e^-) = (-30\%, +80\%)$

# Global constraints

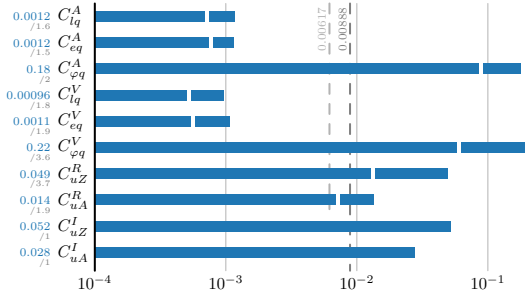
$\sigma + A^{FB}$ :



- in  $\text{TeV}^{-2}$
- white marks: individual constraints
- dashed vertical lines: GDPs
- gray numbers: global/individual ratios



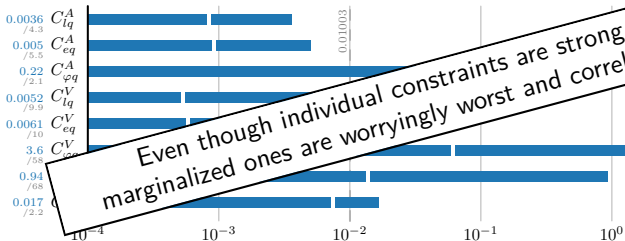
## Statistically optimal observables:



factor of 1.6 improvement  
of the 8-coefficient GDP  
with linear coefficient  
dependence only

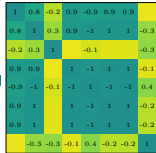
# Global constraints

$\sigma + A^{FB}$ :

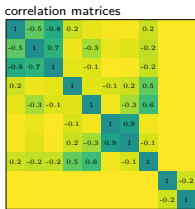
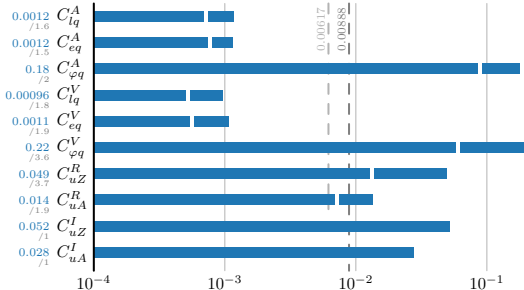


- in  $\text{TeV}^{-2}$
- white marks: individual constraints
- dashed vertical lines: GDPs
- gray numbers: global/individual ratios

Even though individual constraints are strong, marginalized ones are worryingly worst and correlated.



## Statistically optimal observables:

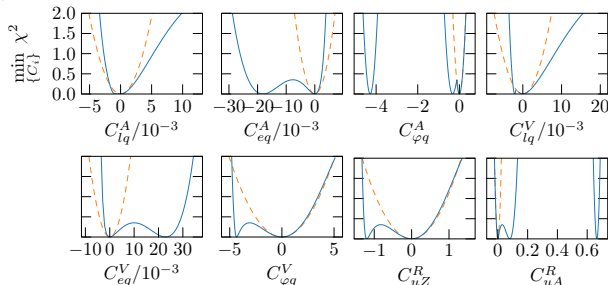


factor of 1.6 improvement of the 8-coefficient GDP with linear coefficient dependence only



# Adding quadratic coefficient dependences

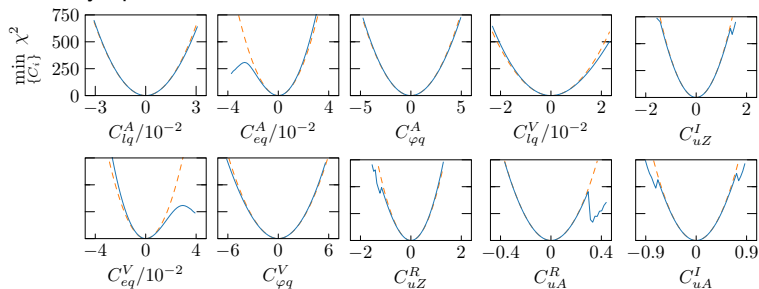
$\sigma + A^{\text{FB}}$ :



dashed: linear dependence only  
plain: linear+quadratic dep.

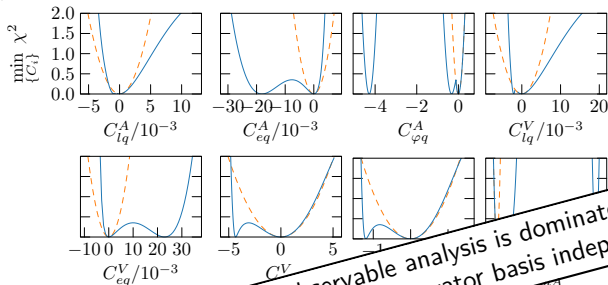
Statistically optimal observables:

Note the vertical scale!



# Adding quadratic coefficient dependences

$\sigma + A^{\text{FB}}$ :

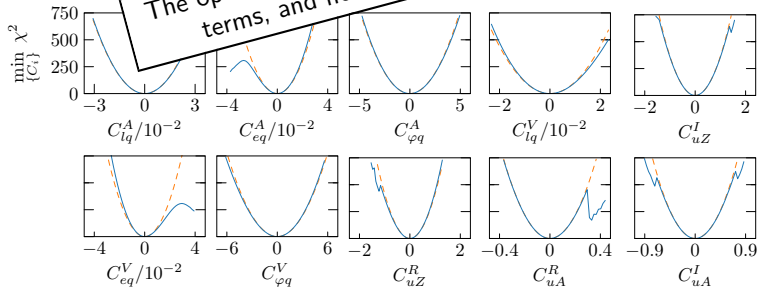


dashed: linear dependence only  
plain: linear+quadratic dep.

Statistically optimal

The optimal-observable analysis is dominated by linear terms, and hence operator basis independent.

Note the vertical scale!



# Getting more realistic

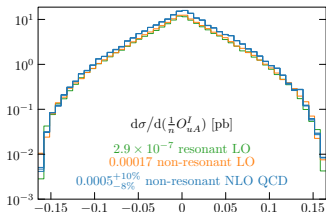
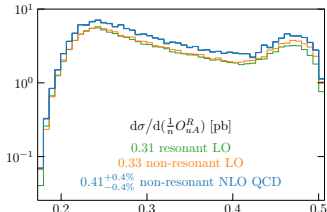
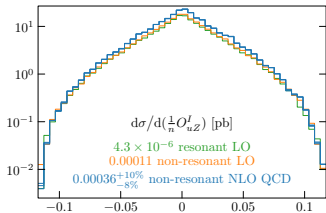
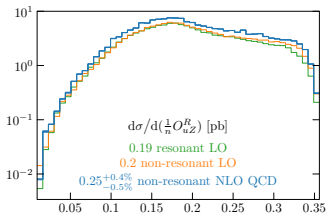
$\sqrt{s} = 500$  GeV,  $P(e^+, e^-) = (+30\%, -80\%)$ ,  
quoted average values of distribution are  $\bar{O}_i/\mathcal{L}$  in pb,  
QCD scale variation from  $m_t/2$  to  $2m_t$

## Theoretical robustness

[similar procedure in Gristan et al, '16]

non-resonant and NLO QCD effects can be studied

e.g.



## Experimental reconstruction

full detector simulations show moderate impact of reconstruction

[details will appear separately]

Run parameters optimization,  
GDP-based

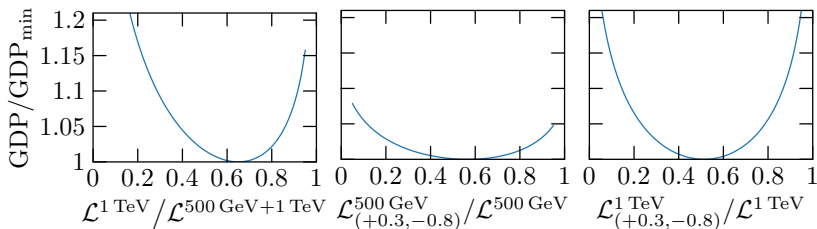
## Examples of run parameters optimization

Given  $1.5 \text{ ab}^{-1}$  to share between two energies and polarizations, the optimal repartition is:

$\sqrt{s} = 500 \text{ GeV}$	$570 \text{ fb}^{-1}$	61% with $P(e^+, e^-) = (+0.3, -0.8)$
$1 \text{ TeV}$	$930 \text{ fb}^{-1}$	52%

→ GDP is 1.02 times better than the benchmark one

for the optimal-observable analysis with all 10 coefficients.



Same performances require  $5.6 \text{ ab}^{-1}$  with only  $\sqrt{s} = 380 + 500 \text{ GeV}$ :

$\sqrt{s} = 380 \text{ GeV}$	$1.7 \text{ ab}^{-1}$	59%	with $P(e^+, e^-) = (+0.3, -0.8)$
$500 \text{ GeV}$	$3.9 \text{ ab}^{-1}$	53%	

## Summary

The EFT parametrizes systematically the theory space in direct vicinity of the standard model.

A global analysis of future-lepton-collider constraints on the top EFT is ongoing.

Statistically optimal observables greatly help covering efficiently the multidimensional phase space.

*Global determinant parameter* ratios assess the strengthening of global constraints, basis independently.