# Cosmological Relaxation of the weak scale

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# The Higgs: Now what? What's Next?

"The experiment worked better than expected and the analysis uncovered a very difficult to find signal"

the words of a string theorist



Profound change in paradigm:

missing SM particle is tool to explore SM and venture into physics landscape beyond

#### Great success...

...but the experimentalists haven't found what the BSM theorists told them they will find in addition to the Higgs boson: no susy, no BH, no extra dimensions, nothing ...



Have the theorists been lying for so many years?

Have the EXP's been too naive to believe the TH's?

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HEP future

exploration/discovery era or consolidation/measurement era?

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Cornell University Library

arXiv.org > physics > arXiv:1503.07735

Physics > Popular Physics

#### **Physics in 100 Years**

Frank Wilczek

(Submitted on 26 Mar 2015)

# SM breakdown

The equations of the [SM] have been tested with far greater accuracy, and under far more extreme conditions, than are required for applications in chemistry, biology, engineering, or astrophysics. While there certainly are many things we don't understand, we do understand the Matter we're made from, and that we encounter in normal life - even if we're chemists, engineers, or astrophysicists (sic: DM!)

## The SM is not free of inadequacies:

Only a description of EW symmetry breaking, not an explanation
 What separates the EW scale from the Planck scale?
 No place for the particle(s) that make up the cosmic DM
 What are the DM particles?
 Does not explain the asymmetry matter-antimatter

Are the conditions realized to allow for EW baryogenesis?

we do not understand the Matter the Universe is made from

Where and how does the SM break down?

— Which machine(s) will reveal (best) this breakdown?

Naturalness (i.e. quantum stability) is the best guide to guess the scale of New Physics

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# Naturalness & TeV scale new physics

Following the arguments of Wilson, 't Hooft (and others): only small numbers associated to the breaking of a symmetry survive quantum corrections ( others are not necessarily theoretically inconsistent but they require some conspiracy at different scales )



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courtesy to N. Craig @ Blois '15

The Higgs mass in the SM doesn't break any (quantum\*) symmetry

\* it does break classical scale invariance, as the running of the gauge couplings does too!

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# Naturalness principle @ work

Following the arguments of Wilson, 't Hooft (and others):

only small numbers associated to the breaking of a symmetry survive quantum corrections

Beautiful examples of naturalness to understand the need of "new" physics see for instance Giudice '13 (and refs. therein) for an account

 $\triangleright$  the need of the positron to screen the electron self-energy:  $\Lambda < m_e/lpha_{
m em}$ 

▶ the rho meson to cutoff the EM contribution to the charged pion mass:  $\Lambda^2 < \delta m_{\pi}^2 / \alpha_{em}$ ▶ the kaon mass difference regulated by the charm quark:  $\Lambda^2 < \frac{\delta m_K}{m_K} \frac{6\pi^2}{G_F^2 f_K^2 \sin^2 \theta_C}$ 

the light Higgs boson to screen the EW corrections to gauge bosons self-energies
 ...

new physics at the weak scale to cancel the UV sensitivity of the Higgs mass?
Apparent fine-tunings have always pointed to new degrees of freedom

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# The Darwinian solution to the Hierarchy

Other origin of small/large numbers according to Weyl and Dirac: hierarchies are induced/created by time evolution/the age of the Universe

Can this idea be formulated in a QFT language?

In which sense is it addressing the stability of small numbers at the quantum level?

Graham, Kaplan, Rajendran '15

- Higgs mass-squared promoted to a field
- The field evolves in time in the early universe and scans a vast range of Higgs mass
- The Higgs mass-squared relaxes to a small negative value
- The electroweak symmetry breaking stops the time-evolution of the dynamical system

## Self-organized criticality

dynamical evolution of a system is stopped at a critical point due to back-reaction

## hierarchies result from dynamics not from symmetries anymore!

important consequences on the spectrum of new physics

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Graham, Kaplan, Rajendran '15

 $\, ar > \,$  slowly rolling field (inflation provides friction) that scans the Higgs mass

 $\Lambda^2 \left( -1 + f\left(\frac{g\phi}{\Lambda}\right) \right) |H|^2 + \Lambda^4 V\left(\frac{g\phi}{\Lambda}\right) + \frac{1}{32\pi^2} \frac{\phi}{f} \tilde{G}^{\mu\nu} G_{\mu\nu}$ Higgs mass potential needed to force depends on  $\phi$  $\phi$  to roll-down in time axion-like coupling (during inflation) that will seed the potential barrier stopping the rolling when the Higgs nd *n* is a positive integer. The first term is rm originate. develops its vev  $\Lambda_{
m QCD}^3 h \cos {\phi \over f}$ e second one corresponds to a Higgs mass time, while such that different values of  $\phi$  scan the Hi pendence on dial have the role  $f_{\rm he}$  have the scale. Finally, the third term plays the role  $\Lambda/q$ 

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Graham, Kaplan, Rajendran '15



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Graham, Kaplan, Rajendran '15



Hierarchy problem solved by light weakly coupled new physics and not by TeV scale physics

a Higgs mass $\phi$ scan the Hi plays the role $\phi$	0 50
<section-header>  PARTICLE PERFECT   WITH ONE SWITCH, EVERYTHING CHANGES   ANNO A LYNKSON AND A LANA AN PARTICLE FYR STON CHANGES AND A LYNKSON KERY ORSTON   MANN A LYNKSON AND A LANA AN PARTICLE FYR STON CHANGES AND A LYNKSON KERY ORSTON   ANNA A LYNKSON AND AND E LANA AN PARTICLE FYR STON CHANGES AND A LYNKSON KERY ORSTON   Particle for MULE AND A LYNKSON AND E LANA AN PARTICLE FYR STON CHANGES AND A LYNKSON KERY ORSTON   Anno Andrea Andre</section-header>	If $\phi$ continues rolling, the Higgs vev increases, the potential barrier increases and ultimately prevents $\phi$ from rolling down further

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Higgs vev stops cosmological rolling

$$\Lambda_{\rm QCD}^3 \frac{v}{f} \sim \frac{\partial}{\partial \phi} \left( \Lambda^4 V(g\phi/\Lambda) \right) \simeq g\Lambda^3$$

**note**: v < A provided that g < 1. It doesn't explain why the coupling is small (that question can be postponed to higher energies) but it ensures that the solution is stable under quantum correction.

ensures that the energy density stored in  $\phi$  does not affect inflation

 $\blacktriangleright$  Classical rolling:  $H_I^3 < g \Lambda^3$ 

Slow rolling:  $H_I > \frac{\Lambda^2}{M_P}$ 

classical displacement over one Hubble time  $\frac{1}{H_{I}}\frac{d\phi}{dt} = \frac{1}{H_{I}^{2}}\frac{dV}{d\phi} = \frac{g\Lambda^{3}}{H_{I}^{2}}$  > quantum fluctuation  $H_{I}$ 

$$\frac{\Lambda^6}{M_P^3} < g\Lambda^3 = \Lambda_{\rm QCD}^3 \frac{v}{f} \qquad \text{ i.e. } \qquad \Lambda < 10^7 \, {\rm GeV} \left(\frac{10^9 \, {\rm GeV}}{f}\right)^{1/6}$$

Important issue:  $\Theta_{QCD} \sim 1 \gg 10^{-10}$ . Can be solved but  $\Lambda < 30$  TeV

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**Classifying relaxing Lagrangians...**  $V(\phi, h) = \Lambda^{3}g\phi - \frac{1}{2}\Lambda^{2}\left(1 - \frac{g\phi}{\Lambda}\right)h^{2} + \epsilon\Lambda_{c}^{4}\left(\frac{h}{\Lambda_{c}}\right)^{n}\cos(\phi/f) + \cdots$ 

▶ n=1: need another source of EWSB

- ▶ QCD condensate <qq>~ Λ<sub>QCD</sub>
- ▶ n=2: no extra source of EWSB needed
  - ▶ quantum stability? h-loops generate extra interactions that will stop  $\phi$ before the Higgs vev develops unless  $\Lambda_c < v$  (coincidence pb and new physics @ TeV again?)  $\epsilon \Lambda_c^4 \cos(\phi/f)$ ,  $\epsilon \Lambda_c^3 g \phi \cos(\phi/f)$

our solution: make the envelop of the oscillatory potential a field

# Cosmological Higgs-Axion Interplay (CHAIN)

Espinosa, Grojean, Panico, Pomarol, Pujolas, Servant'15

$$V(\phi, \sigma, H) = \Lambda^{4} \left(\frac{g\phi}{\Lambda} + \frac{g_{\sigma}\sigma}{\Lambda}\right) - \Lambda^{2} \left(\alpha - \frac{g\phi}{\Lambda}\right) |H|^{2} + \frac{1}{2}\lambda |H|^{4} + A(\phi, \sigma, H) \cos(\phi/f)$$

$$A(\phi, \sigma, H) \equiv \epsilon \Lambda^{4} \left(\beta + c_{\phi} \frac{g\phi}{\Lambda} + \frac{g_{\sigma}\sigma}{\Lambda^{2}} + \frac{|H|^{2}}{\Lambda^{2}}\right)$$
quantum generated
new terms from
the new interaction
the new interaction
the new interaction
the the saves our day
the new interaction
the

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## Same problem, same solution? EX SCALE AS COSMOLOGICAL ERRATIC



okotoks glacial erratic, Alberta, canada

Unnatural large rocks differing in composition from the typical surrounding ones

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## Same problem, same solution? EN SCALE AS COSMOLOGICAL ERRATIC



Unnatural large rocks differing in composition from the typical surrounding ones

Standard geological history:

they were transported by ancient glaciers over hundreds of kilometers

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 $\blacktriangleright$  Quantum stability of the potential  $\,\epsilon \lesssim v^2/\Lambda^2$ 

ensures that terms  $\epsilon^2 \Lambda^4 \cos^2(\phi/f)$  don't affect the tracking solution

Ex.  $cos(\phi/f) = cos(\phi/f) e^{2N4}cos(\phi/f)$ should be subleading compared to  $e^{N^2h^2}cos(\phi/f)$ Requires  $e \lesssim \frac{32^2}{N^2}$ 

courtesy to JR Espinosa

 $\blacktriangleright$  Quantum stability of the potential  $\,\epsilon \lesssim v^2/\Lambda^2$ 

ensures that terms  $\epsilon^2 \Lambda^4 \cos^2(\phi/f)$  don't affect the tracking solution

► Higgs vev stops cosmological rolling  $\frac{\epsilon \Lambda^2 v^2}{f} \sim \frac{\partial}{\partial \phi} \left( \Lambda^4 V(g\phi/\Lambda) \right) \simeq g\Lambda^3$ 

Slow rolling:  $H_I > \frac{\Lambda^2}{M_P}$  ensures that the energy density stored in  $\sigma$  and  $\phi$ does not affect inflation

- ▶ Classical rolling:  $H_I^3 < g\Lambda^3$
- $\phi$  tracks  $\sigma$  in the barrier-free valley before EWSB:  $c_{\phi}g^2 > c_{\sigma}g_{\sigma}^2$
- $\phi$  exits the barrier-free valley after EWSB:  $(c_{\phi} \frac{1}{2\lambda})g^2 < c_{\sigma}g_{\sigma}^2$

to ensure that the Higgs mass scans large field excursions:  $\Delta \phi, \Delta \sigma > \Lambda/g$ values  $\Lambda$  from to the weak scale

 $\left|\frac{\Lambda^3}{M_{\rm Pl}^3} \lesssim g_{\sigma} \lesssim g \lesssim \frac{v^4}{f\Lambda^3}\right| \qquad \Lambda \lesssim \left(v^4 M_{\rm Pl}^3\right)^{1/7} \simeq 2 \times 10^9 \,{\rm GeV}$ 

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 $g = 10g_{\sigma}$ 



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 $g = 10g_{\sigma}$ 



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Nothing to be discovered at the LHC/ILC/CLIC/CepC/SppC/FCC!



only BSM physics below  $\Lambda$ 

## two (very) light and very weakly coupled axion-like scalar fields

$$m_{\phi} \sim (10^{-20} - 10^2) \,\text{GeV}$$
  
 $m_{\sigma} \sim (10^{-45} - 10^{-2}) \,\text{GeV}$ 

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interesting signatures in cosmology and possibly at SHiP

$$\phi \text{ and } \sigma \text{ couple to SM matter via their mixing with the Higgs} \\ \theta_{\phi h} \sim \frac{g\Lambda v}{m_h^2}, \quad \theta_{\sigma \phi} \sim \frac{g_{\sigma} f v^2}{\Lambda^3}, \quad \theta_{\sigma h} \sim \text{Max} \left\{ \theta_{\sigma \phi} \theta_{\phi h}, \frac{g^2}{16\pi^2} \frac{g_{\sigma} \Lambda^7}{f^2 v^3 m_h^2} \right\} \\ \text{from oscillatory potential} & \text{tree-level quantum mixing from } \phi \text{ loop} \\ \text{unsuppressed quartic interaction with the Higgs:} \qquad \phi \phi hh: \quad \epsilon \Lambda^2 / f^2 \\ \phi \text{ and } \sigma \text{ decay to SM particles} \\ \text{(mostly photons in a large region of parameter space)} \\ \Gamma_{\phi} \sim \theta_{\phi h}^2 \Gamma_h(m_{\phi}), \quad \Gamma_{\sigma} \sim \theta_{\sigma h}^2 \Gamma_h(m_{\sigma}) \\ \text{Met Grajean} & \text{Cosmological Relaxation of the used scale 14} & \text{MUST, Jan. 20, 204} \\ \end{array}$$

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vacuum misalignment: (after reheating)

quantum spreading makes the scalars oscillate around their minima

$$\Delta \sigma \sim \Delta \phi \sim \sqrt{N_e} H_I$$

the energy stored in these field oscillations behave like cold DM  $\rho_{\rm ini}^{\sigma} \sim m_{\sigma}^2 (\Delta \sigma)_{\rm ini}^2 \sim H_I^4$  $\rho_{\rm ini}^{\phi} \sim H_I^4$ 

the oscillations start when H~m<sub>i</sub> i.e.  $T_{\rm osc}^i \sim \sqrt{m_i M_{\rm Pl}}$ 

### the energy density is then redshifted till today

 $\Omega_{\sigma} \sim \left(\frac{4 \times 10^{-27}}{a_{\tau}}\right)^{3/2} \left(\frac{\Lambda}{10^8} \,\text{GeV}\right)^{13/2} \qquad \Omega_{\phi} \text{ always very small since } m_{\phi} \gg m_{\sigma} \text{ i.e. } T_{\text{osc}}^{\phi} \gg T_{\text{osc}}^{\sigma}$ 

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 $\phi$  thermal production via interaction with the Higgs  $h + h \to \phi + \phi \qquad \text{or} \qquad SM + SM \to h^{(*)} \to \phi + \phi$ 

single production is subdominant since linear interactions are suppressed by small mixing angle

 $\phi$  almost never in thermal equilibrium (except above  $\Gamma_{BBN}$  line)

number density is obtained from Boltzmann equation

O BBN constraints

 $\bigcirc$  distortions in galactic and extra galactic diffuse X-ray and  $\gamma$ -ray backgrounds

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# Some open questions

▶ large field excursions  Grimet monodromy?	McAllister, Schwaller, Servant, Westphal 'in progress	
non-periodic potential for an axion-like p	Ibanez, Montero, Uranga, Valenzuela '15 article? Gunta Komargodski Perez Ilbaldi '15	
hierarchy of decay constants: F>>f is ~ to	o non-periodid potential	
eternal inflation vs classical evolution?	Choi, Im'15 Kaplan, Rattazzi'15	
Arvanital Iong period of inflation? Riotto et al 'in progress	ci, Dimopoulos, Villadoro 'private communication	
other source of friction to prevent over-	shooting the EW scale? Hardy '15	
UV completion?		
weak gravity conjecture? Heidenreich, Reece, Ru	delius '15 Hebercker, Rompineve, Westphal '15	
Can other scales be relaxed too? SUSY breaking scale?		
signatures in atomic physics?	Datell, Gludice, McCullough 15	
A new playground for model builders at the cross-road between exp/cosmo/pheno/strings Joined forces needed		

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Conclusions

Weither the second s

Now, existence proof that technical naturalness doesn't require new physics at the weak scale.

Is technical naturalness the right criterion?

The energy frontier might be different than what we thought for many years!

let us think further and be prepared to be surprised