



Walter Burke Institute for
Theoretical Physics

On the Singularity Structure of Maximally Supersymmetric Scattering Amplitudes

Jaroslav Trnka
Caltech

N. Arkani-Hamed, J. Bourjaily, F. Cachazo, JT, 1410.0354
Z. Bern, E. Herrmann, S. Litsey, J. Stankowicz, JT, to appear

Motivation

- Goal: New mathematical structures in QFT.
- Ideal test objects: **on-shell scattering amplitudes**.
- Standard approach
 1. Integrand of the amplitude.
 2. Integrated expression.
 3. Cross sections,...
- New formulation of QFT \rightarrow most visible in the **integrand**.
- $\mathcal{N} = 4$ SYM : harmonic oscillator of 21th century.

Overview

- Amplitudes in planar $\mathcal{N} = 4$ SYM
 - ▶ New structures and symmetries, dual formulation.
 - ▶ **Logarithmic singularities.**
- Amplitudes in full $\mathcal{N} = 4$ SYM
 - ▶ No symmetries, no dual formulation,...
 - ▶ Evidence for 4pt: logarithmic singularities.

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Conjecture: [Arkani-Hamed, Bourjaily, Cachazo, JT]

Four point amplitudes in $\mathcal{N} = 4$ SYM have only logarithmic singularities and no poles at infinity.

- In the planar sector: it implies dual conformal invariance.

[Bern, Herrmann, Litsey, Stankowicz, JT]

Planar $\mathcal{N} = 4$ SYM

- Unitary methods successful: high loop results, BDS,...
[Bern, Dixon, Kosower, Roiban, Carrasco, Johansson, Smirnov,...]
- Recursion relations
[Britto, Cachazo, Feng, Witten; Arkani-Hamed, Bourjaily, Cachazo, Caron-Huot, JT]
- Amplitudes at strong coupling
[Alday, Maldacena]
- Dual conformal symmetry
[Drummond, Henn, Korchemsky, Smirnov, Sokatchev; Alday, Maldacena]
- Yangian symmetry
[Drummond, Henn, Plefka]

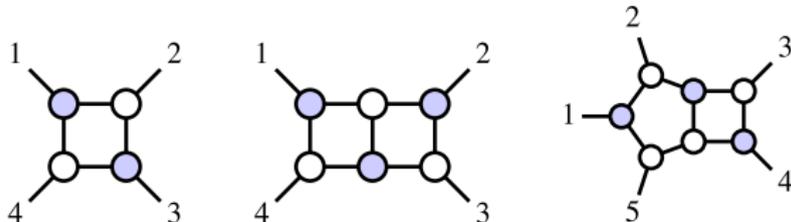
Planar $\mathcal{N} = 4$ SYM

- Amplitudes/super-Wilson loops correspondence
[Caron-Huot; Mason, Skinner]
- Gamma cusp to all orders
[Beisert, Eden, Staudacher]
- OPE and flux-tube S-matrix
[Basso, Sever, Vieira]
- Integrals, symbols, polylogs
[Henn, Smirnov²,...; Golden, Goncharov, Paulos, Spradlin, Volovich]
- Hexagon bootstrap
[Dixon, Drummond, Duhr, Henn, von Hippel, Pennington]
- Many more....

Dual formulation

[Arkani-Hamed, Bourjaily, Cachazo, Goncharov, Postnikov, JT]

- On-shell physics
 - ▶ Feynman diagrams: off-shell physics.
 - ▶ Unitary methods: on-shell data.
 - ▶ **On-shell diagrams**: on-shell objects.



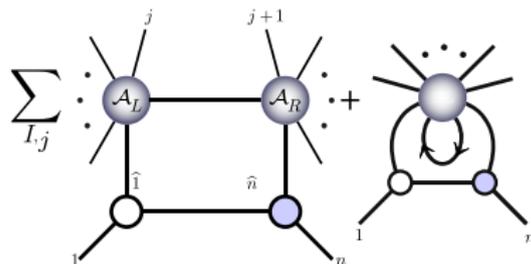
- On-shell diagrams in planar $\mathcal{N} = 4$ SYM
 - ▶ Cells in Positive Grassmannian $G_+(k, n)$.
 - ▶ Logarithmic form

$$\Omega_0 = \int \frac{dx_1}{x_1} \dots \frac{dx_d}{x_d} \delta(C(x_i)Z_j)$$

Dual formulation

[Arkani-Hamed, Bourjaily, Cachazo, Goncharov, Postnikov, JT]

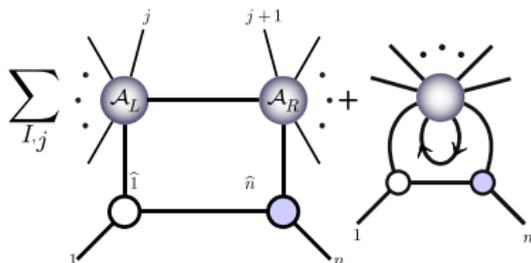
- BCFW recursion relation
 - ▶ Sum of on-shell diagrams.
 - ▶ Yangian invariant term-by-term.



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- **Amplituhedron** [Arkani-Hamed, JT]
 - ▶ Gluing on-shell diagrams together.
 - ▶ Geometric definition of the amplitude:
 - ▶ Defined by the set of inequalities.
 - ▶ Form with logarithmic singularities.
- Crucial property: Logarithmic singularities.

Logarithmic singularities

- Definition: Differential form $\Omega \sim \frac{dx}{x} \tilde{\Omega}$ near $x = 0$.

$$\Omega = \frac{dx}{x} = d \log x \quad \text{vs} \quad \Omega = \frac{dx}{x^2}$$

- ▶ Multiple poles hidden in the cut structure

$$\Omega = \frac{dx dy}{xy(x+y)} \quad \text{Res}_{x=0} \Omega = \frac{dy}{y^2}$$

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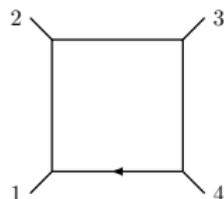
$$\Omega = \frac{dx dy}{xy(x+y)} \quad \text{Res}_{x=0} \Omega = \frac{dy}{y^2}$$

- **Dlog form:** $\Omega = d \log f_1 d \log f_2 \dots d \log f_m$

$$\Omega = \frac{dx dy (x-y)}{xy(x+y)} = d \log \frac{x}{(x+y)} d \log \frac{y}{(x+y)}$$

Singularities of loop integrals

- ▶ Example: box integral



$$I = \frac{d^4 \ell \, st}{\ell^2 (\ell + k_1)^2 (\ell + k_1 + k_2)^2 (\ell - k_4)^2}$$

- ▶ Examples of integrals with non-logarithmic singularities:

$$I = \frac{d^4 \ell}{(\ell^2)^2 (\ell + k_1)^2 (\ell + k_1 + k_2)^2}, \quad I = \frac{d^4 \ell}{\ell^2 (\ell + k_1)^2 (\ell + k_2)^2 (\ell + k_3)^2}$$

- ▶ At higher loops: multiple poles
→ Special **numerator** needed to cancel them.

Poles at infinity

- We consider loop integrals.
- Type of singularities: logarithmic.
- Restriction on positions of singularities:

No poles for $\ell \rightarrow \infty$.

Poles at infinity

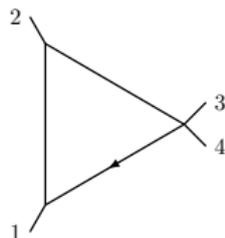
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- Dual conformal symmetry
 - ▶ No infinity twistor \rightarrow no poles at infinity.

Poles at infinity

- Example: triangle integral

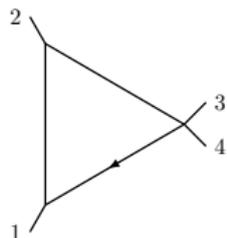


$$I = \frac{d^4\ell \, s}{\ell^2(\ell + k_1)^2(\ell + k_1 + k_2)^2}$$

- ▶ Triple cut: $\ell^2 = (\ell + k_1)^2 = (\ell + k_1 + k_2)^2 = 0$
- ▶ Solution: $\ell - k_1 = \alpha\lambda_1\widetilde{\lambda}_2$

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- ▶ Solution: $\ell - k_1 = \alpha \lambda_1 \widetilde{\lambda}_2$
- ▶ Residue on this cut:
$$I = \frac{d\alpha}{\alpha}$$
- ▶ Pole for $\alpha \rightarrow \infty$ which implies $\ell \rightarrow \infty$.

Planar amplitudes

- Dual formulation using on-shell diagrams:
 - ▶ No poles at infinity.
 - ▶ Logarithmic singularities in the Grassmannian space.
 - ▶ N^k MHV for $k = 0, 1, 2$: logarithmic singularities in **momentum space**.

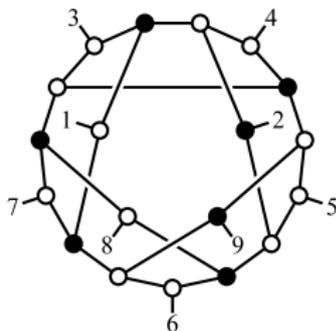
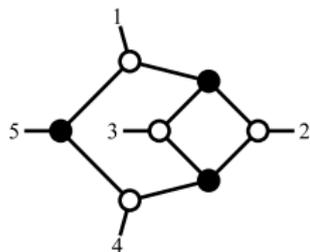
Non-planar amplitudes

- Amplitudes in complete $\mathcal{N} = 4$ SYM
 - ▶ Still have: maximal supersymmetry, UV finiteness.
 - ▶ But: no DCI, no Yangian, no amplitudes/Wilson loop

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 - ▶ Still have: maximal supersymmetry, UV finiteness.
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- On-shell diagrams well defined



- ▶ Even richer mathematical structure.
[Arkani-Hamed, Bourjaily, Cachazo, Postnikov, JT]
- ▶ Logarithmic singularities, no poles at infinity.

Non-planar amplitudes

- No dual formulation for amplitudes yet.
- Suppose we can find it!
- Natural conjecture:

Non-planar amplitudes

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- Natural conjecture:

Amplitudes in complete $\mathcal{N} = 4$ SYM have **logarithmic** singularities and **no poles at infinity**.

- ▶ This is in the Grassmannian space.
- ▶ In momentum space conservative conjecture: **four point**.
- No $1/N$ expansion, property of the full theory.

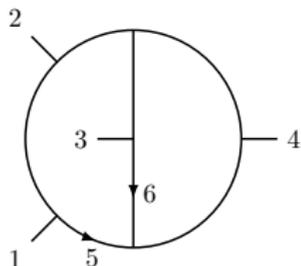
Non-planar amplitudes

- How to test the conjecture?
 - ▶ Analyze data up to 5-loops.
[Bern, Carrasco, Dixon, Johansson, Kosower, Roiban]
 - ▶ Find the new basis with these two properties **manifest**.
 - ▶ Expand the amplitude in this basis.

Non-planar amplitudes

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- Conservative strategy: scalar integrals.
 - ▶ Denominator given by the diagram.
 - ▶ Fix the **numerator**: cancels bad singularities.



Non-planar amplitudes

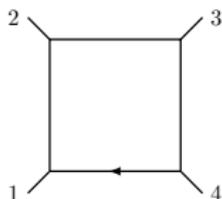
- Existence of such numerators not guaranteed.
 - ▶ Possible cancellations between diagrams.
- New expansion vs. reference
 - ▶ No unique integrand: no algebraic proof.
 - ▶ Match on cuts.

Non-planar amplitudes

- Existence of such numerators not guaranteed.
 - ▶ Possible cancellations between diagrams.
- New expansion vs. reference
 - ▶ No unique integrand: no algebraic proof.
 - ▶ Match on cuts.
- Explicitly constructed and checked up to 3-loops.
- Few more checks at higher loops.

One-loop amplitude

- One-loop amplitude: sum over permutations of box

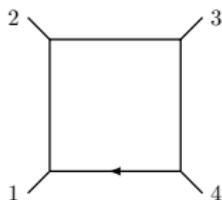


$$= \frac{d^4 \ell \, st}{\ell^2 (\ell + k_1)^2 (\ell + k_1 + k_2)^2 (\ell - k_4)^2}$$

- Same structure as the planar amplitude.

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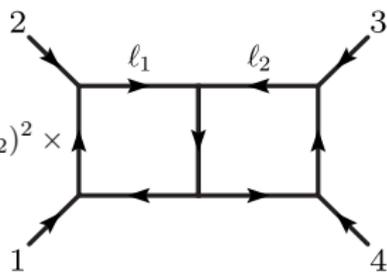
- Same structure as the planar amplitude.
- Dlog form for box

$$\text{dlog} \frac{\ell^2}{(\ell - \ell^*)^2} \text{dlog} \frac{(\ell + k_1)^2}{(\ell - \ell^*)^2} \text{dlog} \frac{(\ell + k_1 + k_2)^2}{(\ell - \ell^*)^2} \text{dlog} \frac{(\ell - k_4)^2}{(\ell - \ell^*)^2}$$

Two-loop amplitude

- Planar double-box

$$\mathcal{I}_{1,2,3,4}^{(P)} \equiv (p_1 + p_2)^2 \times$$



Two-loop amplitude

- Planar double-box

$$\mathcal{I}_{1,2,3,4}^{(P)} \equiv (p_1 + p_2)^2 \times$$

- written in the dlog form

$$st\mathcal{I}_{1234}^{(P)} = \text{dlog } \alpha_1 \text{ dlog } \alpha_2 \text{ dlog } \alpha_3 \dots \text{dlog } \alpha_8$$

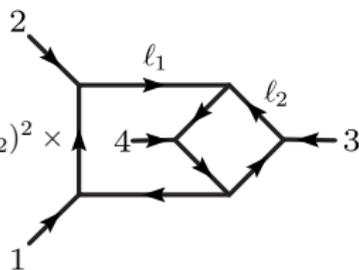
$$\begin{aligned} \alpha_1 &\equiv \ell_1^2 / (\ell_1 - \ell_1^*)^2, & \alpha_5 &\equiv \ell_2^2 / (\ell_2 - \ell_2^*)^2, \\ \alpha_2 &\equiv (\ell_1 - p_2)^2 / (\ell_1 - \ell_1^*)^2, & \alpha_6 &\equiv (\ell_1 + \ell_2)^2 / (\ell_2 - \ell_2^*)^2, \\ \alpha_3 &\equiv (\ell_1 - p_1 - p_2)^2 / (\ell_1 - \ell_1^*)^2, & \alpha_7 &\equiv (\ell_2 - p_3)^2 / (\ell_2 - \ell_2^*)^2, \\ \alpha_4 &\equiv (\ell_1 + p_3)^2 / (\ell_1 - \ell_1^*)^2, & \alpha_8 &\equiv (\ell_2 - p_3 - p_4)^2 / (\ell_2 - \ell_2^*)^2, \end{aligned}$$

Two-loop amplitude

- Non-planar double box

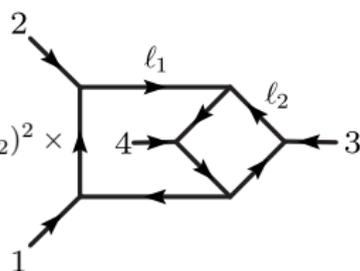
$$\mathcal{I}_{1,2,3,4}^{(NP)}$$

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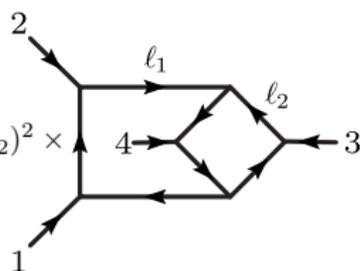
- Generate a **double pole**

- ▶ Quadruple cut on ℓ_2 .
- ▶ Triple cut on $\ell_1 = xp_2$.

$$\text{Res } \mathcal{I} = \frac{dx}{(x+1)x^2tu}$$

Two-loop amplitude

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- Numerator: $N_{old} = (p_1 + p_2)^2 \rightarrow N_{new} = (\ell_1 + p_3)^2 + (\ell_1 + p_4)^2$
 - ▶ On the cut: $N_{new} = (xp_2 + p_3)^2 + (xp_2 + p_4)^2 = -xs$

Two-loop amplitude

- More double poles: N_{new} cancels all of them.
- Also conditions on the absence of the pole at infinity.
 - ▶ Everything resolved by N_{new} .

Two-loop amplitude

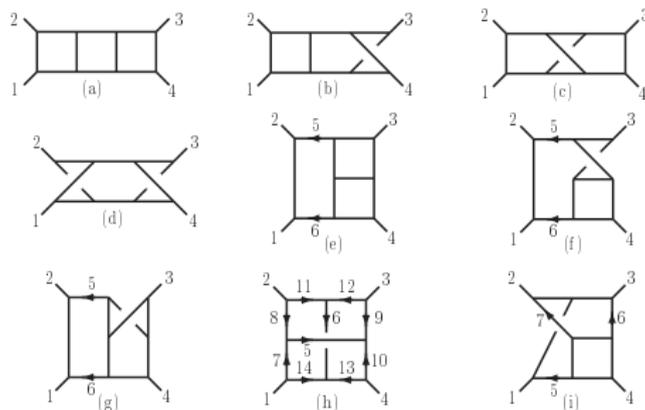
- More double poles: N_{new} cancels all of them.
- Also conditions on the absence of the pole at infinity.
 - ▶ Everything resolved by N_{new} .
- Expand the amplitude in a new basis: YES.
- New result looks differently.
 - ▶ Difference cancels due to color **Jacobi identity**.

$$f^{abe} f^{cde} + f^{ace} f^{dbe} + f^{ade} f^{bce} = 0$$

Three-loop amplitude

- Nine diagrams in the basis.

[Bern, Carrasco, Dixon, Johansson, Kosower, Roiban, 2007]

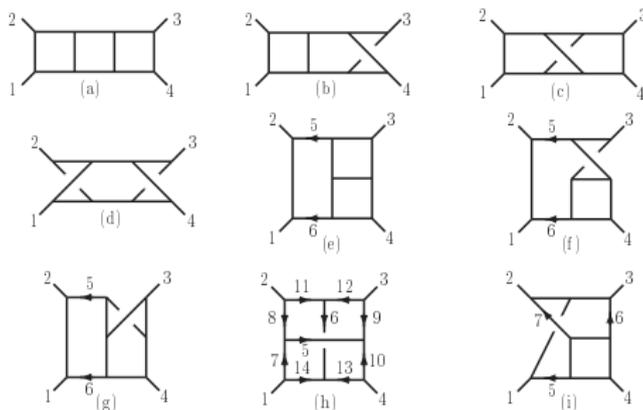


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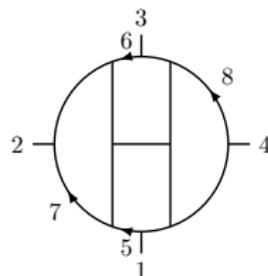


- All except (a) double poles: **new numerators**.
- Check completed: Amplitude expanded in new basis!
→ Talk by Enrico Herrmann

Higher loops

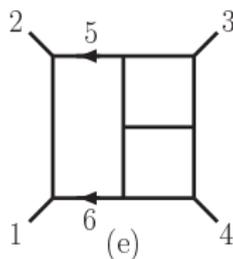
- Partial checks
 - ▶ New basis for some diagrams.
 - ▶ Match the reference on maximal cuts.

- All checks show that the conjecture is correct.



Back to planar sector

- In planar sector
 - ▶ Both properties manifest in the basis of on-shell diagrams.
- What about scalar integrals?
 - ▶ Do not consider DCI.
 - ▶ Follow only our two conditions.

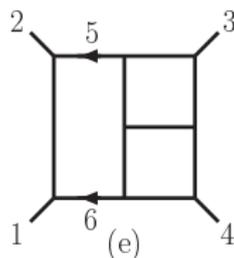


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- Result:

- ▶ In all cases we reproduce DCI numerators.
- ▶ We get even **stronger** restrictions!

Beyond DCI

- Extra condition beyond DCI [Drummond, Korchemsky, Sokatchev]
 - ▶ Starting at 4-loops some DCI integrals **not well-defined**.
 - ▶ All dual loop momenta x_5, x_6, \dots go to external point.

$$\rho^2 = x_{35}^2 + x_{36}^2 + \dots \rightarrow 0$$

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$$\rho^2 = x_{35}^2 + x_{36}^2 + \dots \rightarrow 0$$

- ▶ If the integral behaves like

$$I \sim \int \frac{d\rho}{\rho} \quad \text{IR divergence even off-shell}$$

- ▶ Zero coefficient in the amplitude.

Beyond DCI

- Equivalent to the presence of certain type of **multiple poles!**
- There are more rules of this type
 - ▶ Dual loop variables null separated from each other.

$$\rho^2 = x_{56}^2 + x_{57}^2 + \dots \rightarrow 0$$

- ▶ Null separated from two points

$$\rho^2 = x_{25}^2 + x_{35}^2 + x_{26}^2 + x_{36}^2 + \dots \rightarrow 0$$

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$$\rho^2 = x_{25}^2 + x_{35}^2 + x_{26}^2 + x_{36}^2 + \dots \rightarrow 0$$

- We have data up to 7-loops.

[Bern, Carrasco, Dixon, Johansson, Kosower, Smirnov; Bourjaily, DiRe, Shaikh, Spradlin, Volovich; Eden, Heslop, Korchemsky, Sokatchev]

- These rules explain zeroes in 4-loop and 5-loop expansions.

Beyond DCI

- Many more types of multiple poles.
- Requirement: All of them must be absent.
 - ▶ Explains many zeroes of DCI integrals at 6-loops and 7-loops.
 - ▶ Gives correct relative coefficients at 5-loops.
 - Talk by Enrico Herrmann

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- Cancellation of multiple poles at higher loops.
 - ▶ Possible in each diagram?
 - ▶ Between different diagrams?

Beyond DCI

- Multiple poles from [Drummond, Korchemsky, Sokatchev]
 - ▶ **Real momenta** on the solutions for these cuts.
 - ▶ Problem: they are in the domain of integration.

Beyond DCI

- Multiple poles from [Drummond, Korchemsky, Sokatchev]
 - ▶ **Real momenta** on the solutions for these cuts.
 - ▶ Problem: they are in the domain of integration.
- Other multiple poles in integrals:
 - ▶ Momenta are complex \rightarrow outside integration domain.
 - ▶ Integral well-defined but non-uniform transcendentality.
- Conjecture for $\mathcal{N} = 4$ SYM: no multiple poles at all.

Beyond DCI

- Planar and non-planar integrals:
 - ▶ Fixed by the **same rules**.
- In planar sector:
 - ▶ Integrals with log singularities: automatically DCI.

Beyond DCI

- Planar and non-planar integrals:
 - ▶ Fixed by the **same rules**.
- In planar sector:
 - ▶ Integrals with log singularities: automatically DCI.
- It suggests the existence of the non-planar analogue of DCI!

Beyond DCI

- Planar sector: logarithmic singularities and no poles at infinity

$$I = \sum_j \alpha_j (\text{dlog} \dots \text{dlog})_j$$

Beyond DCI

- Planar sector: logarithmic singularities and no poles at infinity

$$I = \sum_j \alpha_j (\text{dlog} \dots \text{dlog})_j$$

- Two things can go wrong
 1. Arguments of dlogs are non-DCI (without poles at infinity).
 2. Coefficients α_j are non-DCI (mod overall constant).
- Explicit data: None of it happens.

Supergravity

- Extend ideas to other theories:
 - ▶ Lower supersymmetry: new singularity structure.
 - ▶ From Yang-Mills to **gravity**.

Supergravity

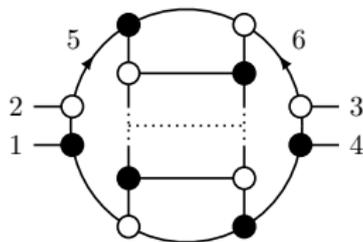
- Extend ideas to other theories:
 - ▶ Lower supersymmetry: new singularity structure.
 - ▶ From Yang-Mills to **gravity**.
- Maximal $\mathcal{N} = 8$ SUGRA
 - ▶ Singularity structure and UV behavior unclear.
 - ▶ Close relation to $\mathcal{N} = 4$ SYM via BCJ relations.
[Bern, Carrasco, Johansson]
 - ▶ Natural idea: explore singularity structure of the integrand.
- At 2-loops logarithmic singularities and no poles at infinity.

Supergravity

- No poles at infinity \rightarrow **UV finiteness**.
- For $\mathcal{N} = 4$ SYM: integrand-based derivation of UV finiteness.
- If true for $\mathcal{N} = 8$ SYM: trivially UV finite as well.

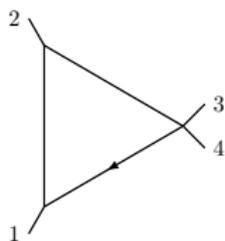
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- For $\mathcal{N} = 4$ SYM: integrand-based derivation of UV finiteness.
- If true for $\mathcal{N} = 8$ SYM: trivially UV finite as well.
- Explicit checks of poles at infinity
 - ▶ No poles at 1-loop and 2-loops.
 - ▶ Logarithmic at 3-loops.
 - ▶ Non-logarithmic at 4-loops,
- Results: Poles at infinity **are present**.



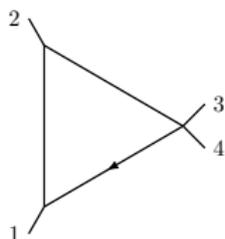
Supergravity

- This is **NOT** a proof of UV divergence.
 - ▶ No poles at infinity \rightarrow UV finiteness.
 - ▶ UV finiteness $\not\rightarrow$ No poles at infinity.
 - ▶ Simplest example: triangle integral.



Supergravity

- This is **NOT** a proof of UV divergence.
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- Possible finiteness of $\mathcal{N} = 8$ SUGRA:
 - ▶ Must depend on **detailed structures** of poles at infinity.
 - ▶ More complicated than $\mathcal{N} = 4$ SYM.
 - ▶ Cancellations between diagrams \rightarrow Talk by Scott Davies.

Conclusion

- Amplitudes in complete $\mathcal{N} = 4$ SYM: singularities of integrand.
- Inspired by properties of on-shell diagrams:

Conjecture: Four point amplitudes in $\mathcal{N} = 4$ SYM have only logarithmic singularities and no poles at infinity.

- Explicit checks up to 3-loops.
- In planar sector: stronger conditions than DCI.
- $\mathcal{N} = 8$ SUGRA: poles at infinity
→ Detailed knowledge needed to study UV behavior.

Conclusion

- Amplitudes in complete $\mathcal{N} = 4$ SYM: singularities of integrand.
- Inspired by properties of on-shell diagrams:

Conjecture: Four point amplitudes in $\mathcal{N} = 4$ SYM have only logarithmic singularities and no poles at infinity.

- Explicit checks up to 3-loops.
- In planar sector: stronger conditions than DCI.
- $\mathcal{N} = 8$ SUGRA: poles at infinity
→ Detailed knowledge needed to study UV behavior.

Thank you for the attention!